



Deutsches Zentrum
für Luft- und Raumfahrt
German Aerospace Center

Institute of Aerodynamics
and Flow Technology

Towards Forced Eddy Simulation for Installed Nozzle-Wing Configurations

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CEAA 2018 Aeroacoustics Conference
18.-23. September 2018
Svetlogorsk, Russia



Knowledge for Tomorrow

Aim of the Work

Installed Jet Noise Simulation

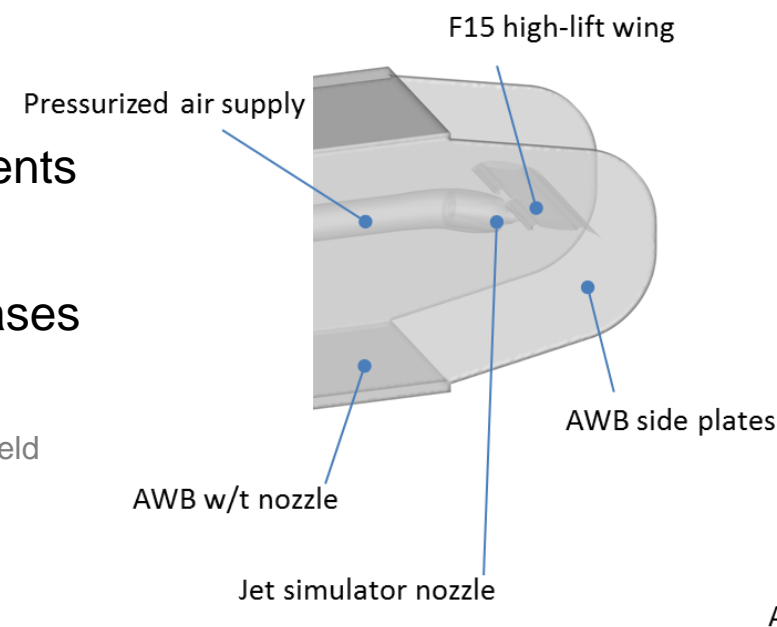
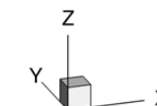
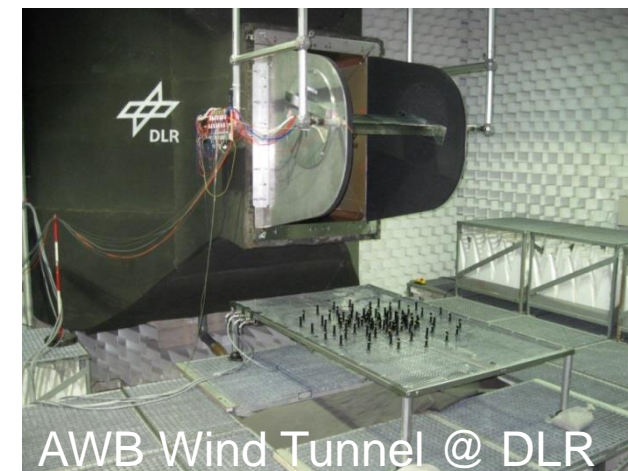
- Direct Noise Computation (DNC¹) of a jet-flap installation noise wind tunnel setup tested in the Acoustic Wind Tunnel Braunschweig (AWB) at DLR
- DNC¹ using a zonal RANS/VLES² approach based on Non-Linear Disturbance Equations (NLDE)³
 - Zonal VLES of the entire wind tunnel setup using a standard CFD/CAA framework
 - Evaluation of the potential to lower resolution requirements using an active stochastic backscatter model
 - Transfer of stochastic backscatter from DHIT⁴ to use cases

¹Bailly, Bogey & Marsden '10: solving the compressible Navier-Stokes equations to determine simultaneously the aerodynamic field and the acoustic field

²Very-large-eddy simulation: filter and grid are too coarse to resolve 80% of the energy (Pope 2000)

³Morris '97, Sagaut & Labourasse '02, Terracol et al. '06, Batten et al. '04

⁴Decaying Homogeneous Isotropic Turbulence



Outline

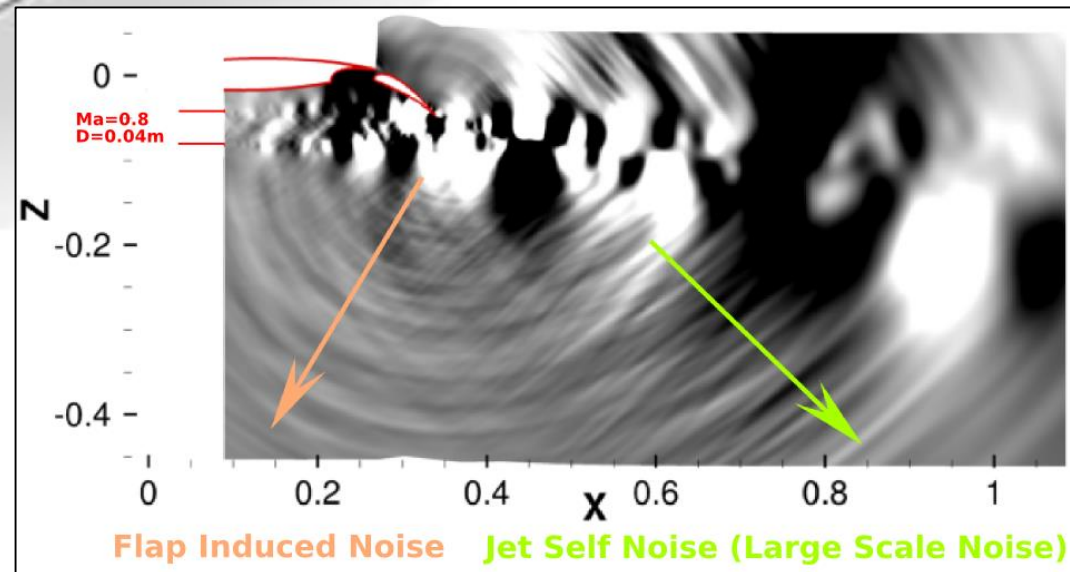
- **Simulation Problem**
- Methodology
- Results of cold isolated single stream jet
- Results of installed UHBR configuration
- Summary & Conclusions



Simulation Problem

UHBR Installation Noise

- Noise prediction for configurations with nozzle mounted closely to wing
- The installation noise from jet-airframe interaction potentially is one of the prominent noise source
- This noise generation mechanism is attributed to coherent hydrodynamic fluctuations passing by the airfoil trailing edge
- Qualitatively different noise generation mechanisms as can be observed in directivity pattern:
 - Airframe → maximum to forward arc
 - Jet Noise → maximum to rearward arc



Outline

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Methodology: Zonal RANS/VLES Approach

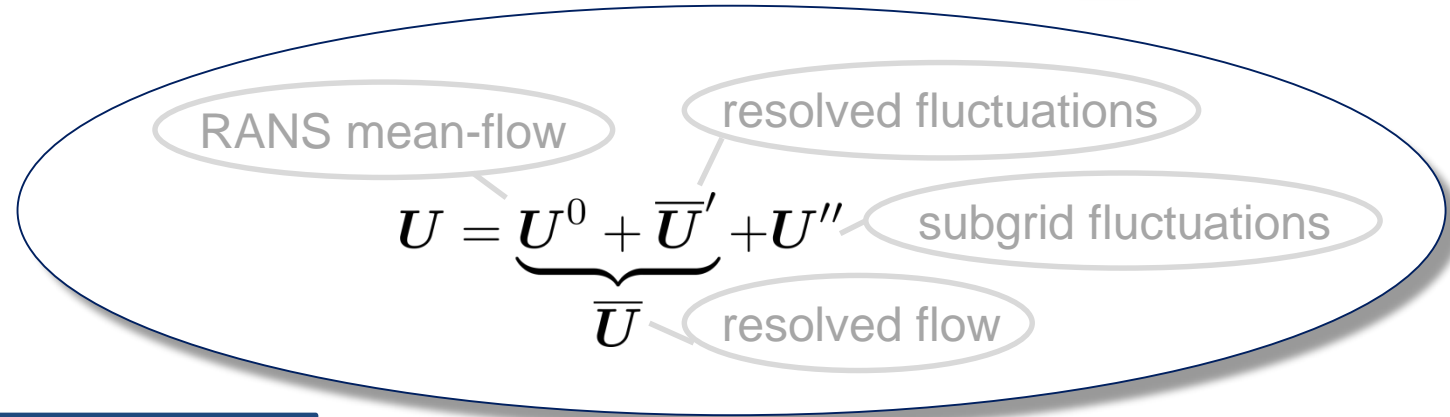
NLDE⁴ (Non-Linear Disturbance Equations)

⁴Morris '97, Sagaut & Labourasse '02, Terracol et al. '06, Batten et al. '04

- Navier-Stokes equations for primitive variables:

$$\frac{\partial \mathbf{U}}{\partial t} + \mathcal{N}(\mathbf{U}) = \mathbf{0} \quad \mathbf{U} = \begin{pmatrix} \rho \\ v_i \\ p \end{pmatrix}$$

- Triple decomposition of variables:



- NLDE Variant 1:

$$\frac{\partial \bar{U}'}{\partial t} + \mathcal{N}(\bar{U}) = f_{sgs}(U'')$$

- NLDE Variant 2:

$$\frac{\partial \bar{U}'}{\partial t} + \mathcal{N}'(U^0, \bar{U}') = f_{visc}^0 + f_{sgs}$$

$$\mathcal{N}(\bar{U}) = \mathcal{N}'(U^0, \bar{U}') + \mathcal{N}^0(U^0)$$

$$\mathcal{N}^0(U^0) = -f_{visc}^0(U^0)$$

≡ RANS



Methodology

NLDE (Non-Linear Disturbance Equations)

- NLDE Variant 1:

$$\frac{\partial \bar{U}'}{\partial t} + \mathcal{N}(\bar{U}) = f_{sgs}(U'')$$

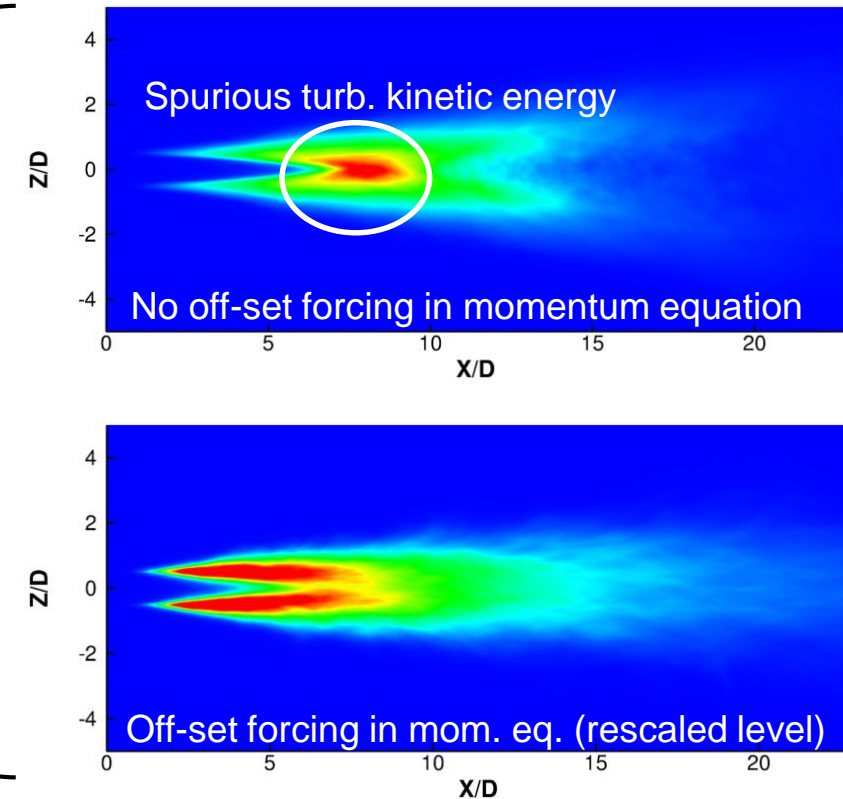
off-set forcing w/o forcing

- NLDE Variant 2:

$$\frac{\partial \bar{U}'}{\partial t} + \mathcal{N}'(U^0, \bar{U}') = f_{visc}^0 + f_{sgs}$$

explicit off-set forcing

- Constant off-set necessary to avoid spurious turbulence in single stream jet
- “Variant 1” used for momentum equation (“off-set forcing w/o forcing”)
- “Variant 2” used for density + pressure equations (no off-set forcing)
- Background RANS flow from DLR CFD solver TAU
- NLDE realized with DLR code PIANO
 - 4th order DRP scheme of Tam & Webb
 - High-order (HO) filter to remove spurious waves
 - Optimized 4th order Runge-Kutta time integration



NLDE + Stochastic Backscatter

= FES (Forced Eddy Simulation)

$$\mathbf{f}_{sgs} = - \underbrace{\nabla \times (\bar{\rho} \nu_r \boldsymbol{\omega}')}_{\mathbf{f}_D^{(5)}} + \underbrace{\nabla \times \mathbf{q}}_{\mathbf{f}_F}$$

r.h.s. vector force dissipation stochastic forcing

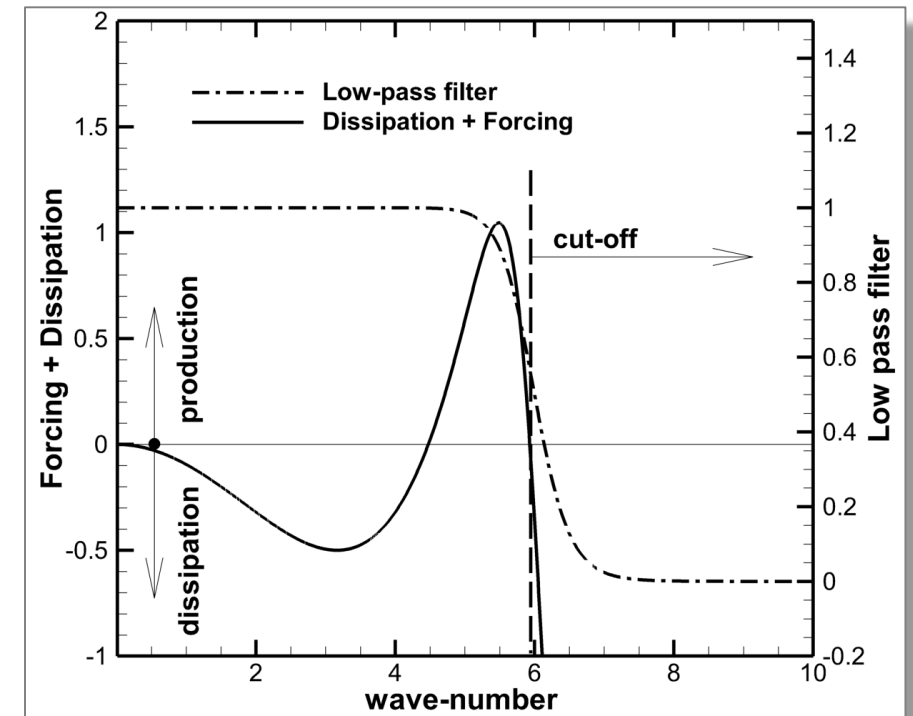
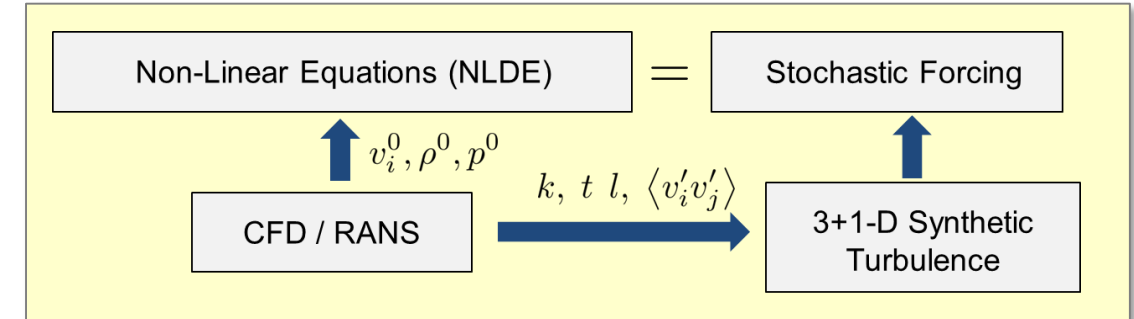
- Smagorinsky-like vector-force sgs-model⁴ of fluctuations
- High-order filtered equations + sgs-model: mixed model⁵
- 3+1-D stochastic forcing provides turbulent backscatter^{1,2,3} from non-resolved scales

⁴Dantinne, Jeanmart, Winckelmans, Legat, Carati, Applied Scientific Research (59), 1998

⁵Bardina, Ferziger, and Reynolds, 1980

¹C.E. Leith, Physics of Fluids A 2, 297, ²U. Schumann, Proc. R. Soc. London A (1995),

³Zamansky et al., J. of Turb. 11, (2010).



NLDE + Stochastic Backscatter

= FES (Forced Eddy Simulation)

$$\mathbf{f}_{sgs} = - \underbrace{\nabla \times (\bar{\rho} \nu_r \boldsymbol{\omega}')}_{\mathbf{f}_D^{(5)}} + \underbrace{\nabla \times \mathbf{q}}_{\mathbf{f}_F}$$

r.h.s. vector force dissipation stochastic forcing

- Backscatter forcing features as discussed in literature

(Kraichnan 1976, Leith 1990, , Mason & Thomson 1992, Schumann 1995, Lesieur 1980, Marstop 2007)

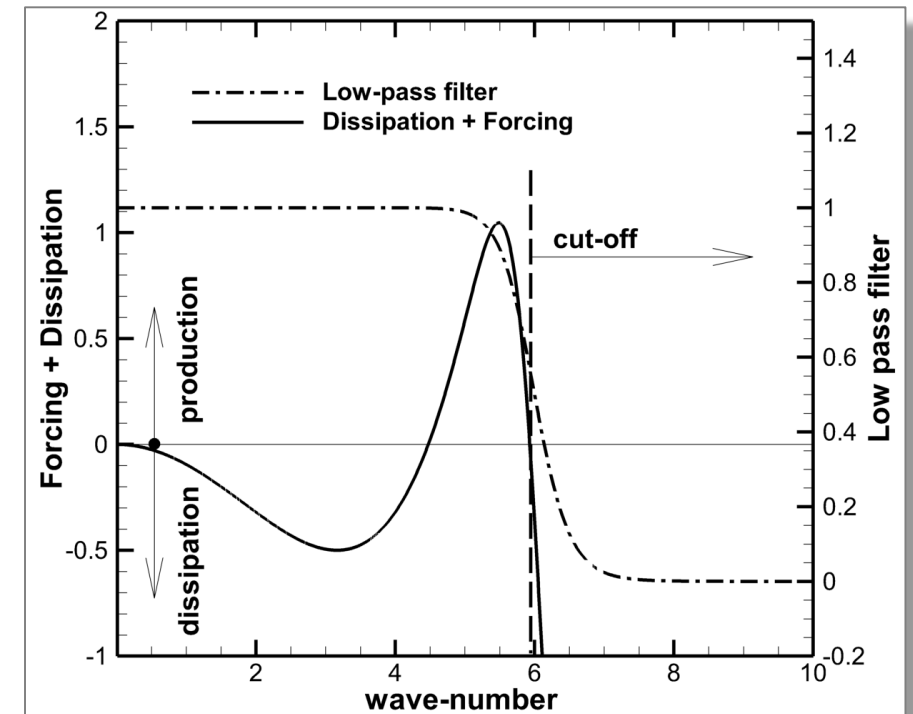
- k^4 -backscatter-spectrum
- divergence-free (solenoidal) forcing
- spatial correlation length scale of forcing from local mesh resolution
- forcing defined in Lagrangian frame
- realization of finite local turbulent time scale

$\mathbf{v}(\mathbf{k},t)$ one obtains an equation for the energy $E(\mathbf{k},t)$. In principle, all of the well-known theories are capable of casting the energy equation into the general form

$$\left\{ \frac{\partial}{\partial t} + k^2 [\nu + \nu_e(\mathbf{k},t)] \right\} E(\mathbf{k},t) = k^4 A(\mathbf{k},t) \quad (47)$$

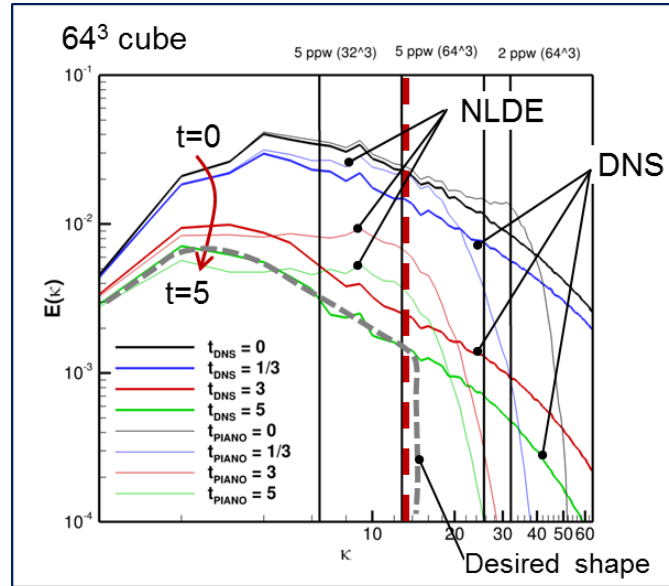
where the nonlinearities are absorbed into an eddy-viscosity, $\nu_e(\mathbf{k},t)$ and a forcing term, $A(\mathbf{k},t)$, where ν and A depend only weakly on \mathbf{k} . The ν_e term in equation (47) is readily understood,

Marcus, 'NUMERICAL MODELING OF SUBGRID-SCALE FLOW IN TURBULENCE ROTATION AND CONVECTION', 1986

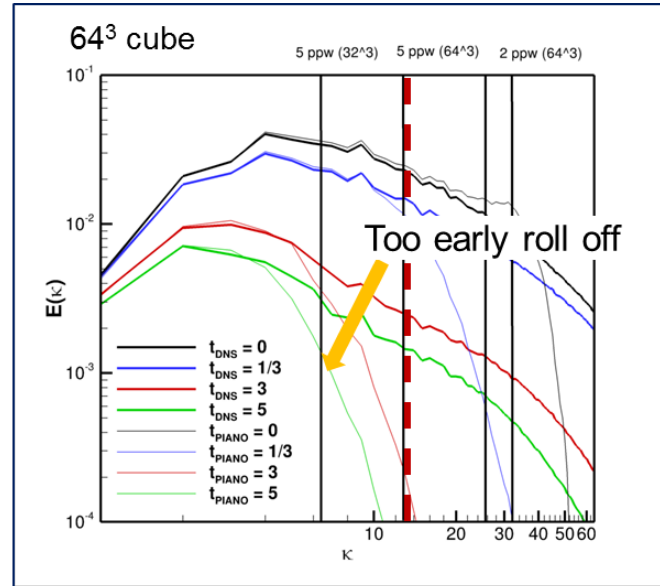


Backscatter Calibration for DHIT (AIAA 2017-3017)

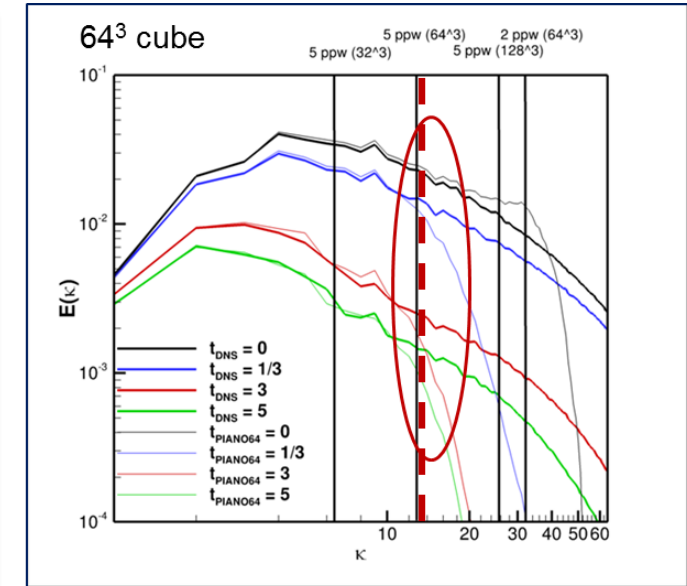
Status Accomplished



No eddy viscosity, no forcing



Calibrated eddy viscosity, no forcing



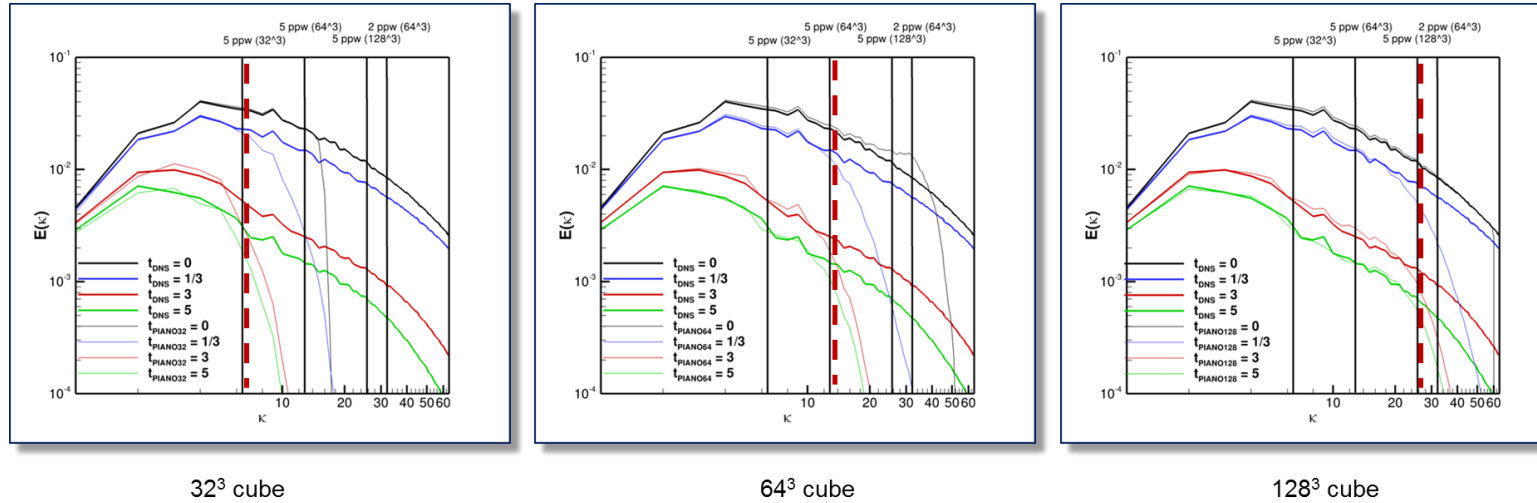
Calibrated eddy viscosity + forcing

- Results derived from Decaying Homogeneous Isotropic Turbulence (DHIT) simulation (Wray data*)
 - i. k^4 -scaling of backscatter forcing from literature confirmed
 - ii. Forcing scaling derived (as a function of cut-off wave number / length scale)
 - iii. Proper scaling of residual eddy viscosity derived from calibration
 - iv. High order filtering of equations not sufficient, eddy-viscosity needed to avoid energy pile-up in spectrum (subgrid forcing comprises dissipative and productive parts)

*DNS Data available on AGARD database; A. Wray 1997;
<http://torroja.dmt.upm.es/turbdata/agard/chapter3/HOM02/>

Backscatter Calibration for DHIT (AIAA 2017-3017)

- Methods works for different mesh resolutions



- Eddy viscosity calibration yields a modified damped eddy viscosity model (Speziale '98)
- Scaling function based on the RANS length scale L_R instead of Kolmogorov scale

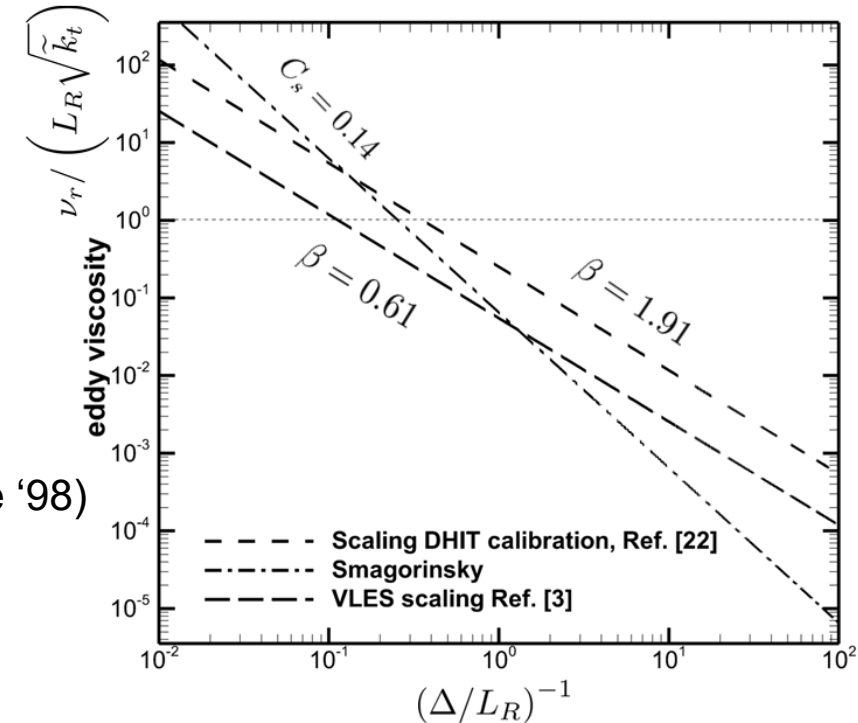
$$\nu_r = F_r \nu_t$$

$$F_r = [1 - \exp(-\beta \Delta / L_R)]^n$$

$$\Delta = \text{cut-off length scale}, L_R = k_t^{3/2} / \epsilon, n = 4/3, \beta \approx 1.91$$

- Corresponds with VLES model of Han et al.¹, albeit with larger beta-value

⁴J. Sci China-Phys Mech Astron, 2012



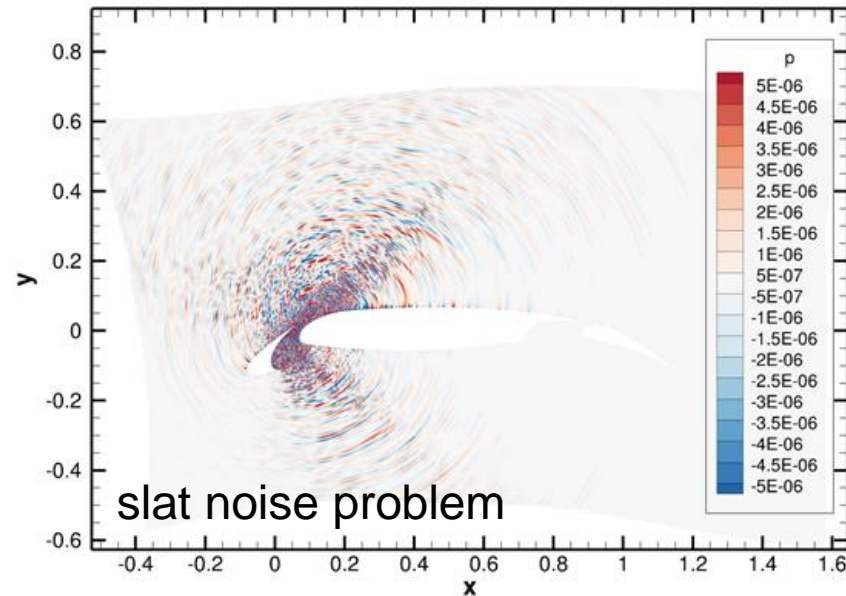
Cut-off depending eddy-viscosity scaling; Smagorinsky vs. damped eddy viscosity model (from calibration)

Extension of DHIT calibration to use cases

Modification of dissipative part

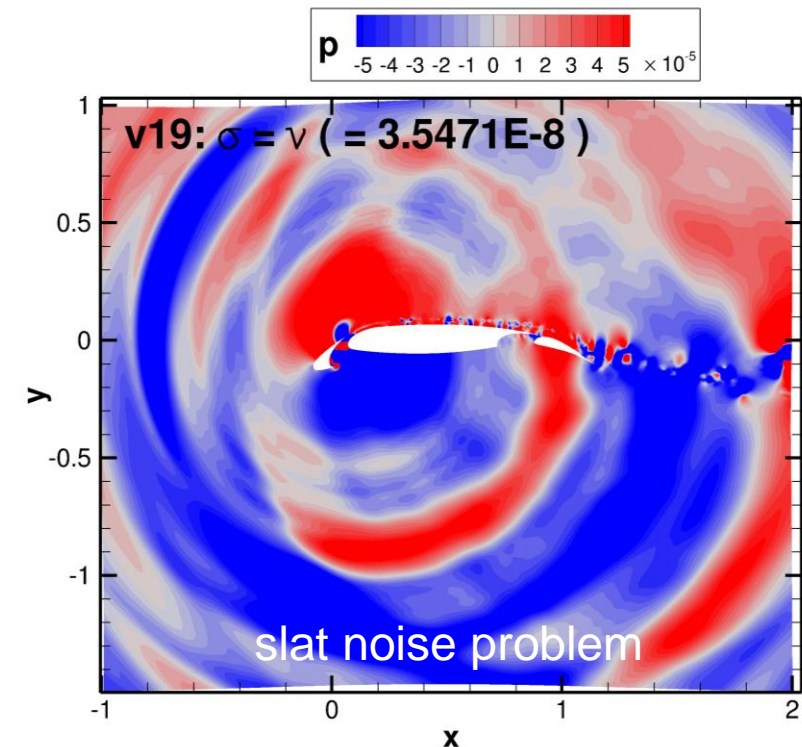
- Slat noise with calibrated backscatter model (Heitmann & Ewert 2018)
 - Observation: calibrated damped residual eddy viscosity at start of comp. too high (no resolved fluctuations)
 - Solution: eddy viscosity scaled with resolved TKE from moving-averaging of solution
 - Moving-average time scale from background RANS

$$\tilde{k}_t^{n+1}(\mathbf{x}) = \theta \tilde{k}_t^n(\mathbf{x}) + (1 - \theta) \frac{|(\mathbf{v}')^n|^2}{2}, \quad \theta := \exp\left(-\frac{\Delta t}{\tau_s}\right)$$



Example high-lift slat-noise: frequency content too high at slat, i.e., low frequency noise missing

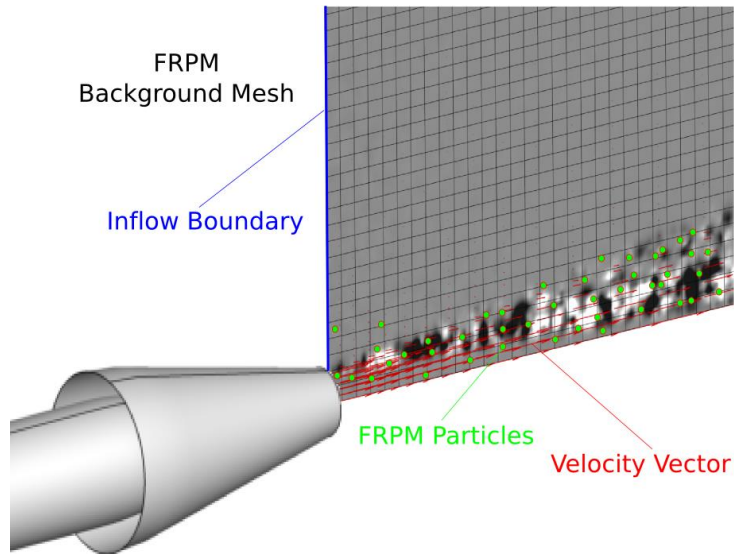
$$\nu_r^* = \sqrt{\frac{\tilde{k}_t}{k_t}} \nu_r$$



Modified dissipative model: proper sound generation @slat

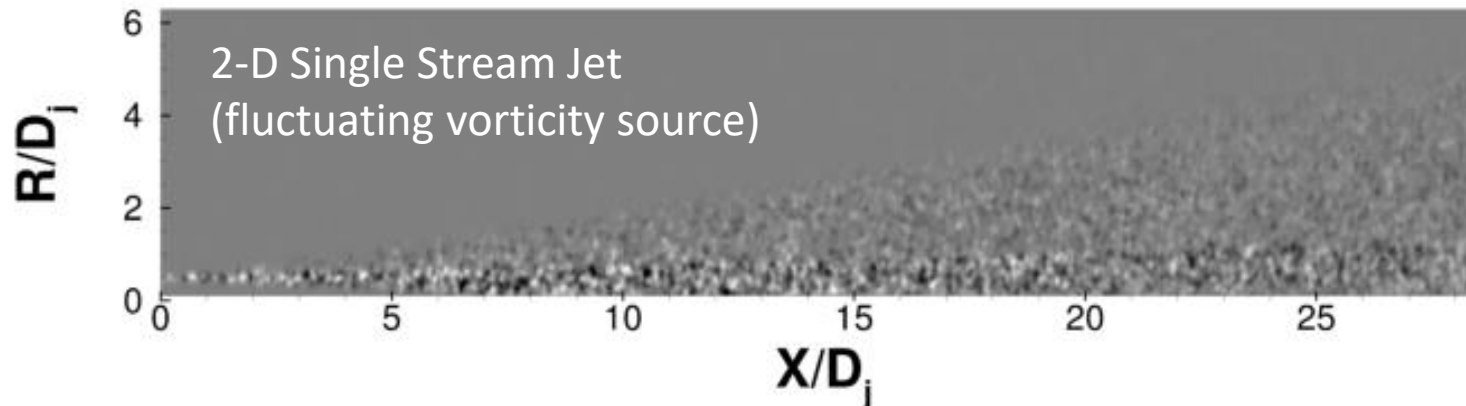
3+1-D Stochastic Backscatter Forcing

FRPM: Fast Random Particle-Mesh Method



- Velocity field divergence-free (solenoidal)
- Convection property of vorticity in non-uniform mean-flow
- Local reconstruction of target turbulent kinetic energy from RANS
- Local reconstruction of turbulent time- & length-scales from RANS
- Broadband spectra realized
- Computationally efficient / parallelized

AIAA Pap. 2005-2862 / AIAA Pap. 2006 / Comp. & Fluids 37 (2008) / AIAA 2007-3506 / AIAA 2009-3369 / AIAA 2009-3175 / JSV 330 (2011) / Comp. & Fluids 132 (2016)



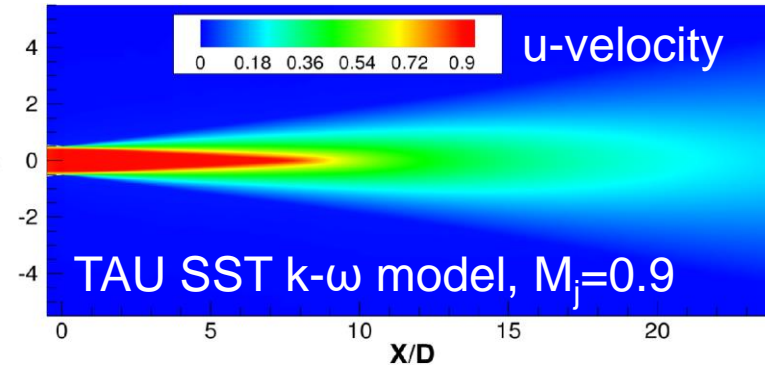
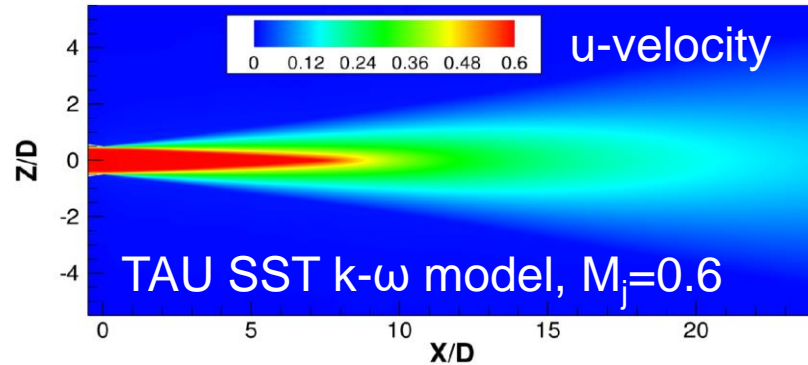
Outline

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- **Results for cold isolated single stream jets**
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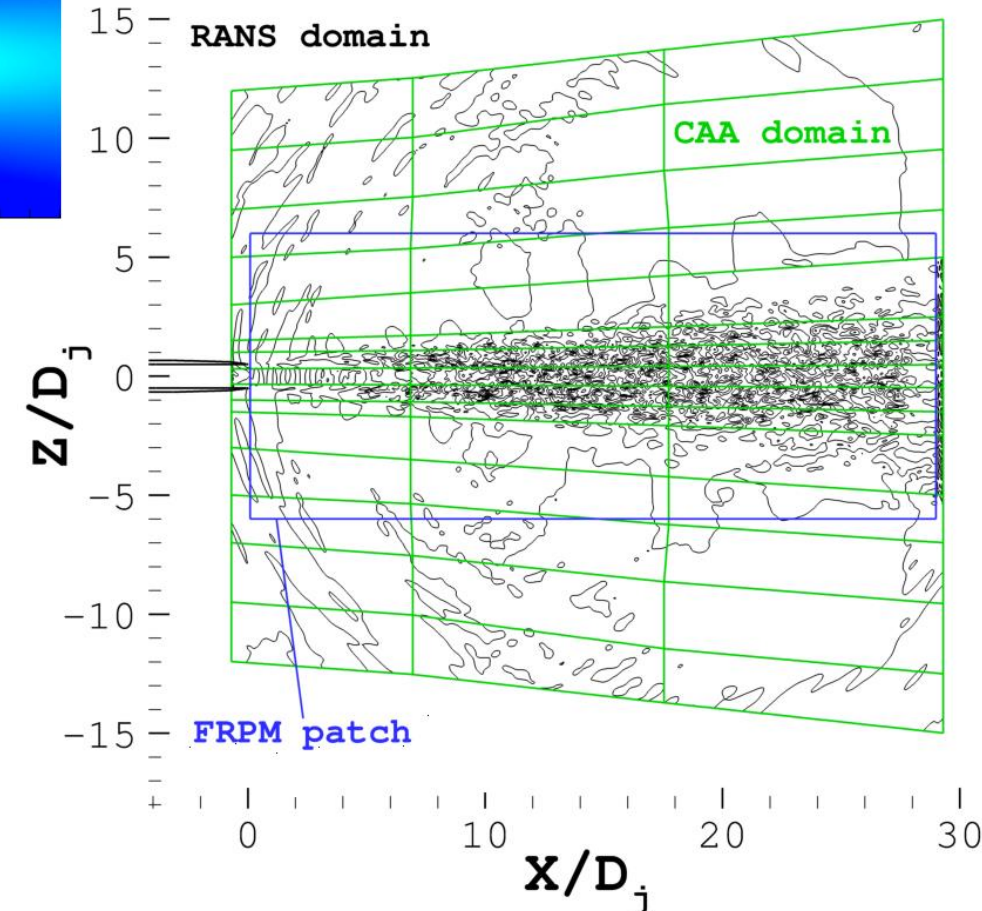


Results for cold isolated single stream jet

Computational Setup for Single Stream Jet FES



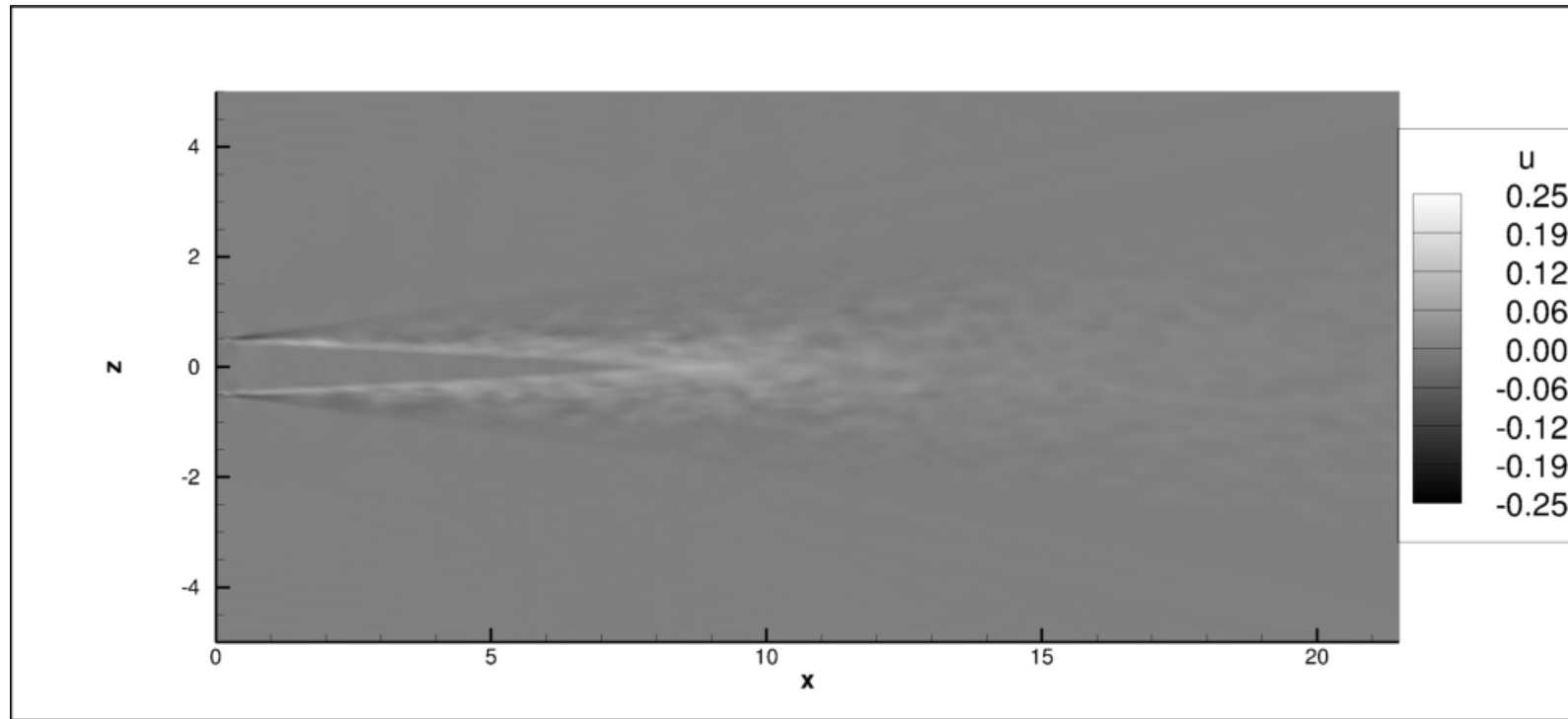
- Two computational cases are used to validate the FES approach
- 3-D cold isolated single stream jet (SSJ) computation for nozzle exit Mach numbers 0.6 ($Re=0.68M$) and 0.9 ($Re=1M$)
- Both SSJ in static condition
- CAA mesh with 87 grid blocks and approx. 9.7 Mio. grid points (O-mesh topology, approx. 20kHz Resolution, $St \sim 3.6$)
- FRPM patch is resolved with approx. 2.6 Mio grid points
- Computations are performed on 88 CPUs with runtime of 4-7 days
- FES domain excludes nozzle interior



Results for cold isolated single stream jets II

Flow Properties

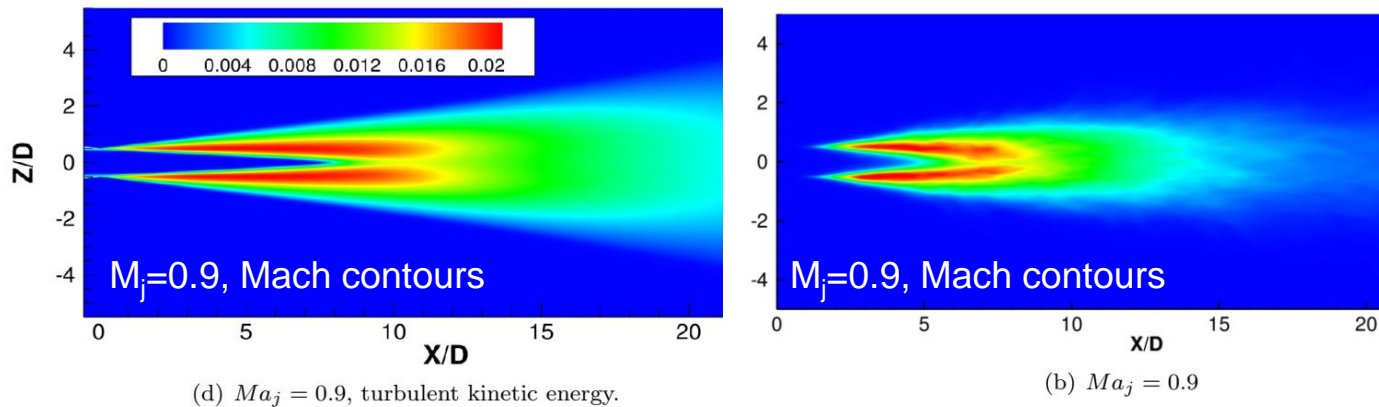
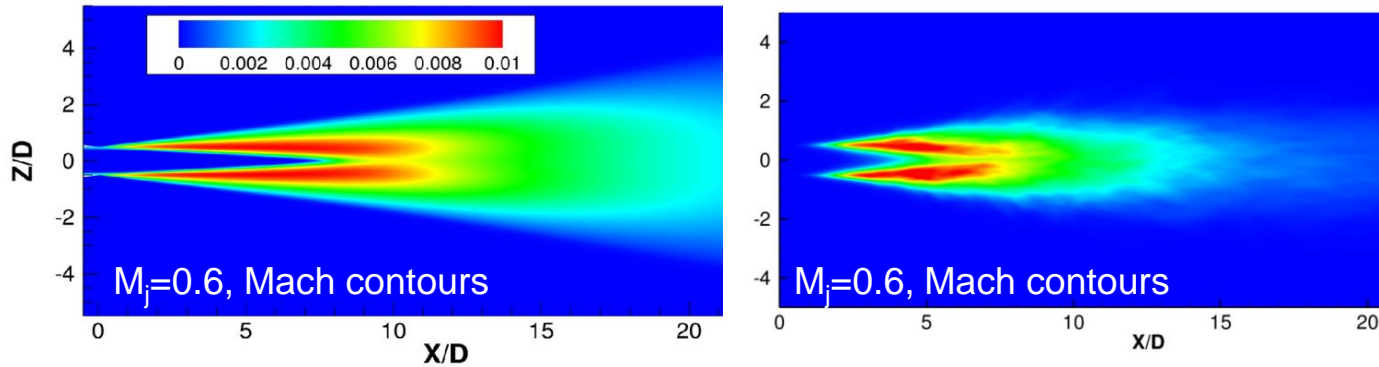
u-velocity $M_j=0.9$ jet



Results for cold isolated single stream jets III

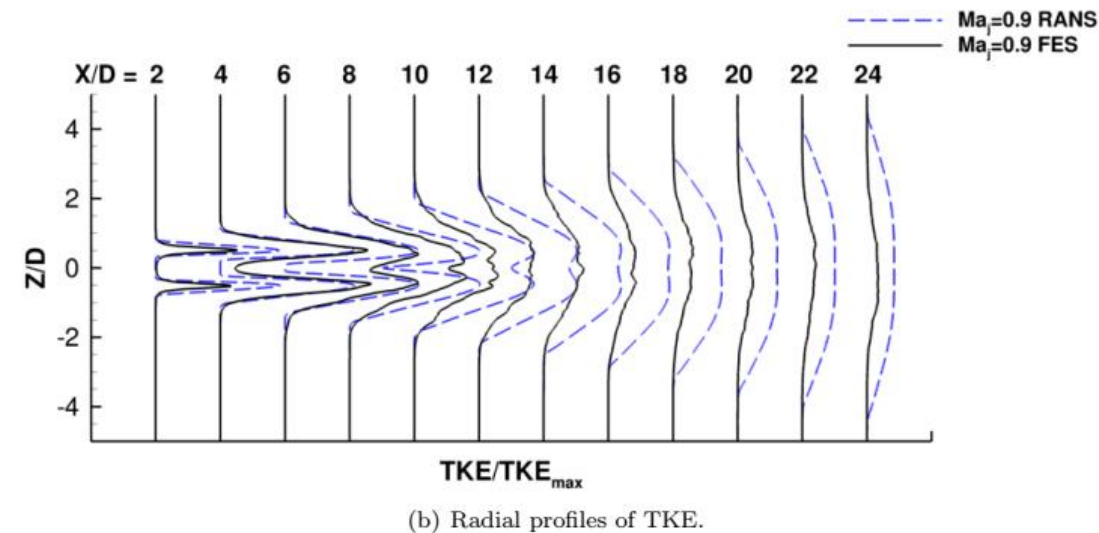
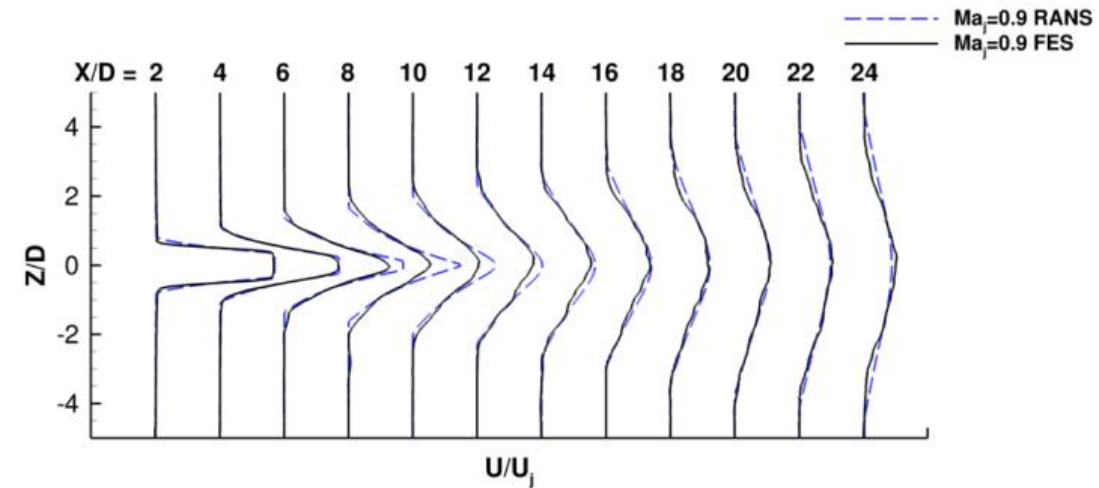
Flow Properties

- Dissipation extended by mean-part: too fast decay observed



RANS TKE

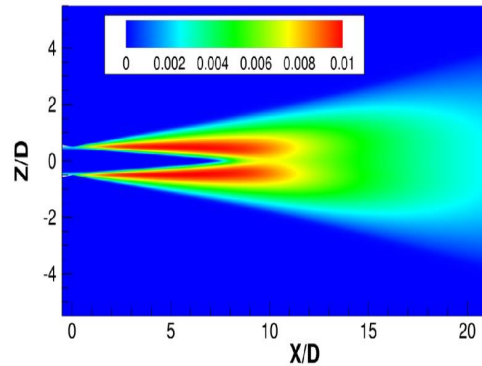
FES Resolved TKE



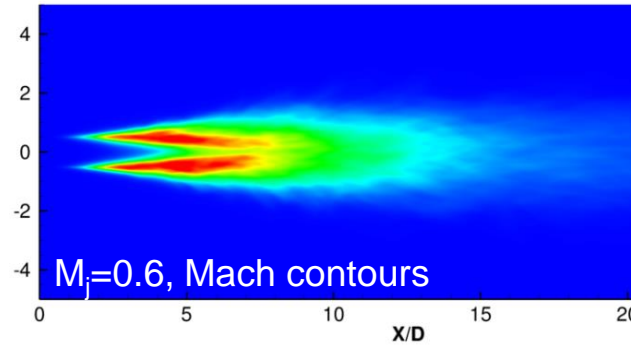
Results for cold isolated single stream jets III

Flow Properties

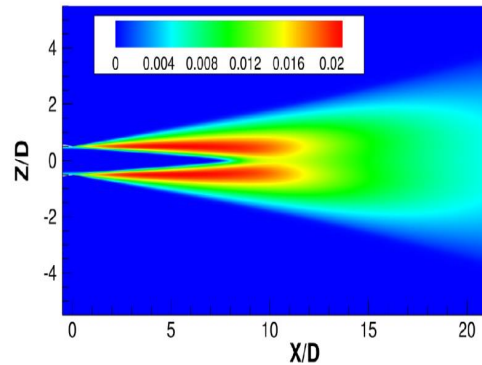
- Dissipation extended by mean-part: too fast decay observed



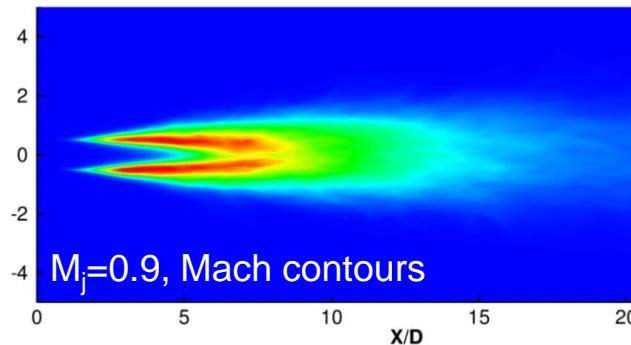
(c) $Ma_j = 0.6$, turbulent kinetic energy.



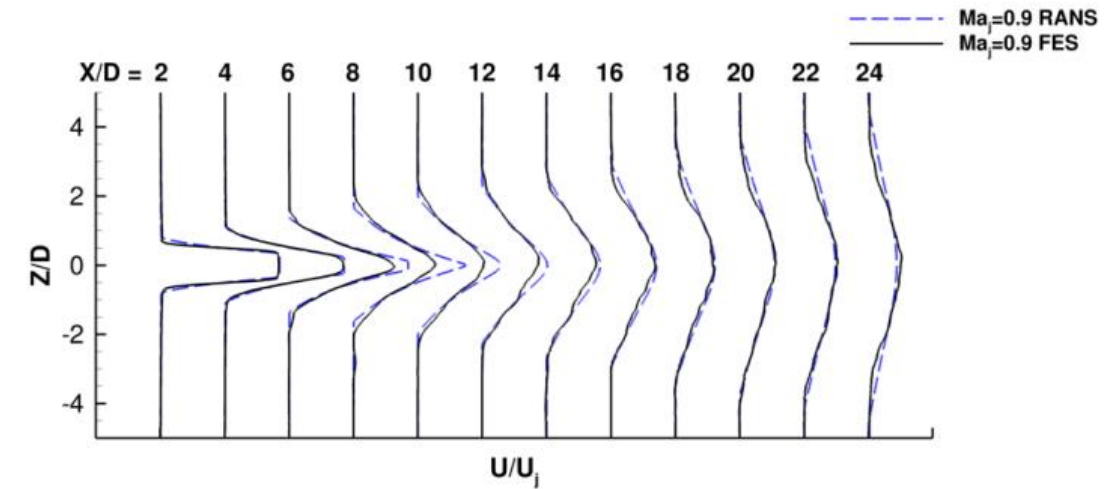
(a) $Ma_j = 0.6$



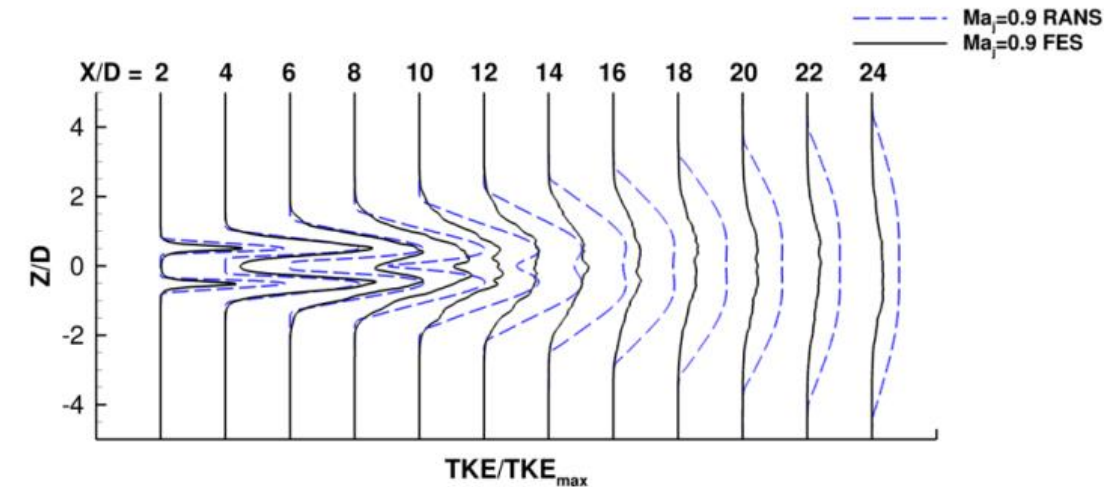
(d) $Ma_j = 0.9$, turbulent kinetic energy.



(b) $Ma_j = 0.9$



(a) Radial profiles of axial velocity.



(b) Radial profiles of TKE.

RANS TKE, 33.3% squeezed

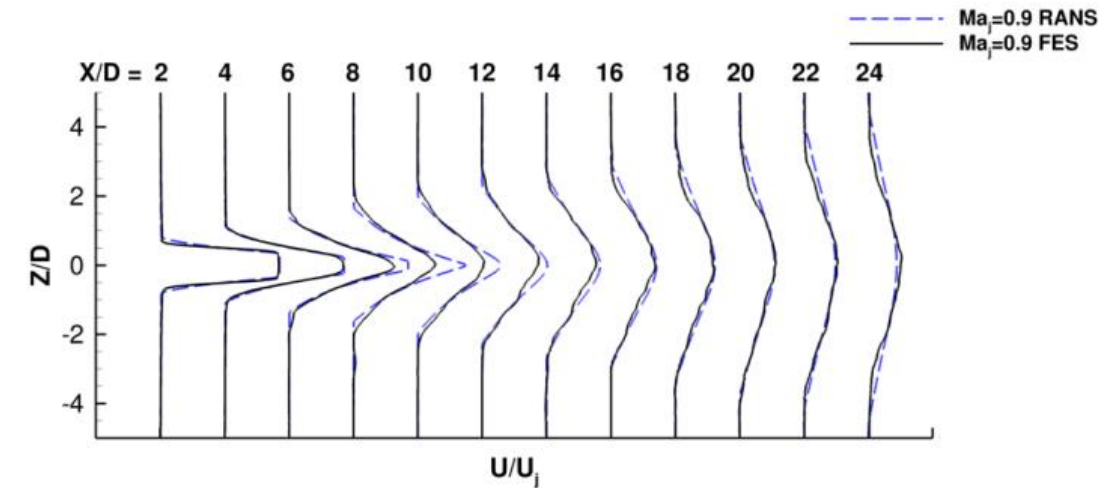
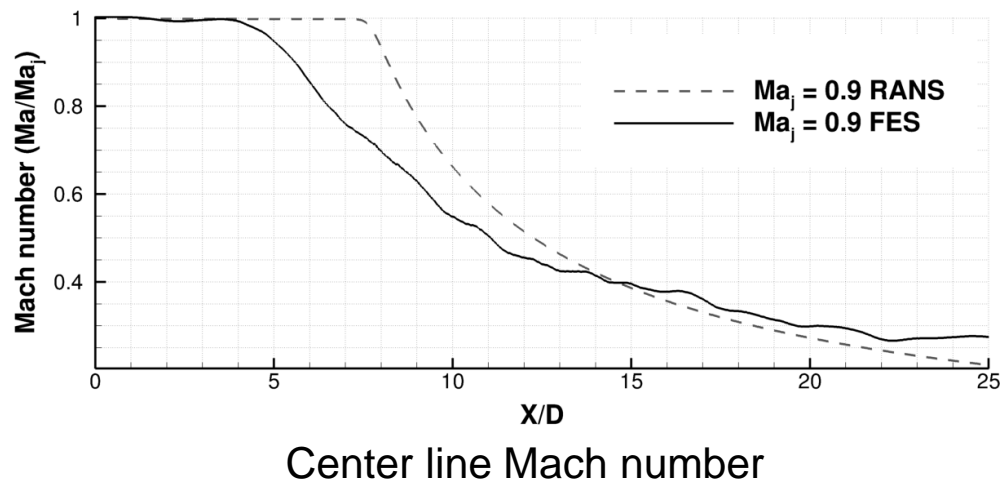
FES Resolved TKE



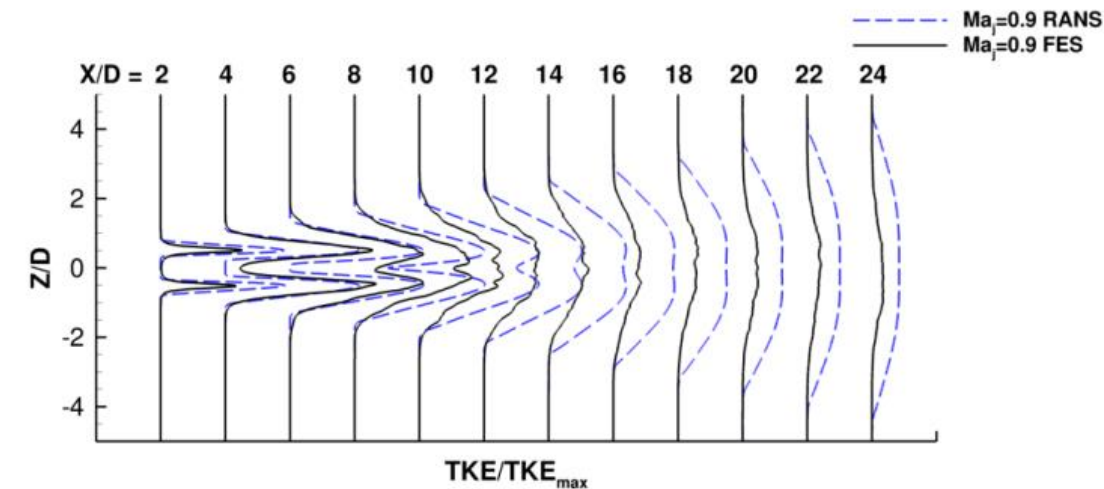
Results for cold isolated single stream jets IV

Flow Properties

- Potential core of RANS too long (7.5D) and drop-off too steep
- RANS mean-flow corrected by NLDE-FES
- Qualitatively right correction of center line Mach number
- FES Potential core (4.5D) slightly too short (4-6D expected)
→ may indicate slightly too strong shear layer mixing
 - stochastic forcing slightly too strong
 - too much filtering



(a) Radial profiles of axial velocity.



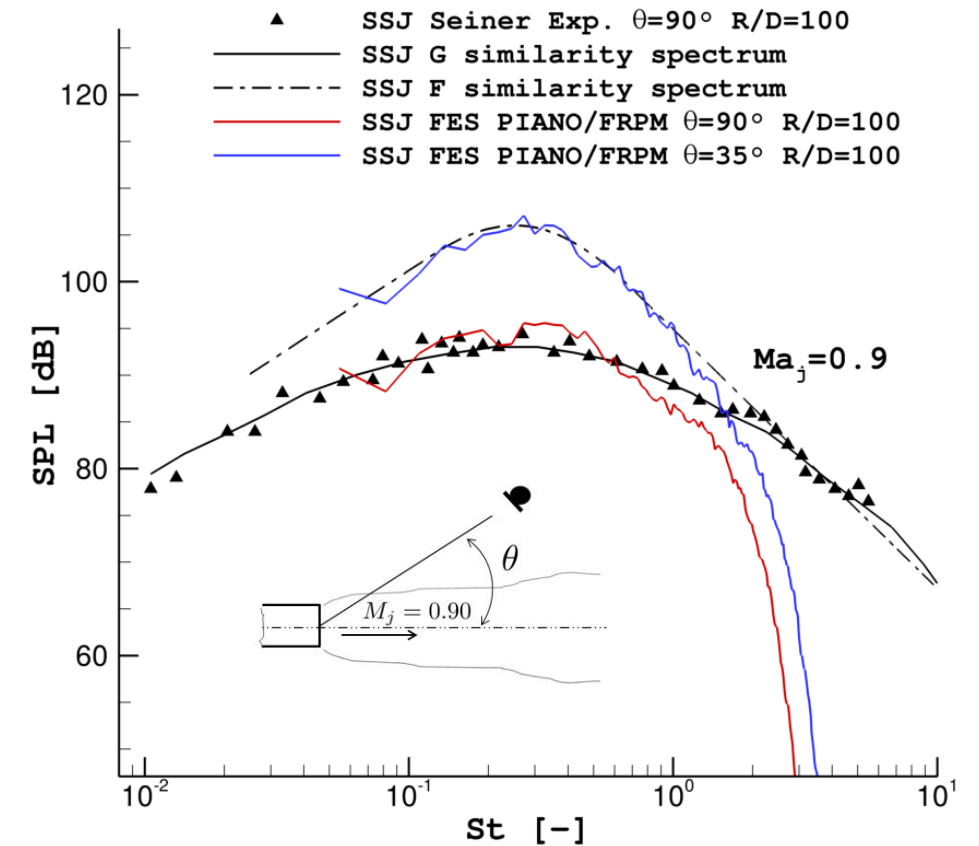
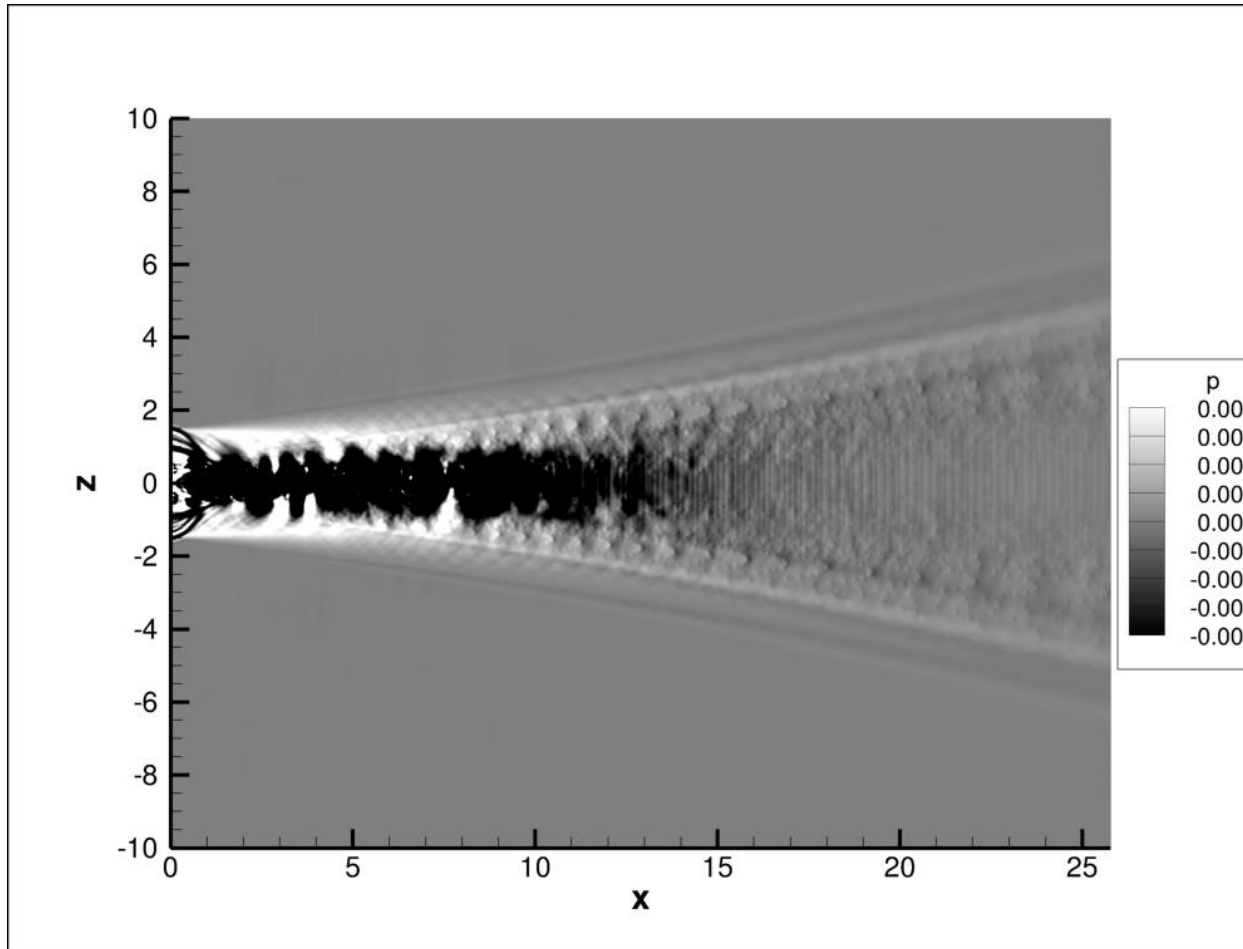
(b) Radial profiles of TKE.



Results for cold isolated single stream jets V

Resolved Sound Field $M_j=0.9$

p' , fluctuating pressure $M_j=0.9$ jet



Comparison of spectra with reference

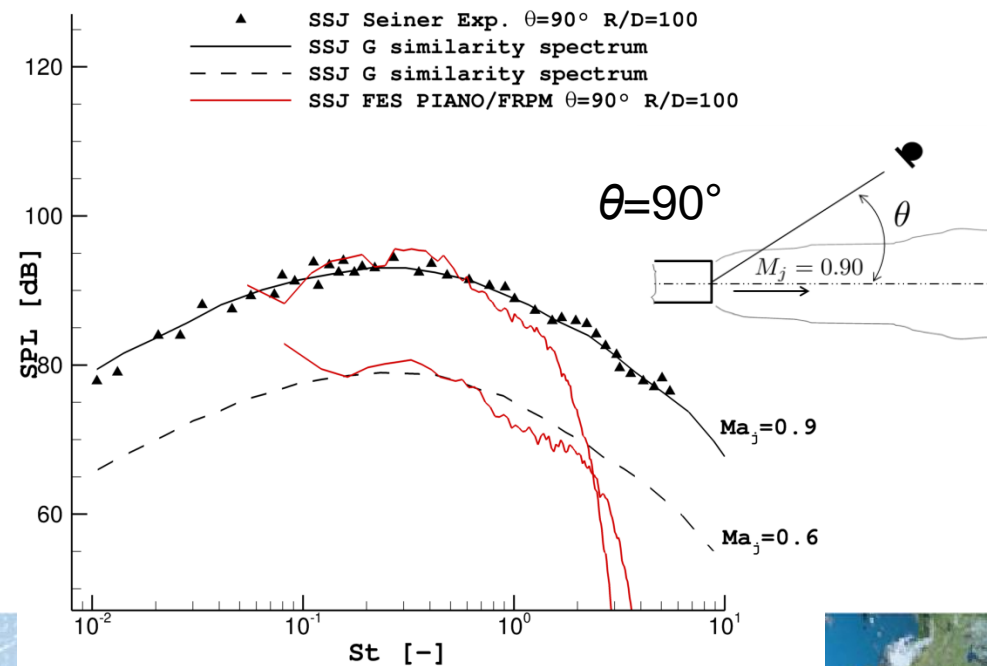
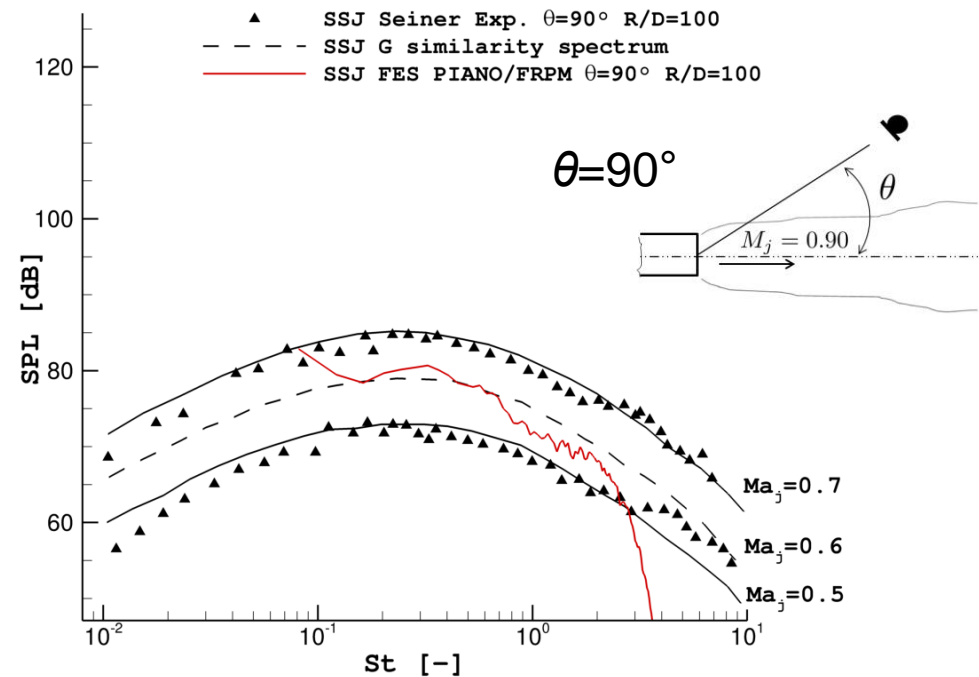
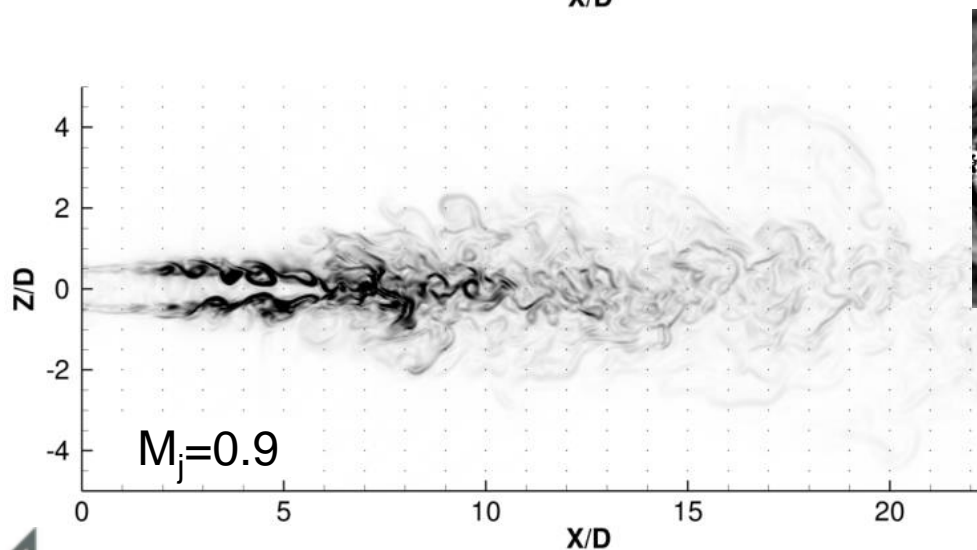
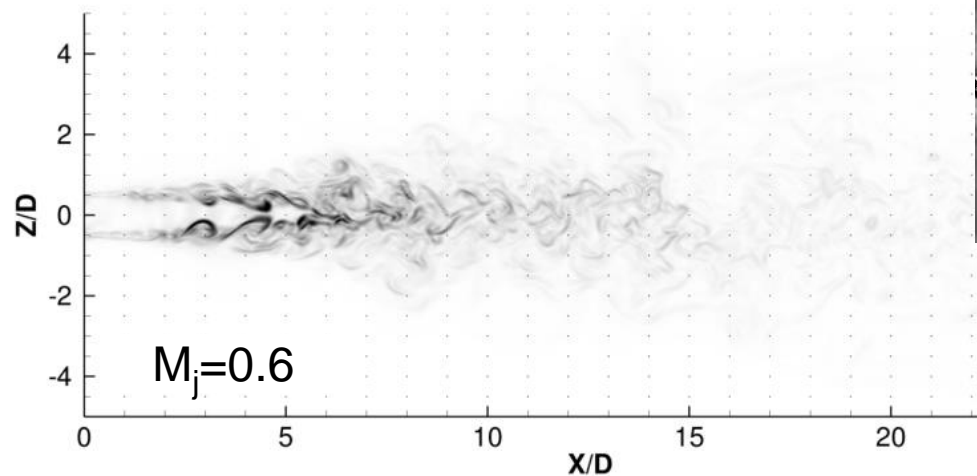
- No enhanced laminar-to-turbulent shear-layer transition and increased noise due to vortex pairing



Results for cold isolated single stream jets VI

Flow Properties

- Right Mach-number scaling



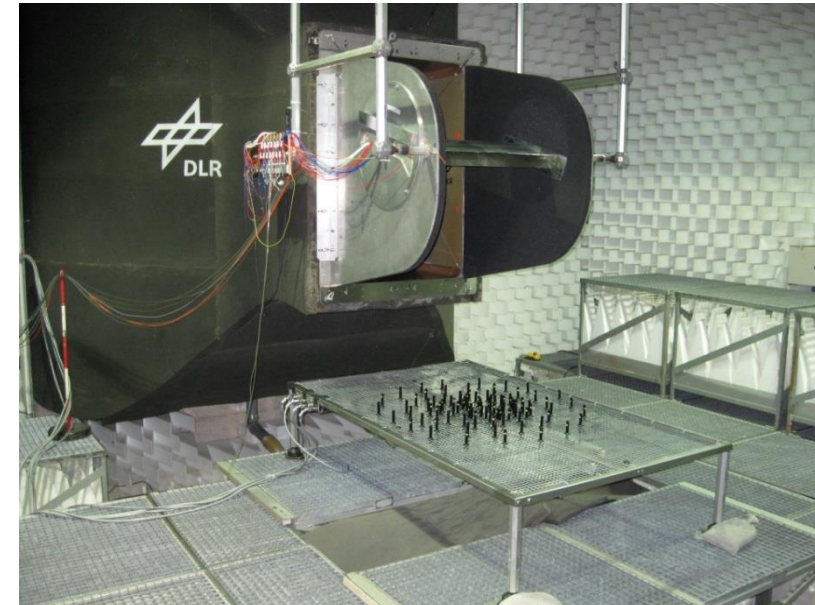
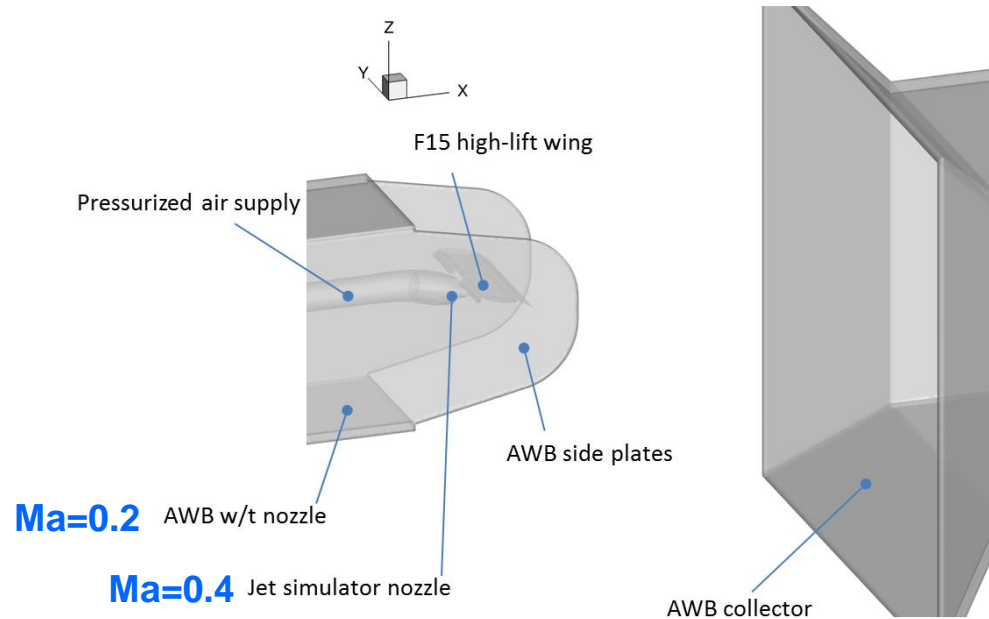
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Results of installed UHBR configuration

Installed UHBR configuration in AWB wind tunnel I



Experimental setup of LIST configuration

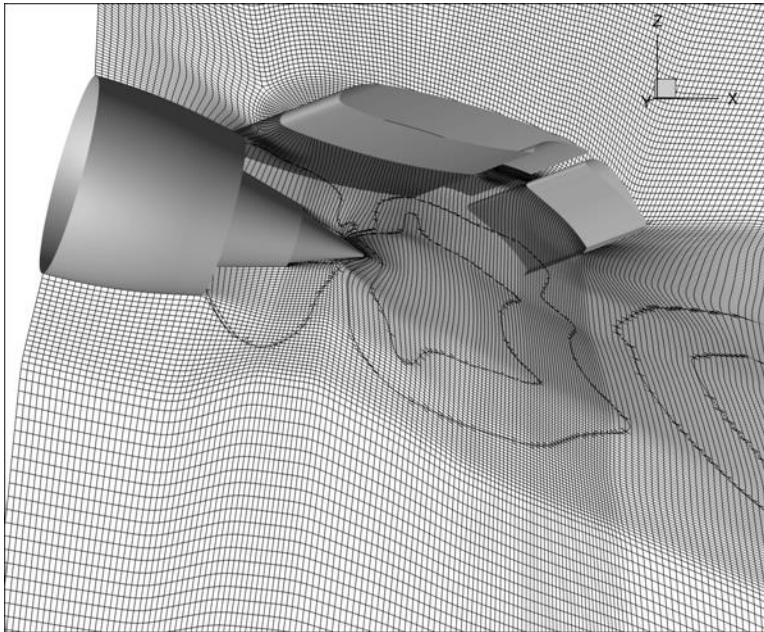
- Numerical setup is reproducing the experimental setup of AWB for installed UHBR configuration
- Computational domain is limited to area of interest, i.e. between the AWB nozzle and collector; in spanwise direction between the side plates
- Flow conditions are emulating the approach conditions with flap angle 25°



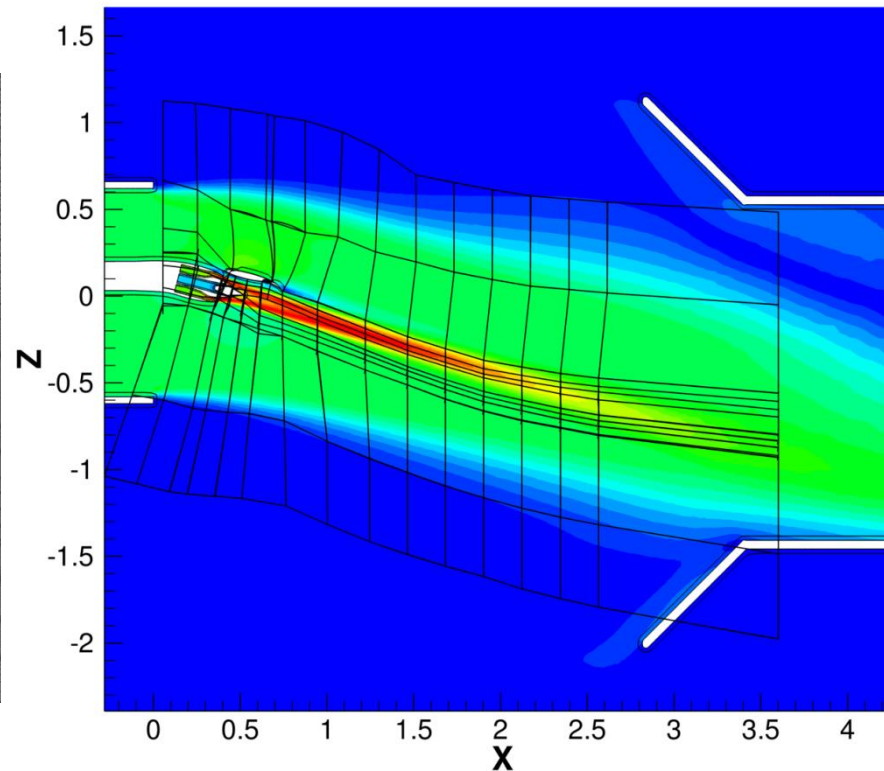
Results of installed UHBR configuration

Computational Chain

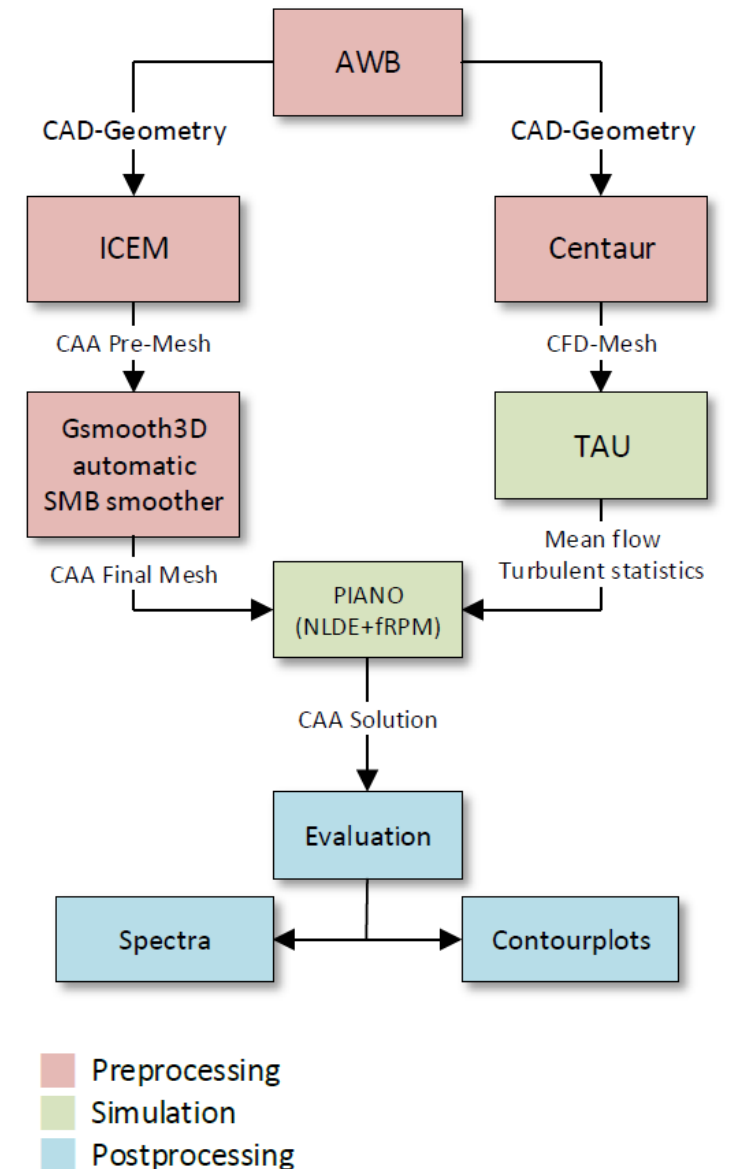
- Smoothed SMB mesh, AWB jet-wing installation, 350 blocks, 30mio points



CAA mesh

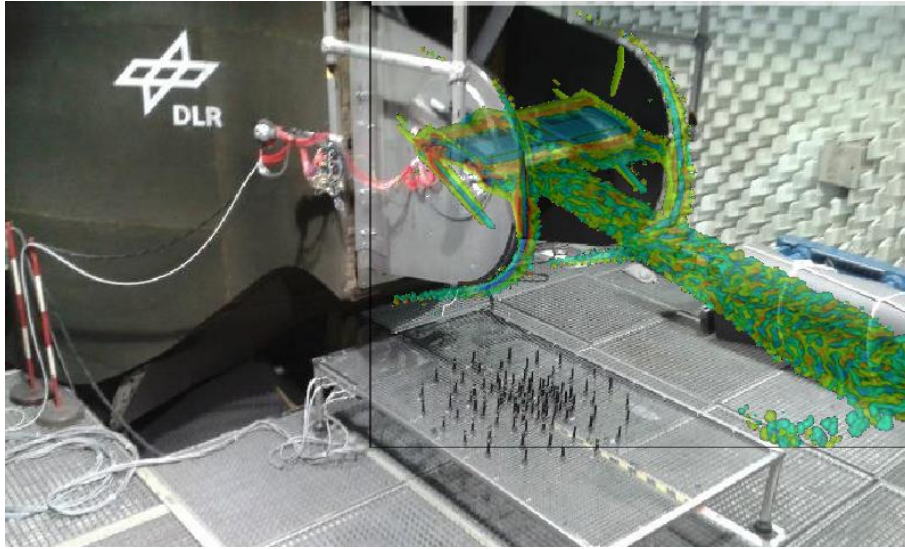


TAU RANS mean-flow and CAA mesh

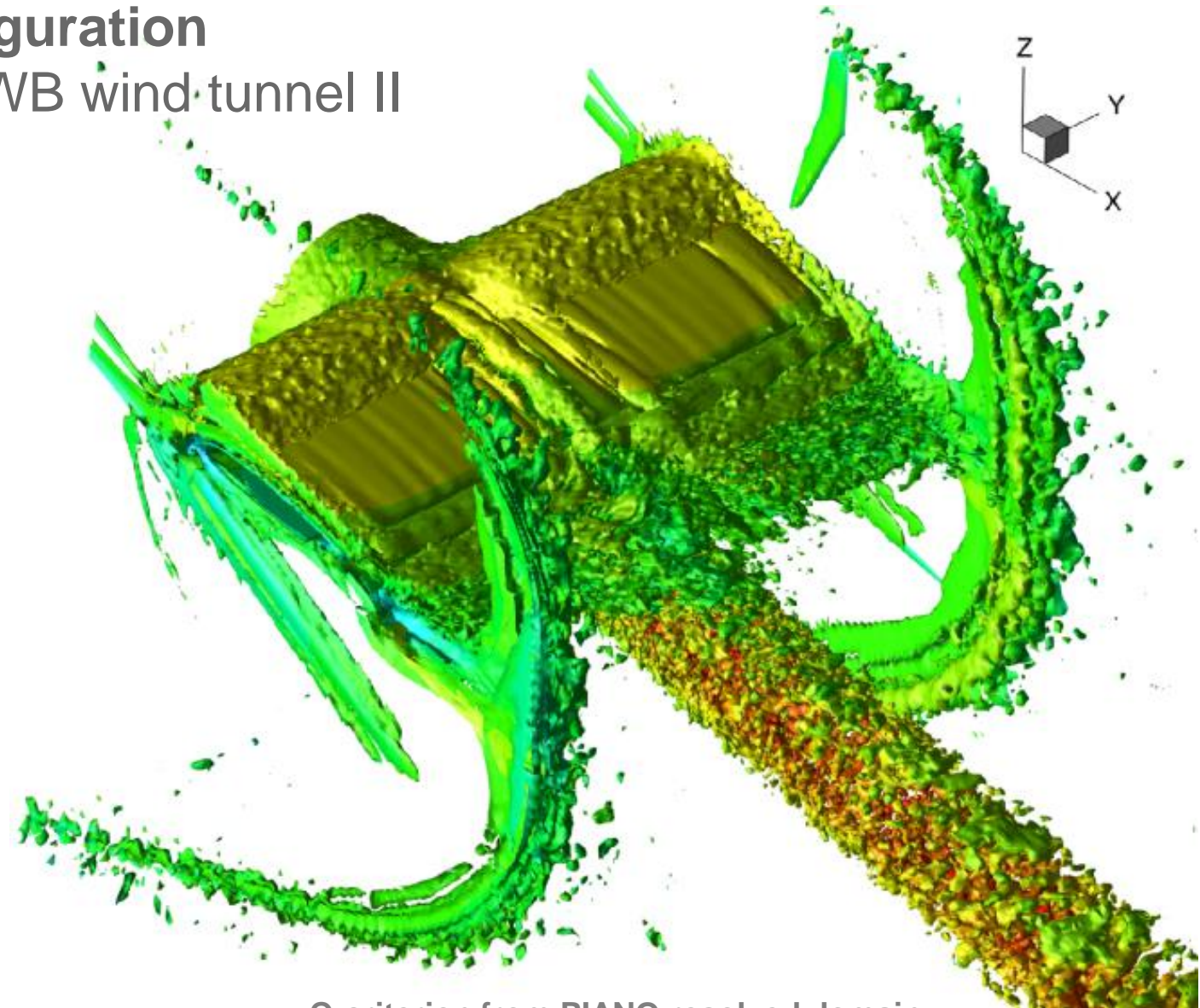


Results of installed UHBR configuration

Installed UHBR configuration in AWB wind tunnel II



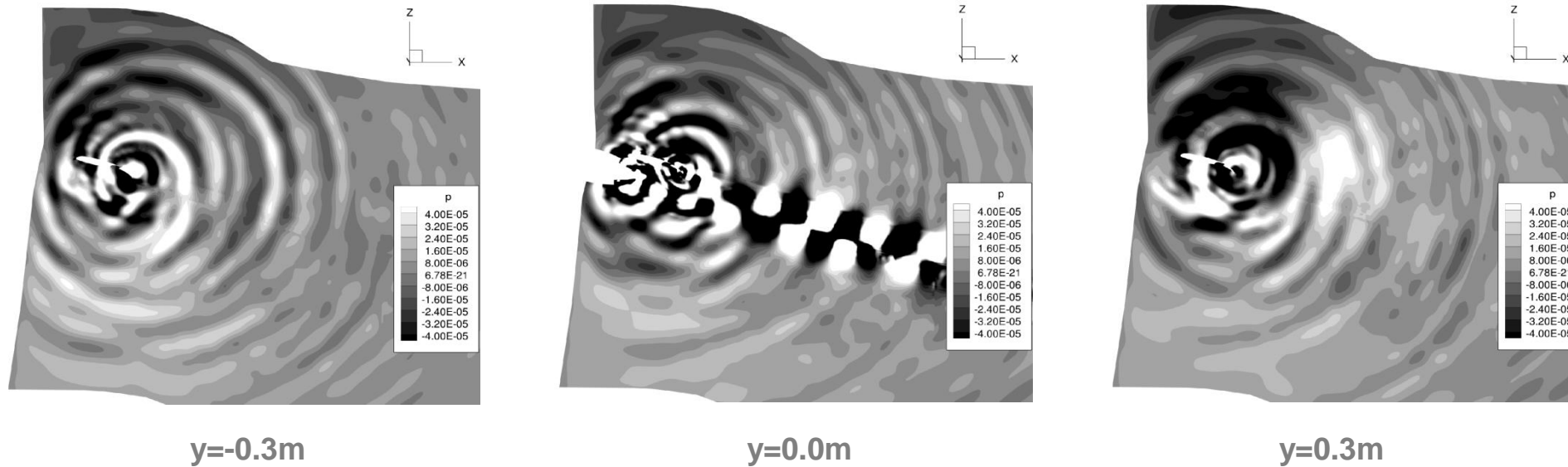
- Q-criterion (iso-surfaces) colored with the flow velocity
- Spanwise extend of computational domain does not include w/t side plates but are closely positioned to it
- Appearance of wake structures from the trailing edges of side plates
- Vortex shedding on the upper side of the wing is triggered by close position of the UBHR nacelle0



Q-criterion from PIANO resolved domain

Results of installed UHBR configuration

Installed UHBR configuration in AWB wind tunnel III



- Contour plots of pressure fluctuations at three different spanwise cuts; in the symmetry plane of UHBR nozzle and to the left/right with 0.3m separation
- Strong circular noise radiation is observed with its origin at the flap trailing edge
- The hydrodynamic pressure fluctuations of the UHBR jet are observed in the cut at the centerline
- Noise radiation to the bottom side appears to be less intensive compared to region above wing



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Summary & Conclusions

- **Method**

- Application of stochastic backscatter method to technical problem: Forced Eddy Simulation
- Successful transfer of DHIT calibration to single stream jet test case

- **Single stream jet**

- NLDE provide correction to RANS mean flow (potential core length)
- TKE level properly reproduced in magnitude by resolved fluctuations with respect to RANS
- Slight modification of the TKE topology
- Proper prediction of absolute level, spectra, directivity and Mach number scaling

- **Installed jet noise**

- Tool chain established
- RANS simulation of entire wind tunnel set up
- FES simulations running
- Jet-flap interactions identified in results



Thank You!



Knowledge for Tomorrow

