



DLR

Deutsches Zentrum
für Luft- und Raumfahrt
German Aerospace Center

Institute of Aerodynamics
and Flow Technology

Towards Forced Eddy Simulation for Installed Nozzle-Wing Configurations

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Svetlogorsk, Russia



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Knowledge for Tomorrow



Aim of the Work

Installed Jet Noise Simulation

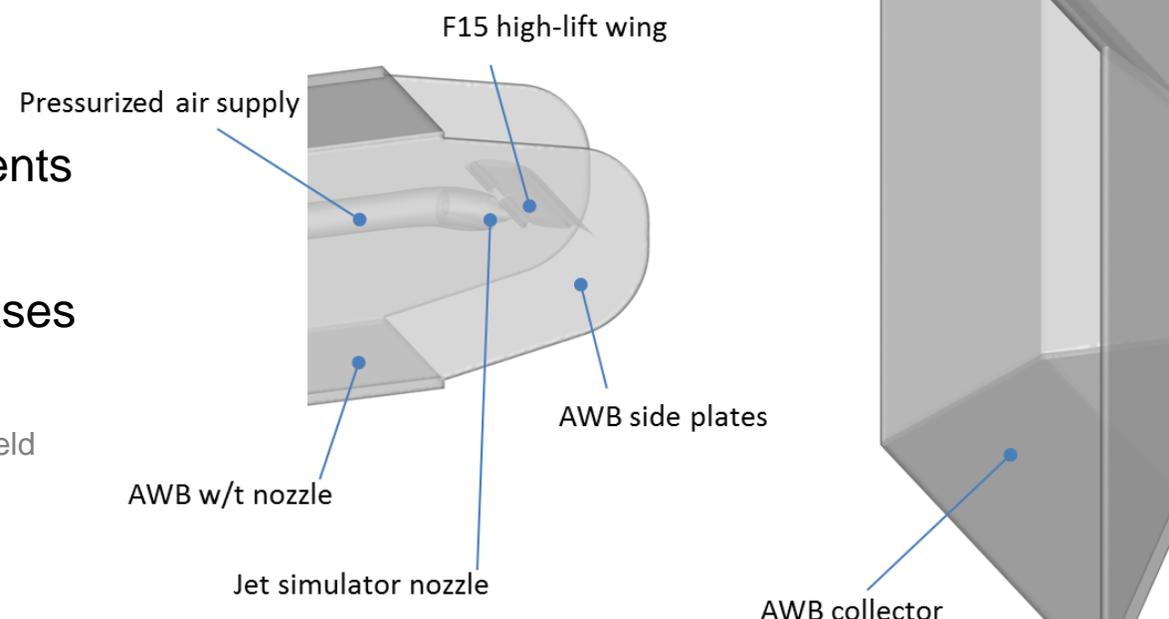
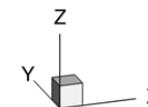
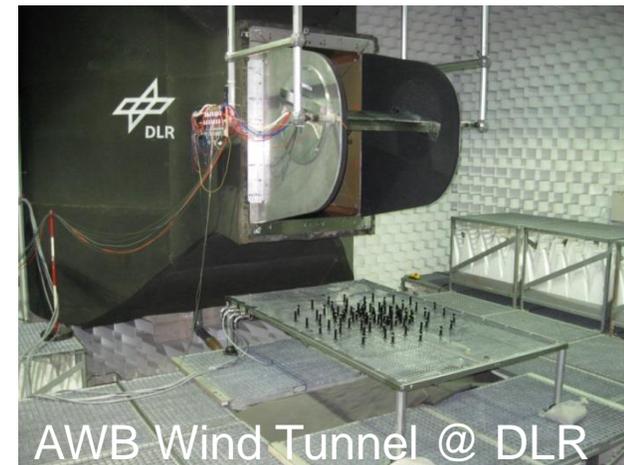
- Direct Noise Computation (DNC¹) of a jet-flap installation noise wind tunnel setup tested in the Acoustic Wind Tunnel Braunschweig (AWB) at DLR
- DNC¹ using a zonal RANS/VLES² approach based on Non-Linear Disturbance Equations (NLDE)³
 - Zonal VLES of the entire wind tunnel setup using a standard CFD/CAA framework
 - Evaluation of the potential to lower resolution requirements using an active stochastic backscatter model
 - Transfer of stochastic backscatter from DHIT⁴ to use cases

¹Bailly, Bogey & Marsden '10: solving the compressible Navier-Stokes equations to determine simultaneously the aerodynamic field and the acoustic field

²Very-large-eddy simulation: filter and grid are too coarse to resolve 80% of the energy (Pope 2000)

³Morris '97, Sagaut & Labourasse '02, Terracol et al. '06, Batten et al. '04

⁴Decaying Homogeneous Isotropic Turbulence



Outline

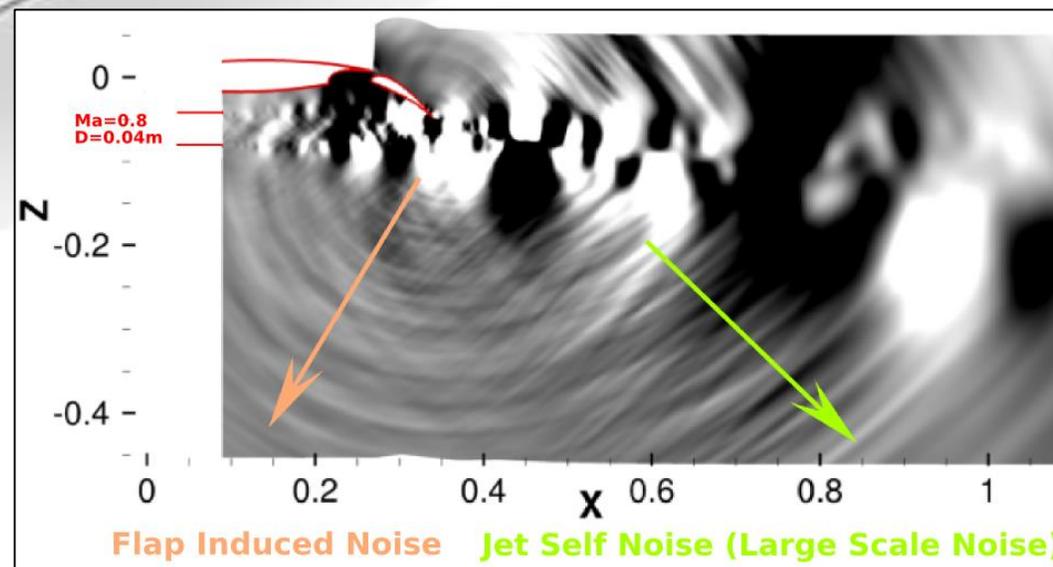
- **Simulation Problem**
- Methodology
- Results of cold isolated single stream jet
- Results of installed UHBR configuration
- Summary & Conclusions



Simulation Problem

UHBR Installation Noise

- Noise prediction for configurations with nozzle mounted closely to wing
- The installation noise from jet-airframe interaction potentially is one of the prominent noise source
- This noise generation mechanism is attributed to coherent hydrodynamic fluctuations passing by the airfoil trailing edge
- Qualitatively different noise generation mechanisms as can be observed in directivity pattern:
 - Airframe → maximum to forward arc
 - Jet Noise → maximum to rearward arc



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Methodology: Zonal RANS/VLES Approach

NLDE⁴ (Non-Linear Disturbance Equations)

⁴Morris '97, Sagaut & Labourasse '02, Terracol et al. '06, Batten et al. '04

- Navier-Stokes equations for primitive variables:

$$\frac{\partial \mathbf{U}}{\partial t} + \mathcal{N}(\mathbf{U}) = \mathbf{0} \quad \mathbf{U} = \begin{pmatrix} \rho \\ v_i \\ p \end{pmatrix}$$

- Triple decomposition of variables:

$$\mathbf{U} = \underbrace{\mathbf{U}^0 + \bar{\mathbf{U}}'}_{\bar{\mathbf{U}}} + \mathbf{U}''$$

Labels in diagram: RANS mean-flow (points to \mathbf{U}^0), resolved flow (points to $\bar{\mathbf{U}}$), resolved fluctuations (points to $\bar{\mathbf{U}}'$), subgrid fluctuations (points to \mathbf{U}'')

- NLDE Variant 1:

$$\frac{\partial \bar{\mathbf{U}}'}{\partial t} + \mathcal{N}(\bar{\mathbf{U}}) = \mathbf{f}_{sgs}(\mathbf{U}'')$$

- NLDE Variant 2:

$$\frac{\partial \bar{\mathbf{U}}'}{\partial t} + \mathcal{N}'(\mathbf{U}^0, \bar{\mathbf{U}}') = \mathbf{f}_{visc}^0 + \mathbf{f}_{sgs}$$

$$\mathcal{N}(\bar{\mathbf{U}}) = \mathcal{N}'(\mathbf{U}^0, \bar{\mathbf{U}}') + \mathcal{N}^0(\mathbf{U}^0)$$

$$\mathcal{N}^0(\mathbf{U}^0) = -\mathbf{f}_{visc}^0(\mathbf{U}^0)$$

≡RANS



Methodology

NLDE (Non-Linear Disturbance Equations)

- NLDE Variant 1:

$$\frac{\partial \bar{U}'}{\partial t} + \mathcal{N}(\bar{U}) = f_{sgs}(U'')$$

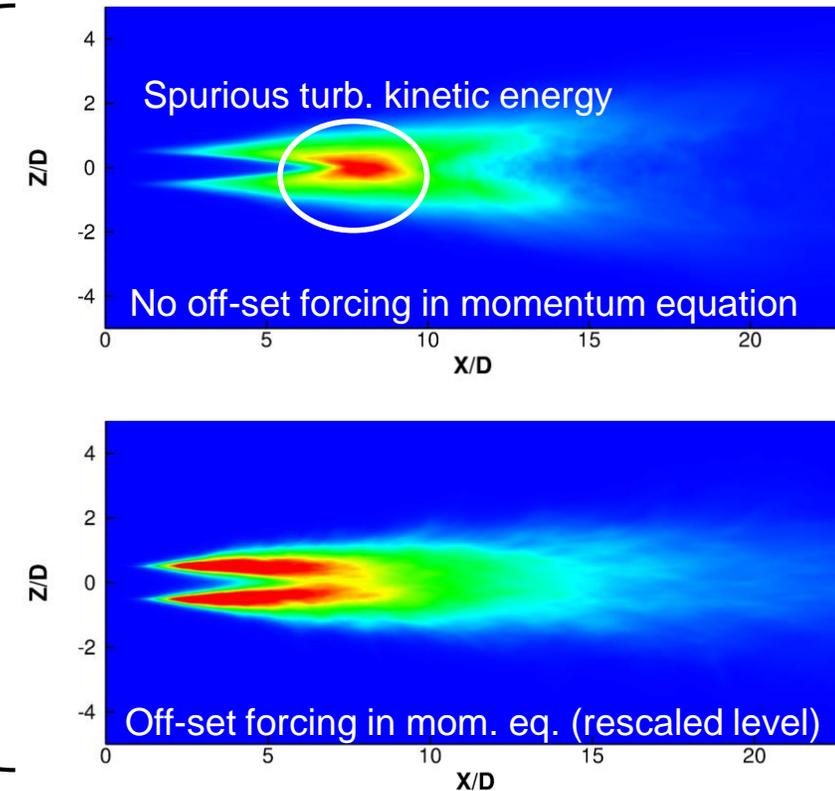
off-set forcing w/o forcing

- NLDE Variant 2:

$$\frac{\partial \bar{U}'}{\partial t} + \mathcal{N}'(U^0, \bar{U}') = f_{visc}^0 + f_{sgs}$$

explicit off-set forcing

- Constant off-set necessary to avoid spurious turbulence in single stream jet
- “Variant 1” used for momentum equation (“off-set forcing w/o forcing”)
- “Variant 2” used for density + pressure equations (no off-set forcing)
- Background RANS flow from DLR CFD solver TAU
- NLDE realized with DLR code PIANO
 - 4th order DRP scheme of Tam & Webb
 - High-order (HO) filter to remove spurious waves
 - Optimized 4th order Runge-Kutta time integration



NLDE + Stochastic Backscatter = FES (Forced Eddy Simulation)

$$\mathbf{f}_{sgs} = - \underbrace{\nabla \times (\bar{\rho} \nu_r \boldsymbol{\omega}')}_{\mathbf{f}_D^{(5)}} + \underbrace{\nabla \times \mathbf{q}}_{\mathbf{f}_F}$$

r.h.s. vector force dissipation stochastic forcing

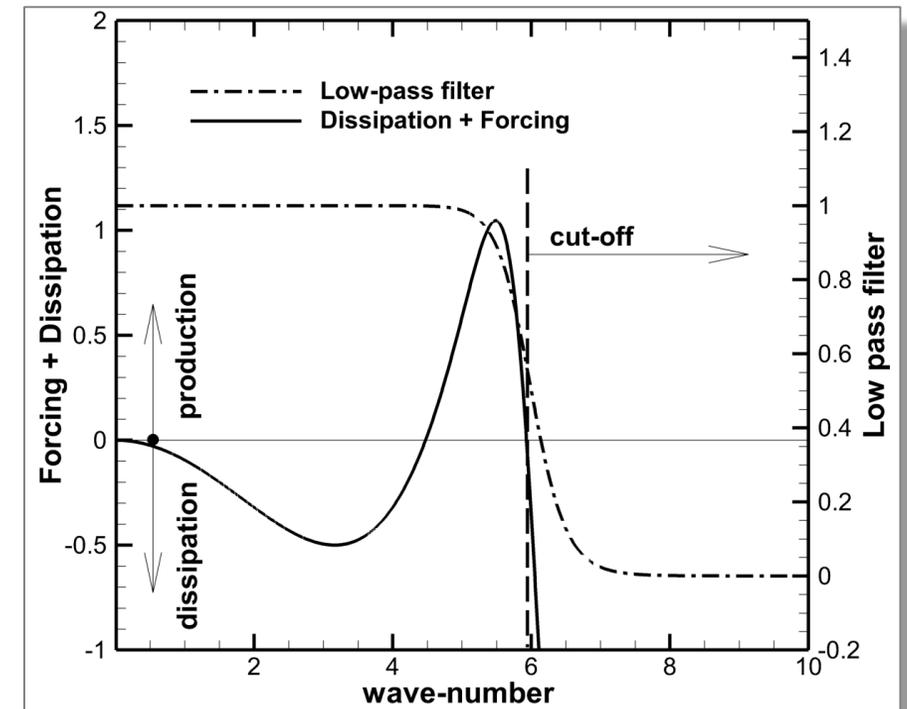
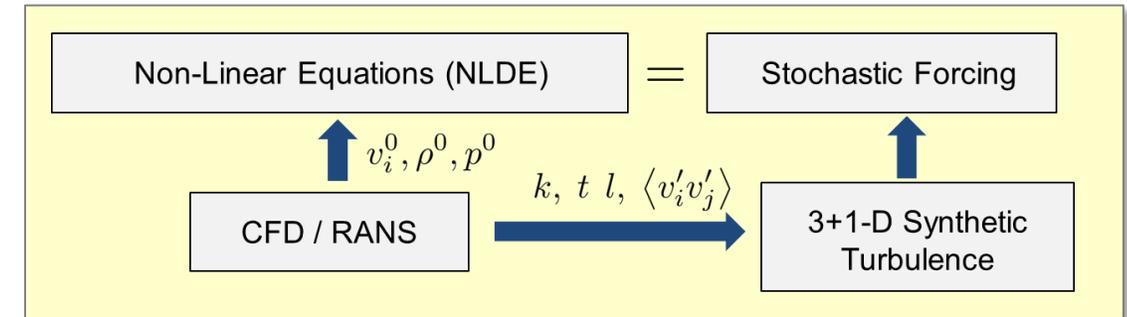
- Smagorinsky-like vector-force sgs-model⁴ of fluctuations
- High-order filtered equations + sgs-model: mixed model⁵
- 3+1-D stochastic forcing provides turbulent backscatter^{1,2,3} from non-resolved scales

⁴Dantinne, Jeanmart, Winkelmanns, Legat, Carati, Applied Scientific Research (59), 1998

⁵Bardina, Ferziger, and Reynolds, 1980

¹C.E. Leith, Physics of Fluids A 2, 297, ²U. Schumann, Proc. R. Soc. London A (1995),

³Zamansky et al., J. of Turb. 11, (2010).



NLDE + Stochastic Backscatter = FES (Forced Eddy Simulation)

$$\mathbf{f}_{sgs} = - \underbrace{\nabla \times (\bar{\rho} \nu_r \boldsymbol{\omega}')}_{\mathbf{f}_D^{(5)}} + \underbrace{\nabla \times \mathbf{q}}_{\mathbf{f}_F}$$

r.h.s. vector force dissipation stochastic forcing

- Backscatter forcing features as discussed in literature

(Kraichnan 1976, Leith 1990, , Mason & Thomson 1992, Schumann 1995, Lesieur 1980, Marstop 2007)

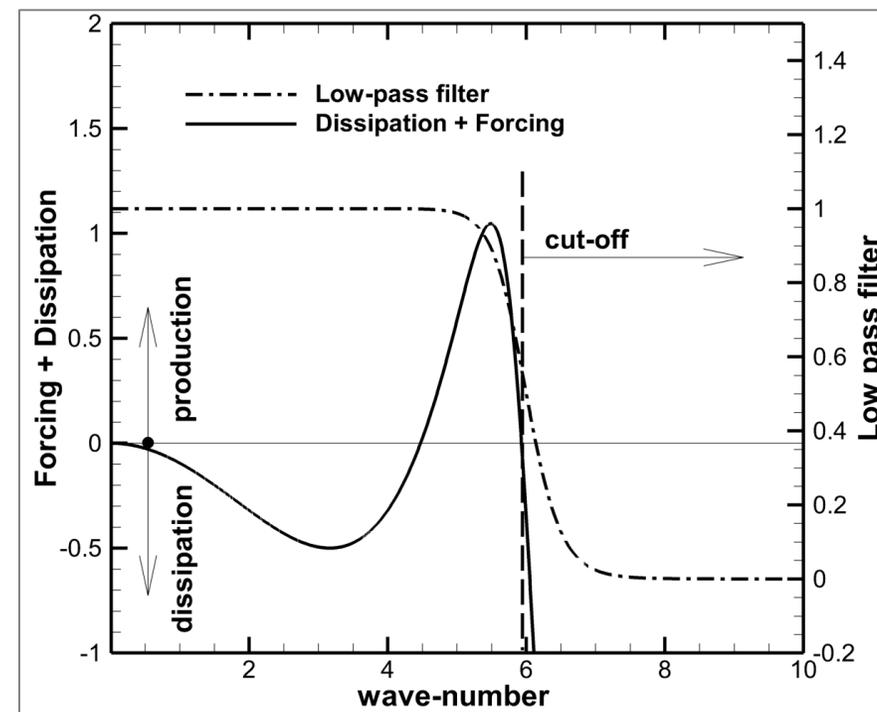
- k^4 -backscatter-spectrum
- divergence-free (solenoidal) forcing
- spatial correlation length scale of forcing from local mesh resolution
- forcing defined in Lagrangian frame
- realization of finite local turbulent time scale

$v(k,t)$ one obtains an equation for the energy $E(k,t)$. In principle, all of the well-known theories are capable of casting the energy equation into the general form

$$\left\{ \frac{\partial}{\partial t} + k^2 [\nu + \nu_e(k,t)] \right\} E(k,t) = k^4 A(k,t) \quad (47)$$

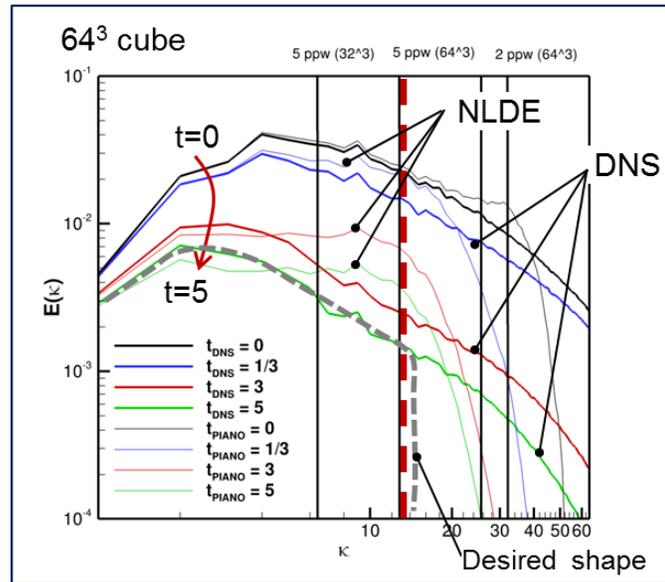
where the nonlinearities are absorbed into an eddy-viscosity, $\nu_e(k,t)$ and a forcing term, $A(k,t)$, where ν and A depend only weakly on k . The ν_e term in equation (47) is readily understood,

Marcus, 'NUMERICAL MODELING OF SUBGRID-SCALE FLOW IN TURBULENCE ROTATION AND CONVECTION', 1986

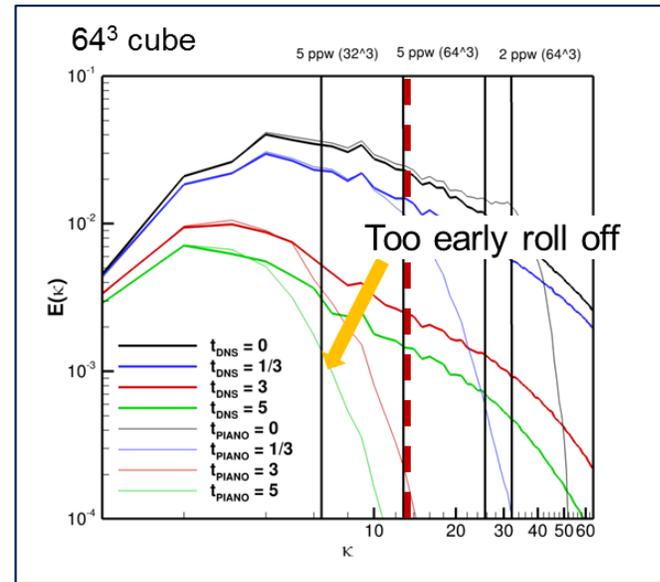


Backscatter Calibration for DHIT (AIAA 2017-3017)

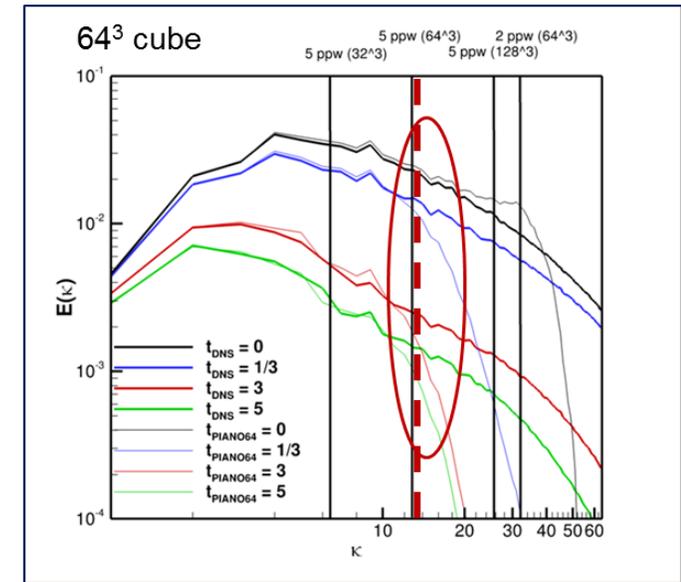
Status Accomplished



No eddy viscosity, no forcing



Calibrated eddy viscosity, no forcing



Calibrated eddy viscosity + forcing

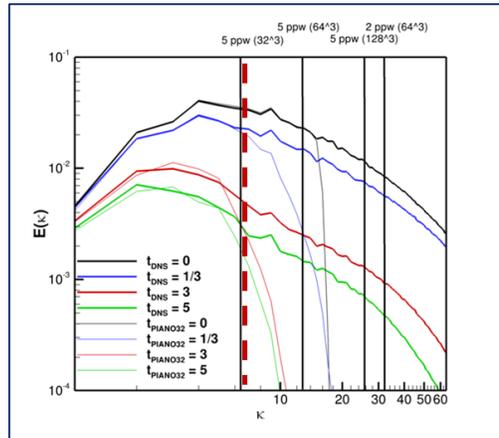
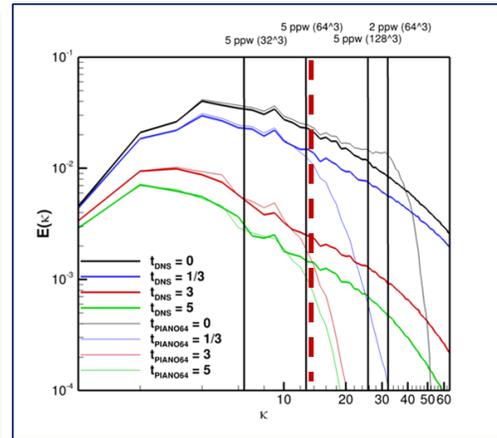
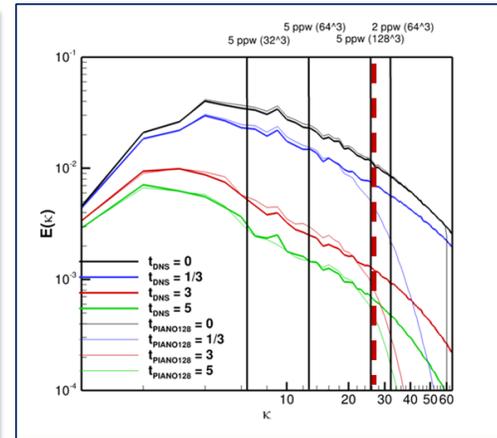
- Results derived from Decaying Homogeneous Isotropic Turbulence (DHIT) simulation (Wray data*)
 - i. k^4 -scaling of backscatter forcing from literature confirmed
 - ii. Forcing scaling derived (as a function of cut-off wave number / length scale)
 - iii. Proper scaling of residual eddy viscosity derived from calibration
 - iv. High order filtering of equations not sufficient, eddy-viscosity needed to avoid energy pile-up in spectrum (subgrid forcing comprises dissipative and productive parts)

*DNS Data available on AGARD database; A. Wray 1997;
<http://torroja.dmt.upm.es/turbdata/agard/chapter3/HOM02/>



Backscatter Calibration for DHIT (AIAA 2017-3017)

- Methods works for different mesh resolutions

32³ cube64³ cube128³ cube

- Eddy viscosity calibration yields a modified damped eddy viscosity model (Speziale '98)
- Scaling function based on the RANS length scale L_R instead of Kolmogorov scale

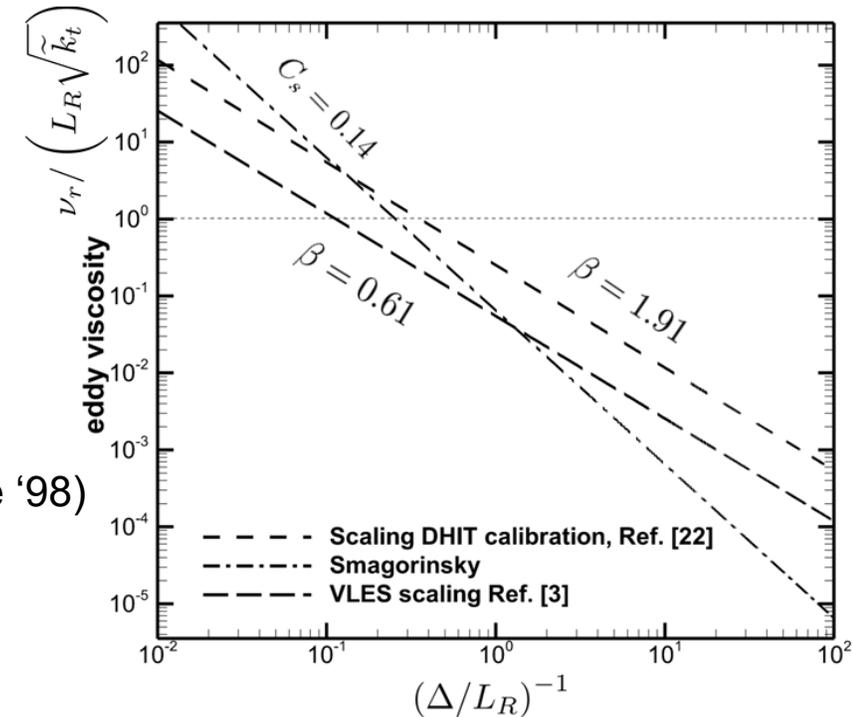
$$\nu_r = F_r \nu_t$$

$$F_r = [1 - \exp(-\beta \Delta / L_R)]^n$$

$$\Delta = \text{cut-off length scale}, L_R = k_t^{3/2} / \epsilon, n = 4/3, \beta \approx 1.91$$

- Corresponds with VLES model of Han et al.¹, albeit with larger beta-value

⁴J. Sci China-Phys Mech Astron, 2012



Cut-off depending eddy-viscosity scaling;
Smagorinsky vs. damped eddy viscosity
model (from calibration)

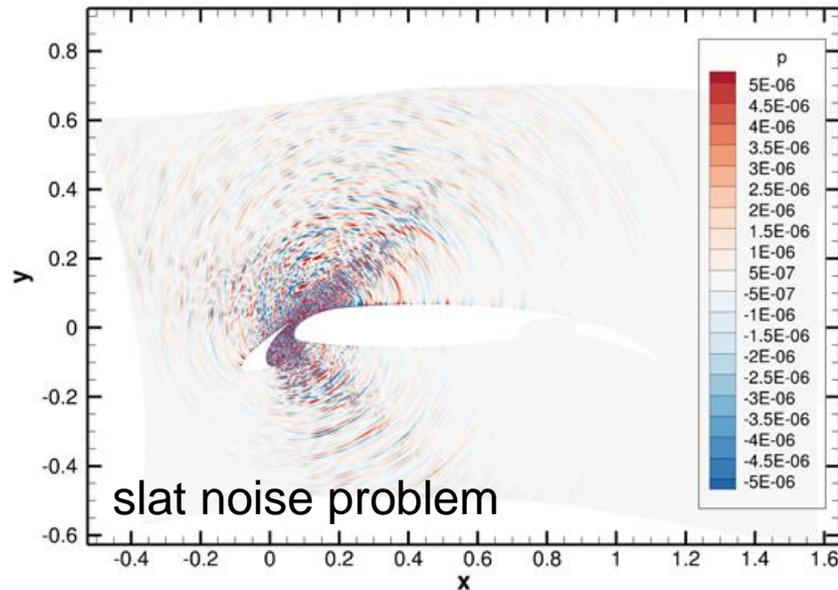
Extension of DHIT calibration to use cases

Modification of dissipative part

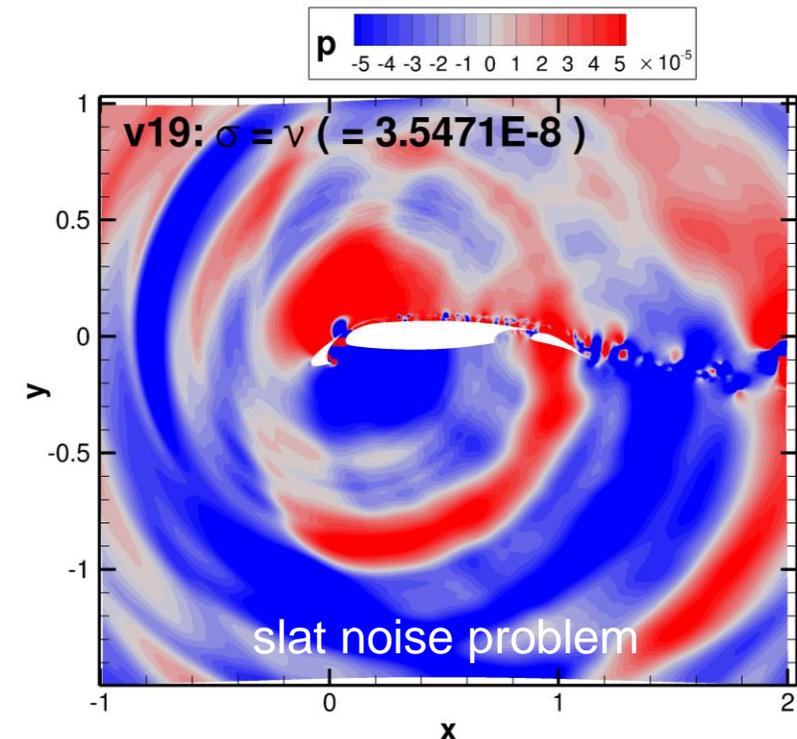
$$\nu_r^* = \sqrt{\frac{\tilde{k}_t}{k_t}} \nu_r$$

- Slat noise with calibrated backscatter model (Heitmann & Ewert 2018)
 - Observation: calibrated damped residual eddy viscosity at start of comp. too high (no resolved fluctuations)
 - Solution: eddy viscosity scaled with resolved TKE from moving-averaging of solution
 - Moving-average time scale from background RANS

$$\tilde{k}_t^{n+1}(\mathbf{x}) = \theta \tilde{k}_t^n(\mathbf{x}) + (1 - \theta) \frac{|(\mathbf{v}')^n|^2}{2}, \quad \theta := \exp\left(-\frac{\Delta t}{\tau_s}\right)$$



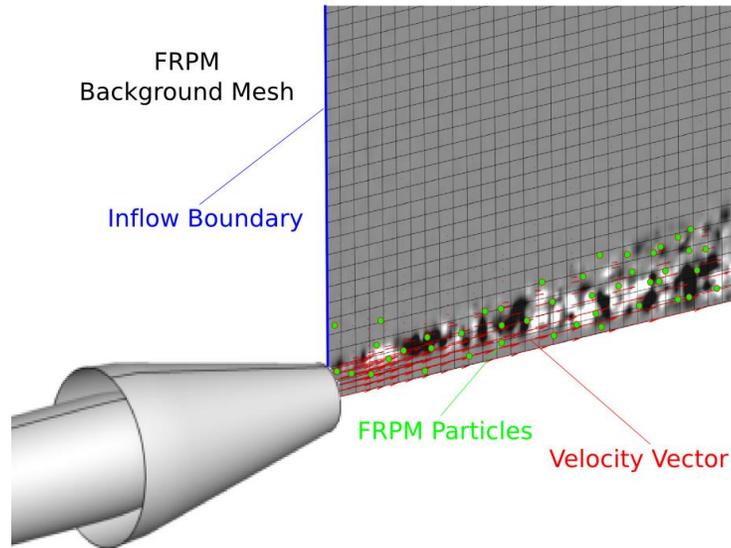
Example high-lift slat-noise: frequency content too high at slat, i.e., low frequency noise missing



Modified dissipative model: proper sound generation @slat

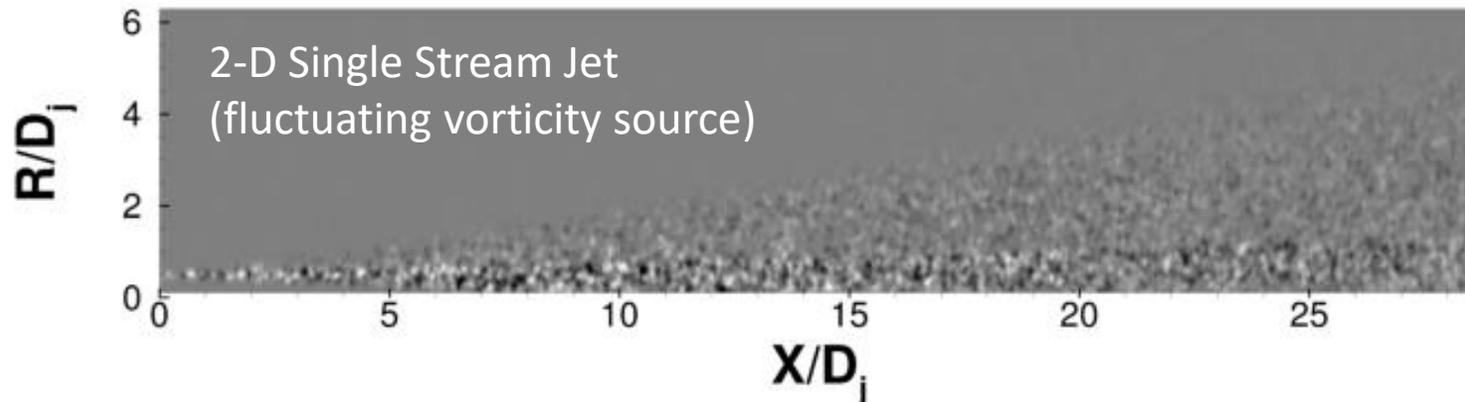
3+1-D Stochastic Backscatter Forcing

FRPM: Fast Random Particle-Mesh Method



- Velocity field divergence-free (solenoidal)
- Convection property of vorticity in non-uniform mean-flow
- Local reconstruction of target turbulent kinetic energy from RANS
- Local reconstruction of turbulent time- & length-scales from RANS
- Broadband spectra realized
- Computationally efficient / parallelized

AIAA Pap. 2005-2862 / AIAA Pap. 2006 / Comp. & Fluids 37 (2008) / AIAA 2007-3506 / AIAA 2009-3369 / AIAA 2009-3175 / JSV 330 (2011) / Comp. & Fluids 132 (2016)



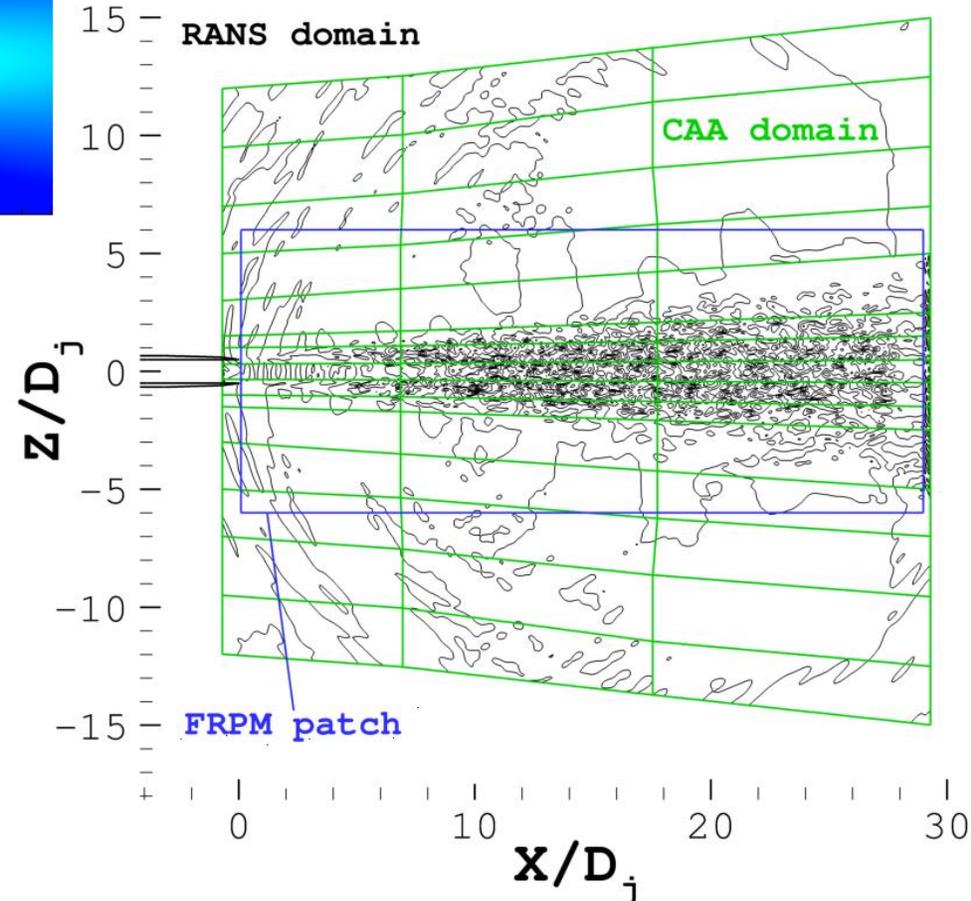
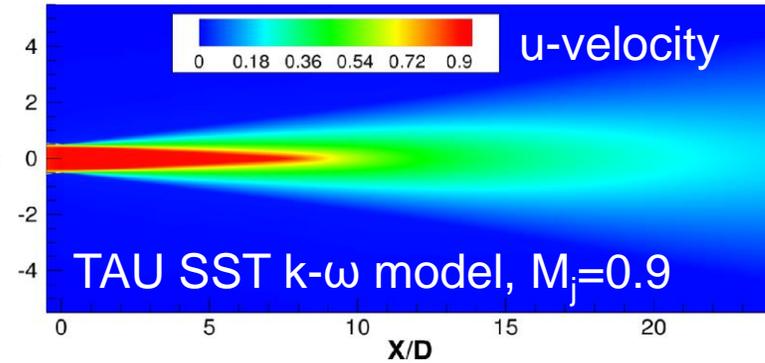
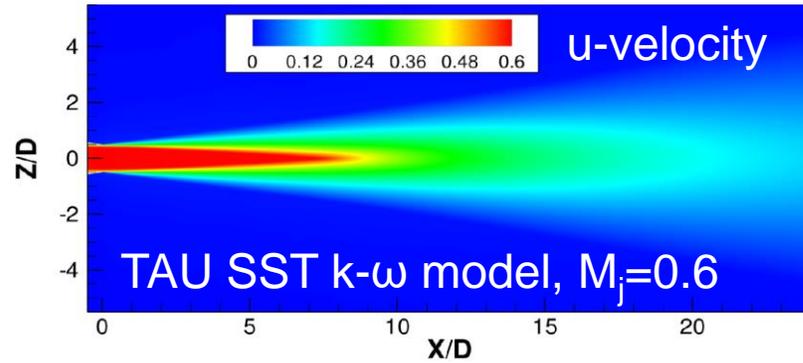
Outline

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- **Results for cold isolated single stream jets**
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Results for cold isolated single stream jet

Computational Setup for Single Stream Jet FES



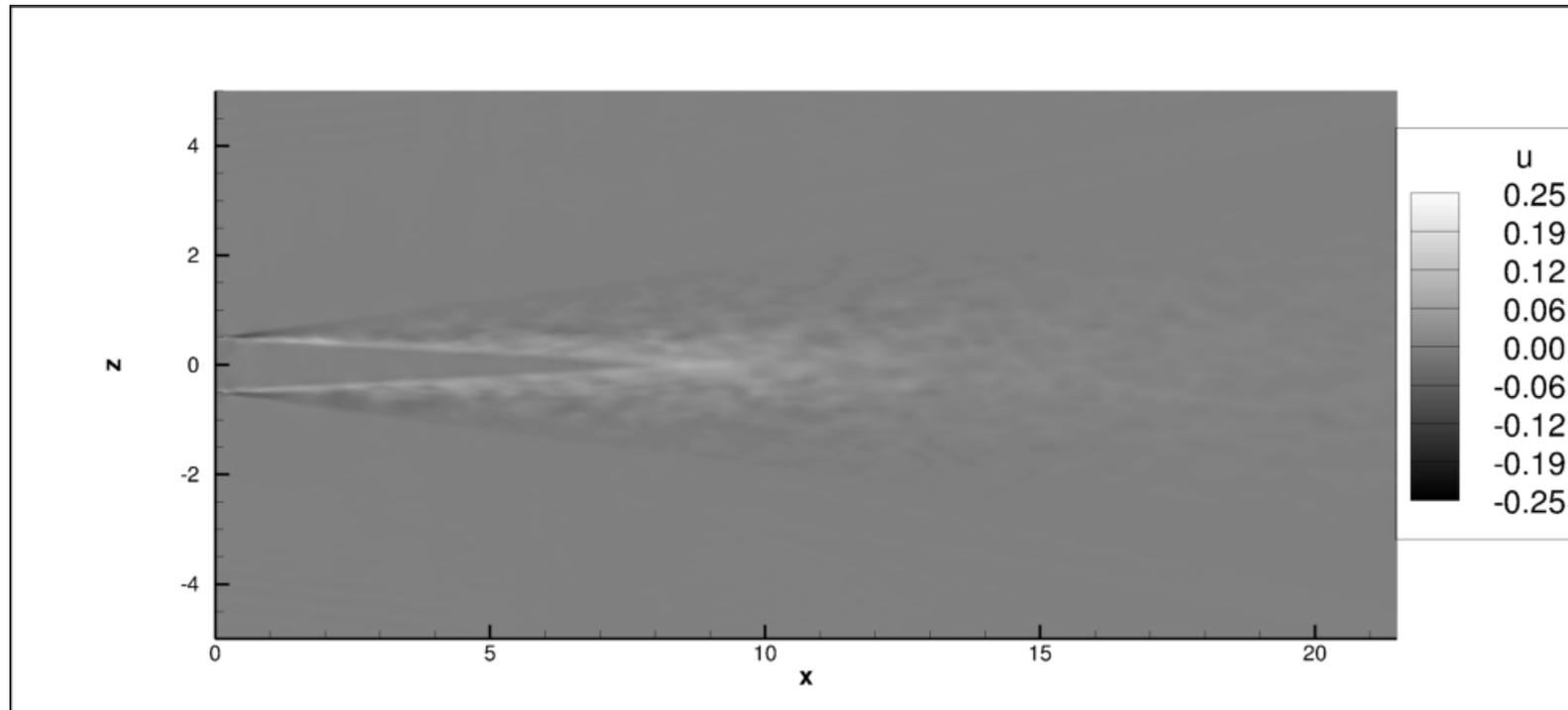
- Two computational cases are used to validate the FES approach
- 3-D cold isolated single stream jet (SSJ) computation for nozzle exit Mach numbers 0.6 ($Re=0.68M$) and 0.9 ($Re=1M$)
- Both SSJ in static condition
- CAA mesh with 87 grid blocks and approx. 9.7 Mio. grid points (O-mesh topology, approx. 20kHz Resolution, $St\sim 3.6$)
- FRPM patch is resolved with approx. 2.6 Mio grid points
- Computations are performed on 88 CPUs with runtime of 4-7 days
- FES domain excludes nozzle interior



Results for cold isolated single stream jets II

Flow Properties

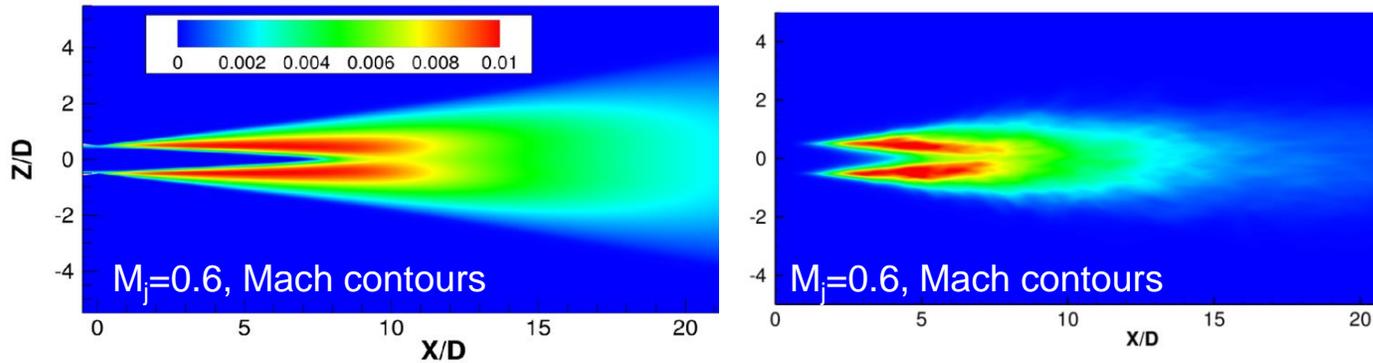
u-velocity $M_j=0.9$ jet



Results for cold isolated single stream jets III

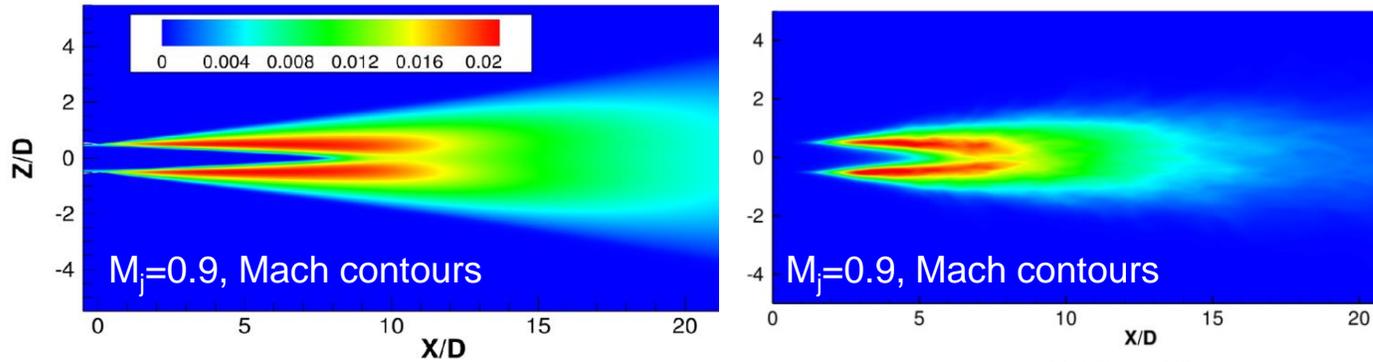
Flow Properties

- Dissipation extended by mean-part: too fast decay observed



(c) $Ma_j = 0.6$, turbulent kinetic energy.

(a) $Ma_j = 0.6$

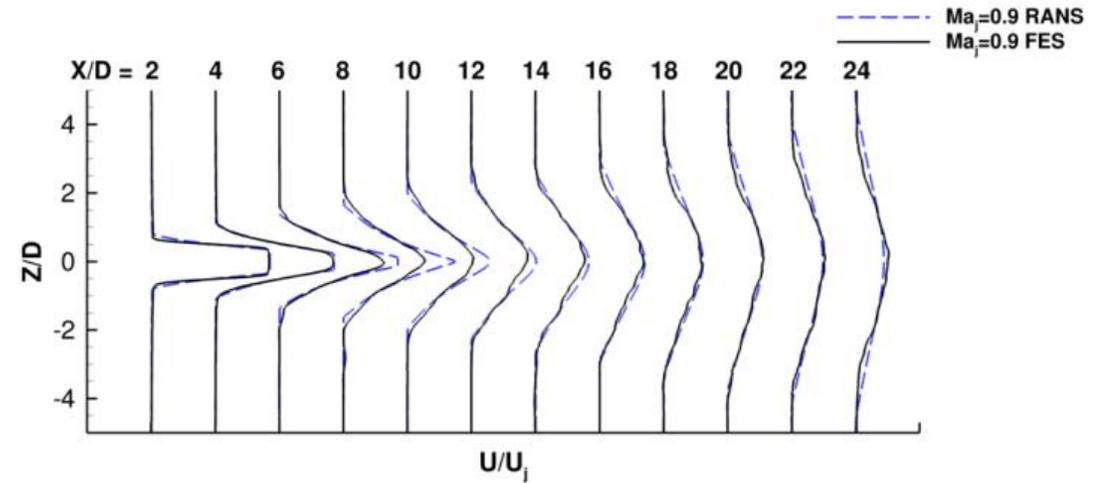


(d) $Ma_j = 0.9$, turbulent kinetic energy.

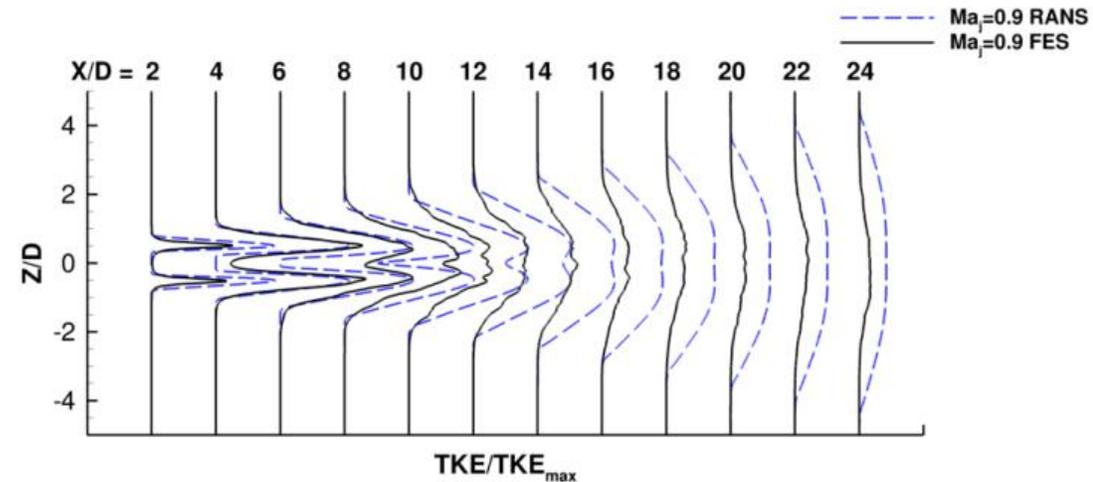
(b) $Ma_j = 0.9$

RANS TKE

FES Resolved TKE



(a) Radial profiles of axial velocity.



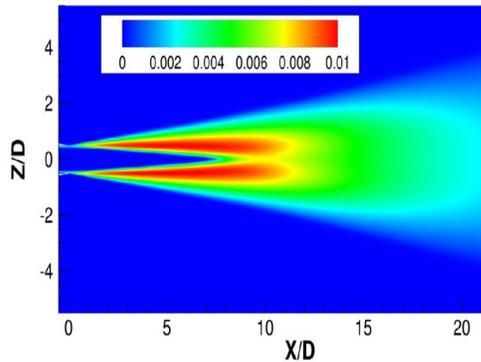
(b) Radial profiles of TKE.



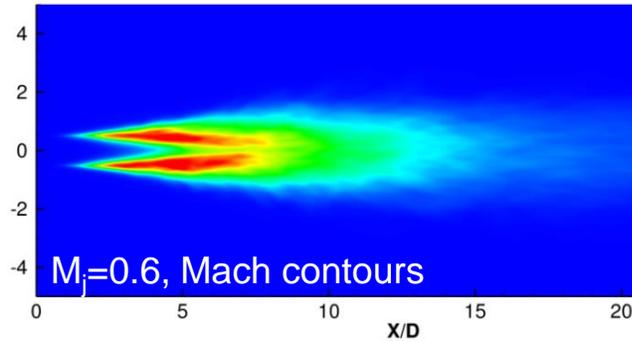
Results for cold isolated single stream jets III

Flow Properties

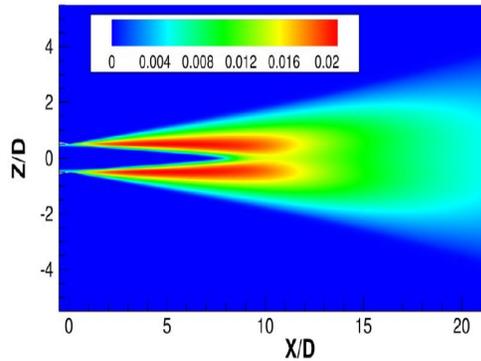
- Dissipation extended by mean-part: too fast decay observed



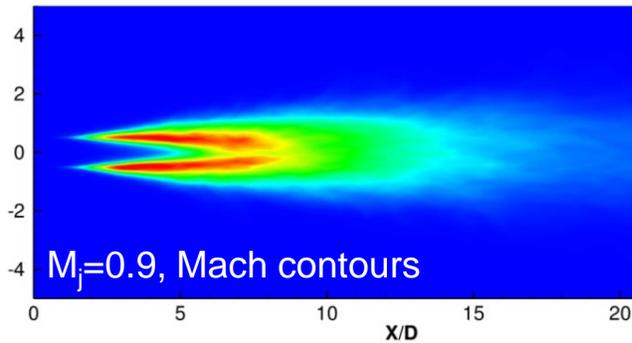
(c) $Ma_j = 0.6$, turbulent kinetic energy.



(a) $Ma_j = 0.6$



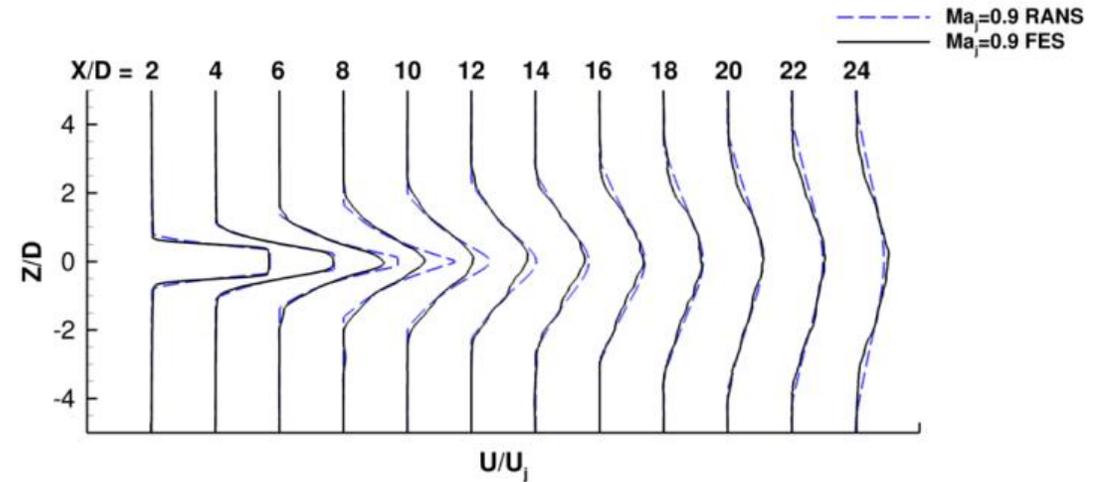
(d) $Ma_j = 0.9$, turbulent kinetic energy.



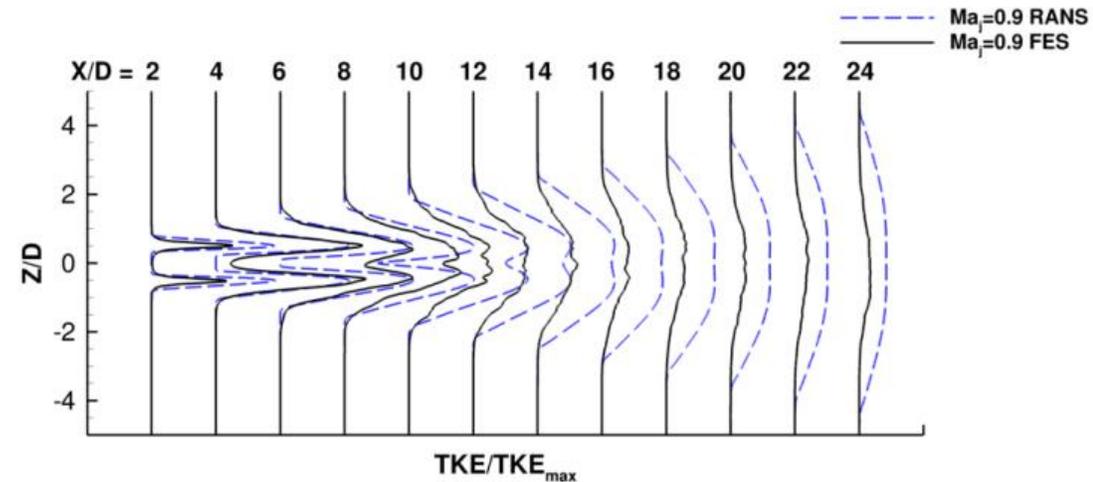
(b) $Ma_j = 0.9$

RANS TKE, 33.3% squeezed

FES Resolved TKE



(a) Radial profiles of axial velocity.



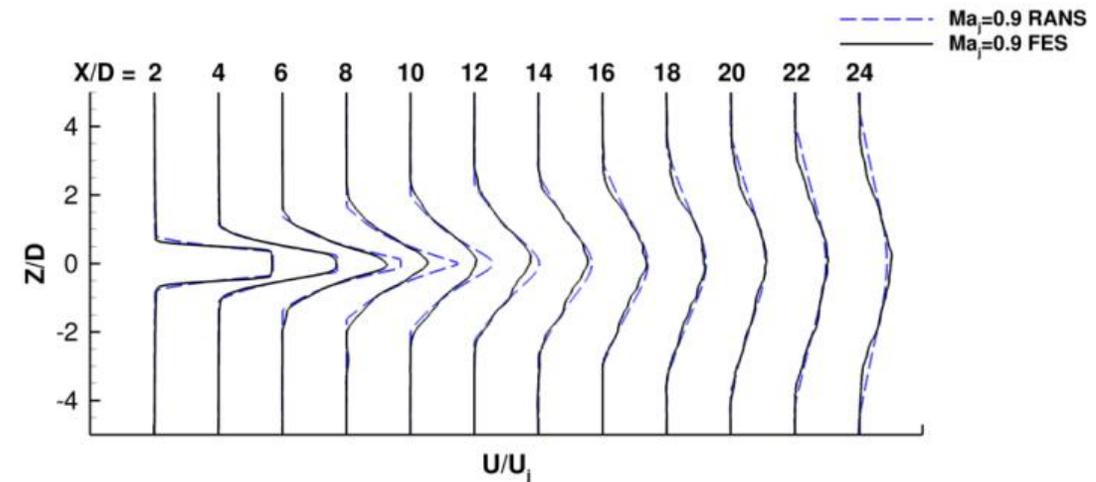
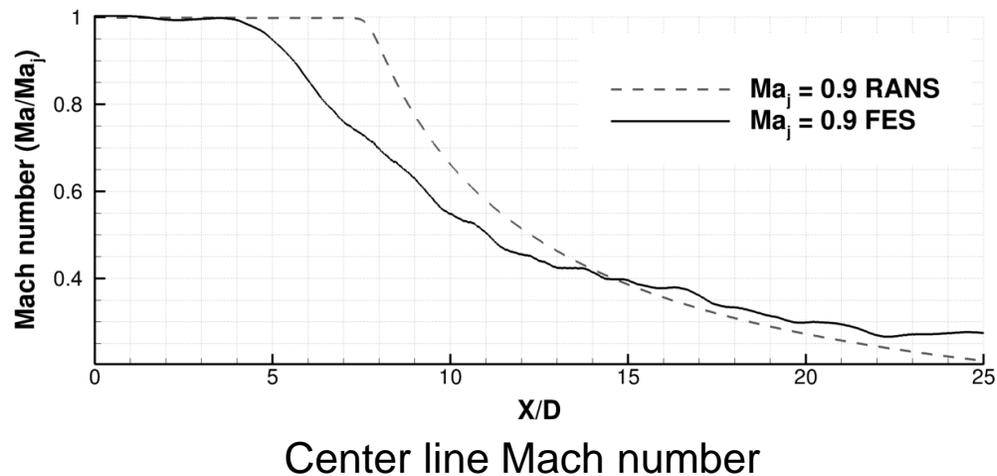
(b) Radial profiles of TKE.



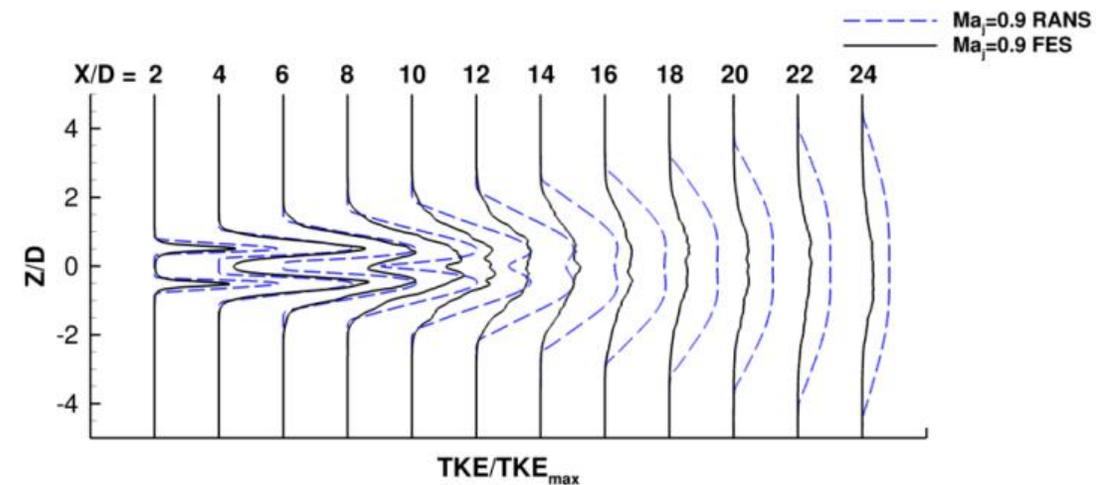
Results for cold isolated single stream jets IV

Flow Properties

- Potential core of RANS too long (7.5D) and drop-off too steep
- RANS mean-flow corrected by NLDE-FES
- Qualitatively right correction of center line Mach number
- FES Potential core (4.5D) slightly too short (4-6D expected)
 - may indicate slightly too strong shear layer mixing
 - stochastic forcing slightly too strong
 - too much filtering



(a) Radial profiles of axial velocity.



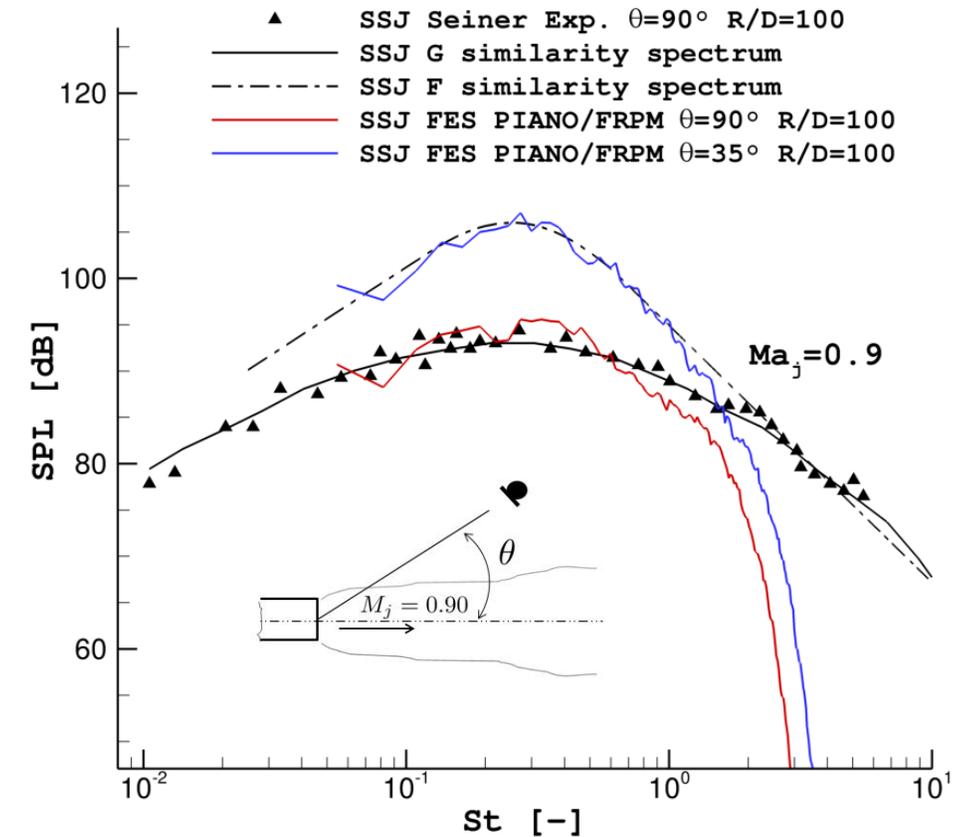
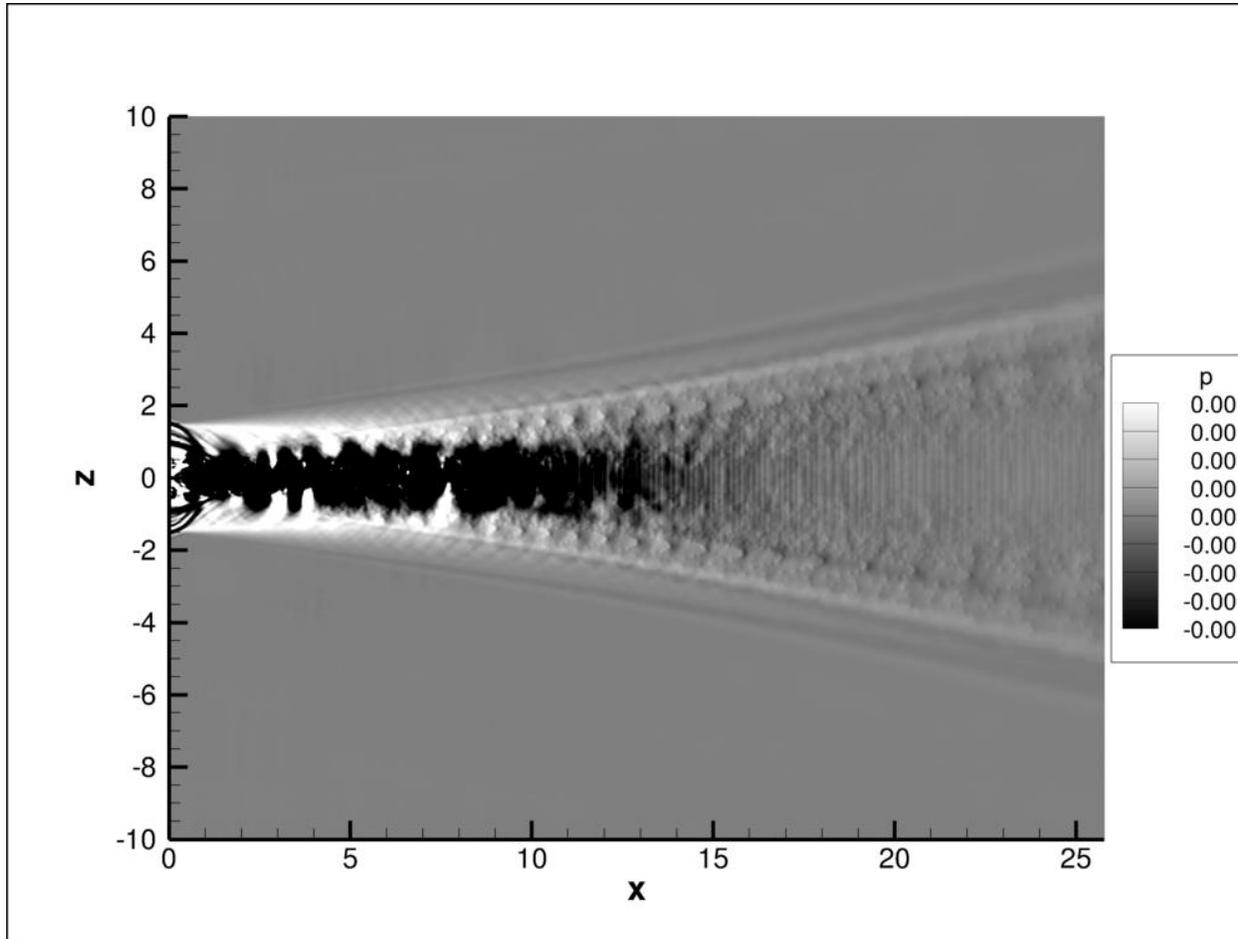
(b) Radial profiles of TKE.



Results for cold isolated single stream jets V

Resolved Sound Field $M_j=0.9$

p' , fluctuating pressure $M_j=0.9$ jet



Comparison of spectra with reference

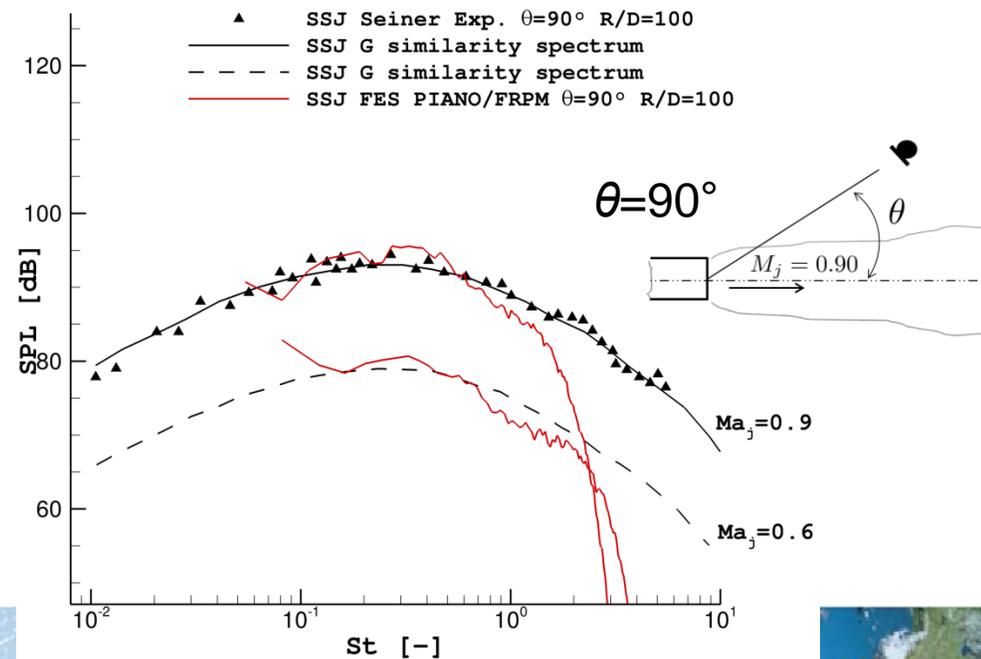
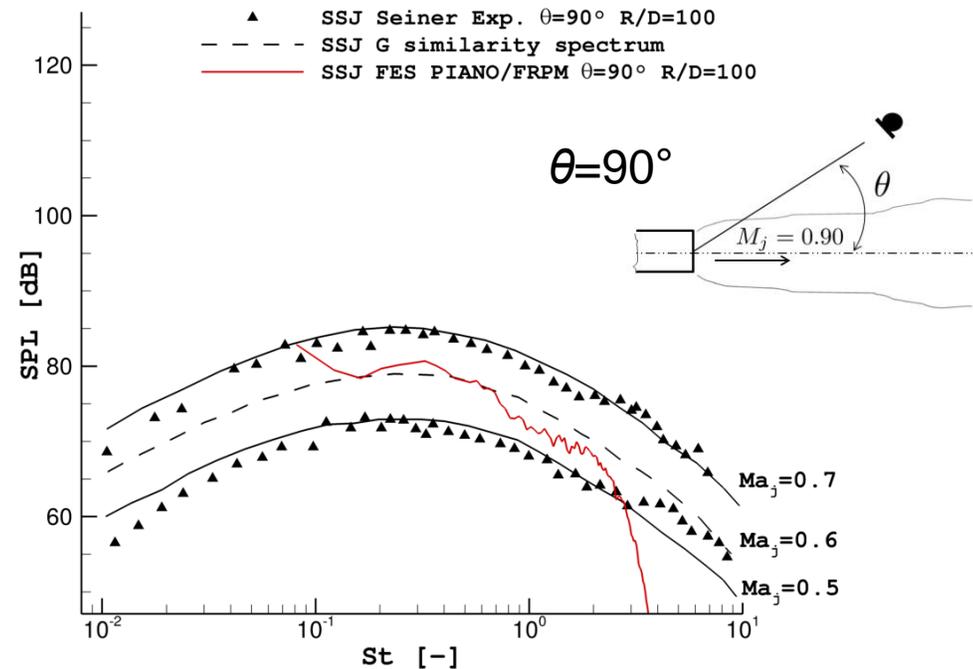
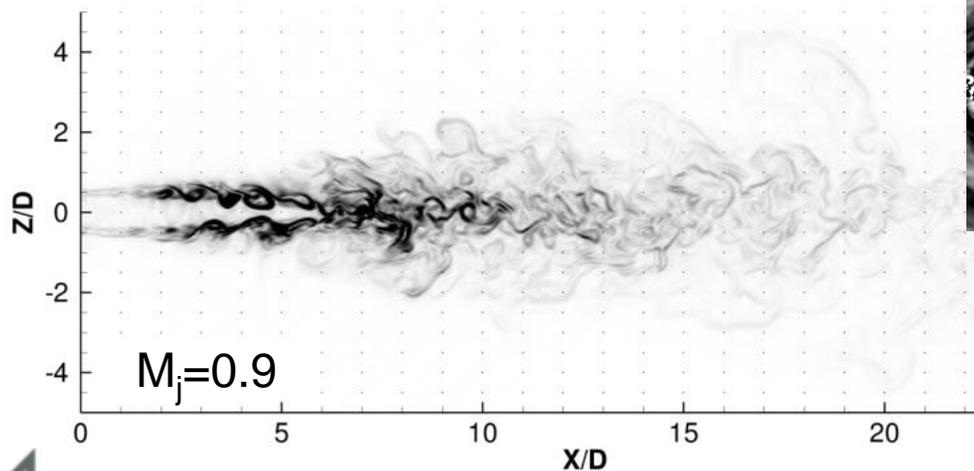
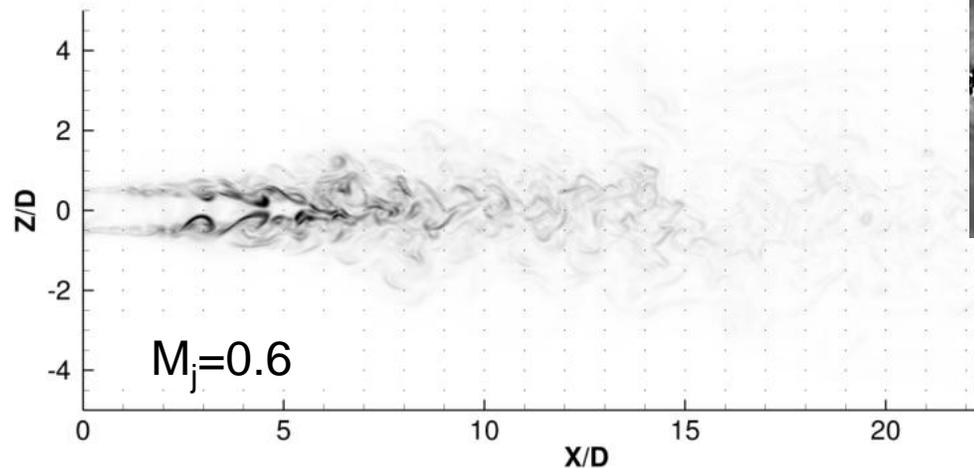
- No enhanced laminar-to-turbulent shear-layer transition and increased noise due to vortex pairing



Results for cold isolated single stream jets VI

Flow Properties

- Right Mach-number scaling



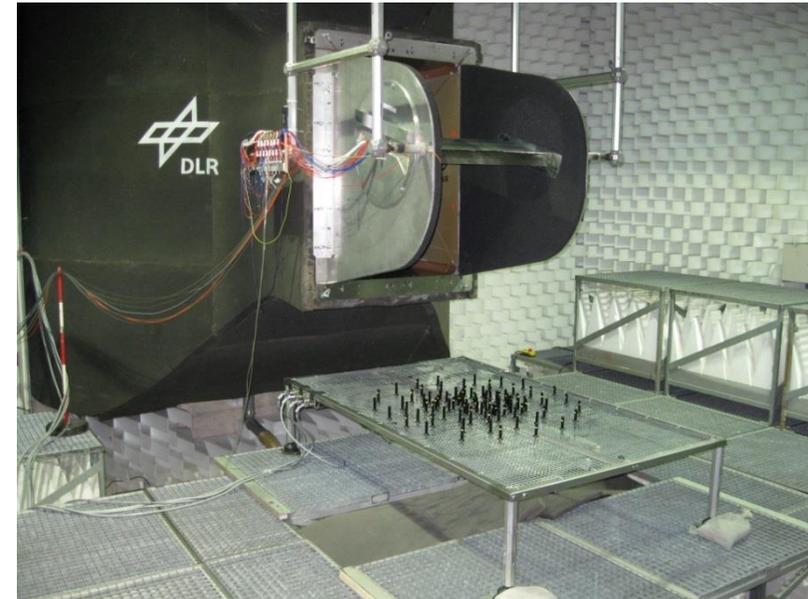
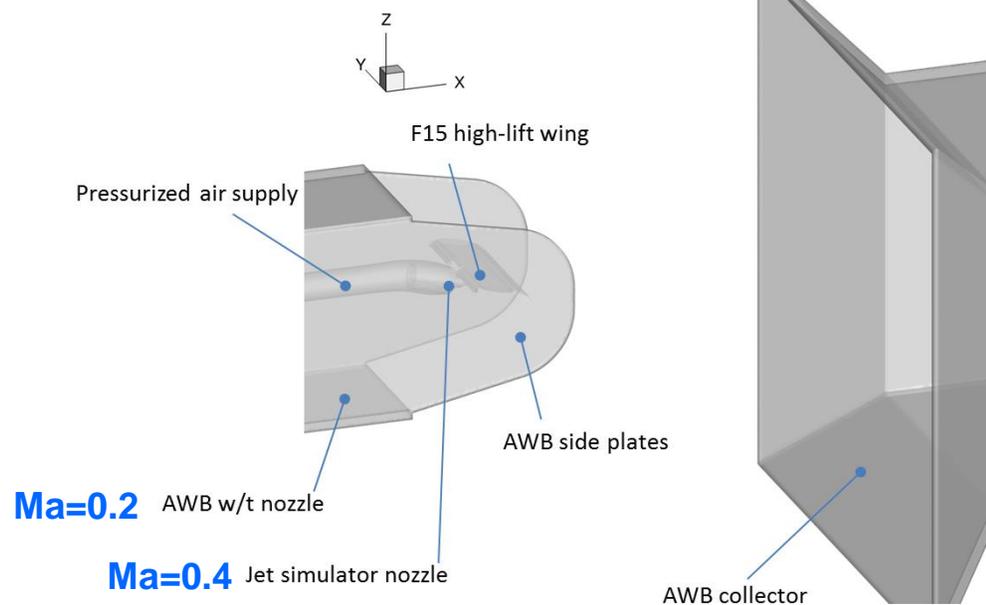
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Results of installed UHBR configuration

Installed UHBR configuration in AWB wind tunnel I



Experimental setup of LIST configuration

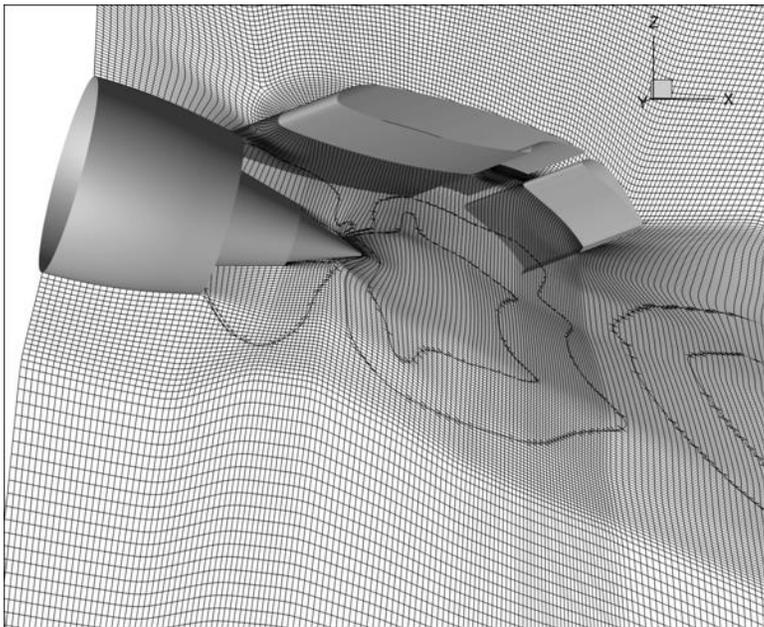
- Numerical setup is reproducing the experimental setup of AWB for installed UHBR configuration
- Computational domain is limited to area of interest, i.e. between the AWB nozzle and collector; in spanwise direction between the side plates
- Flow conditions are emulating the approach conditions with flap angle 25°



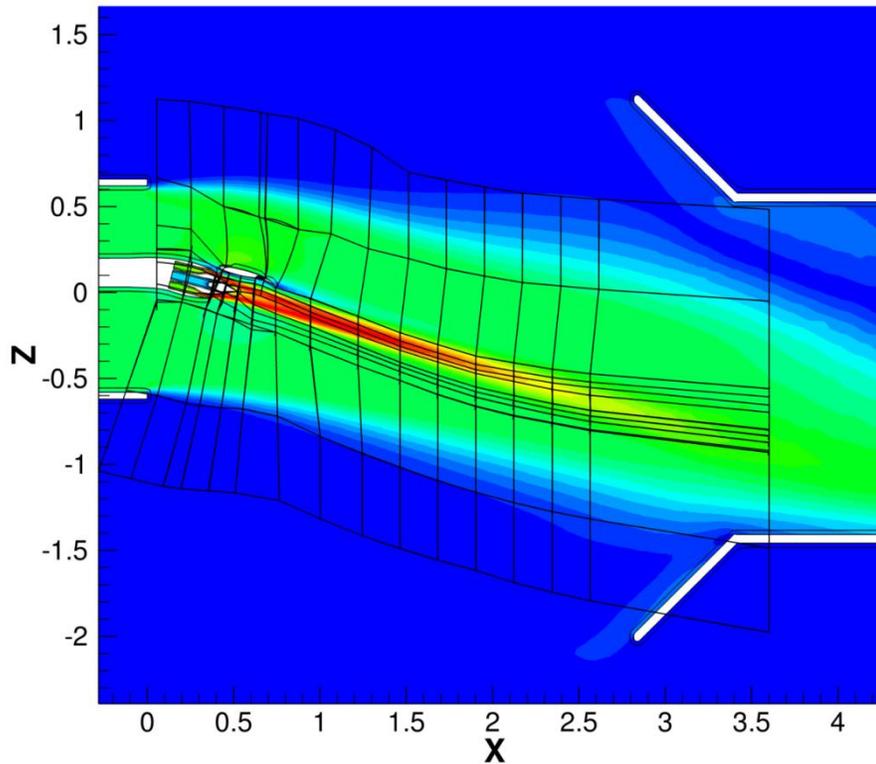
Results of installed UHBR configuration

Computational Chain

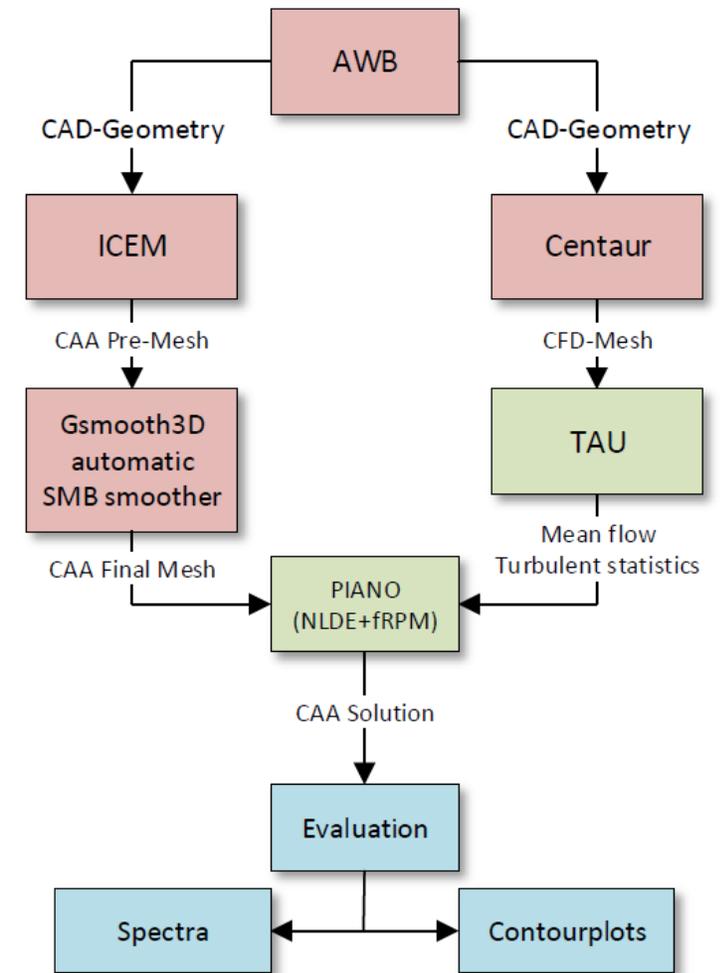
- Smoothed SMB mesh, AWB jet-wing installation, 350 blocks, 30mio points



CAA mesh



TAU RANS mean-flow and CAA mesh

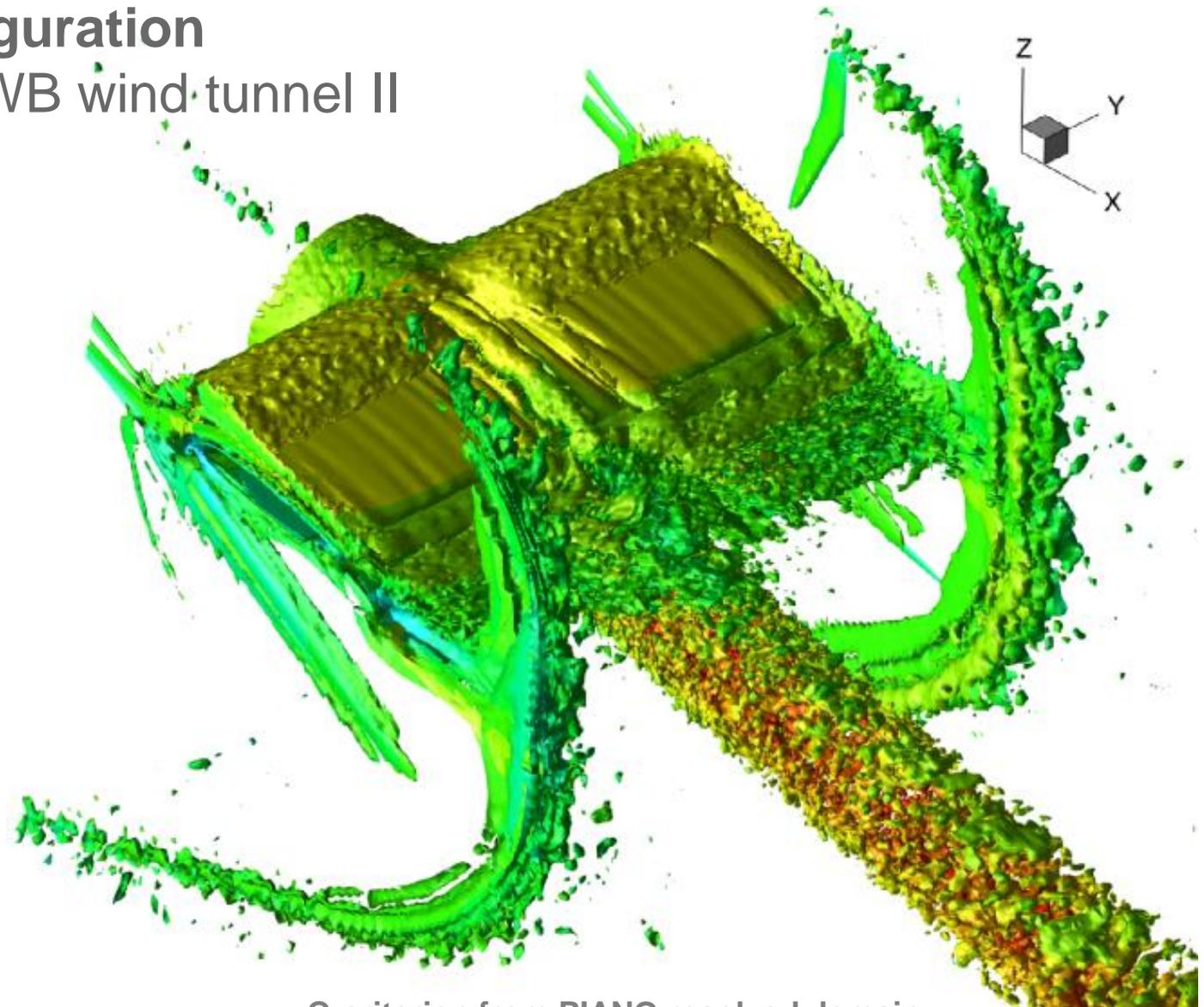
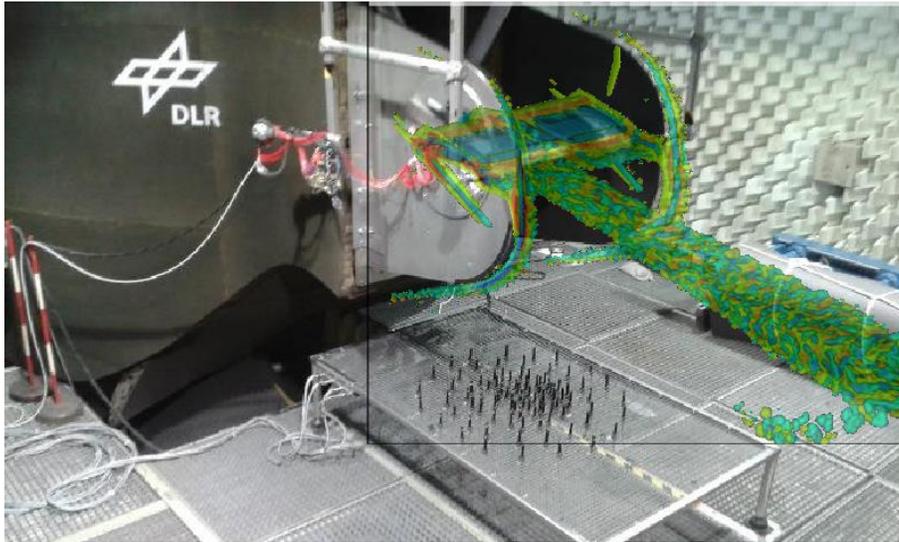


- Preprocessing
- Simulation
- Postprocessing



Results of installed UHBR configuration

Installed UHBR configuration in AWB wind tunnel II

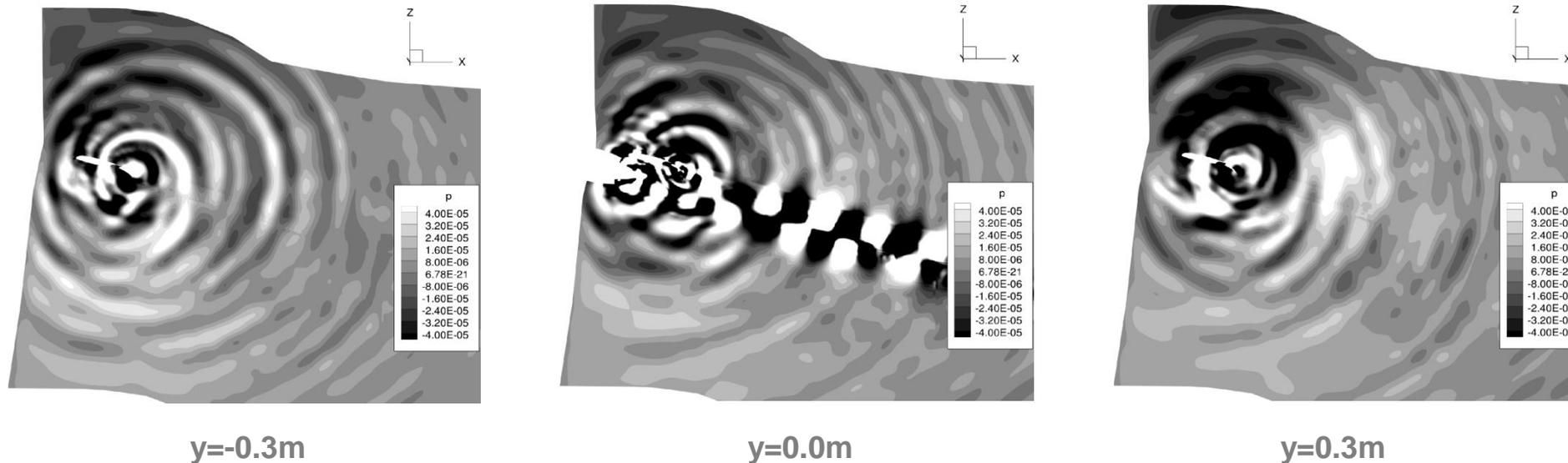


Q-criterion from PIANO resolved domain

- Q-criterion (iso-surfaces) colored with the flow velocity
- Spanwise extend of computational domain does not include w/t side plates but are closely positioned to it
- Appearance of wake structures from the trailing edges of side plates
- Vortex shedding on the upper side of the wing is triggered by close position of the UBHR nacelle0

Results of installed UHBR configuration

Installed UHBR configuration in AWB wind tunnel III



- Contour plots of pressure fluctuations at three different spanwise cuts; in the symmetry plane of UHBR nozzle and to the left/right with 0.3m separation
- Strong circular noise radiation is observed with its origin at the flap trailing edge
- The hydrodynamic pressure fluctuations of the UHBR jet are observed in the cut at the centerline
- Noise radiation to the bottom side appears to be less intensive compared to region above wing



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Summary & Conclusions

- **Method**

- Application of stochastic backscatter method to technical problem: Forced Eddy Simulation
- Successful transfer of DHIT calibration to single stream jet test case

- **Single stream jet**

- NLDE provide correction to RANS mean flow (potential core length)
- TKE level properly reproduced in magnitude by resolved fluctuations with respect to RANS
- Slight modification of the TKE topology
- Proper prediction of absolute level, spectra, directivity and Mach number scaling

- **Installed jet noise**

- Tool chain established
- RANS simulation of entire wind tunnel set up
- FES simulations running
- Jet-flap interactions identified in results



Thank You!



Knowledge for Tomorrow

