

# A dispersion improved CABARET scheme for linear acoustic propagation problems

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## Acknowledgement:

Drs Anton Markesteijn and Vasily Semiletov

## Summary

- Challenges of modelling in computational aeroacoustics
- CABARET solver for flow and noise modelling: our state-of-the art
- CABARET, dispersion improved version and flux correction possibilities
- Numerical examples: 1d and 2d test problems

Jet noise example: turbulent jet is a very non-efficient sound generator, less than  $\sim 1/100,000$ th of its mechanical energy transforms into sound: effect of sound cancellation

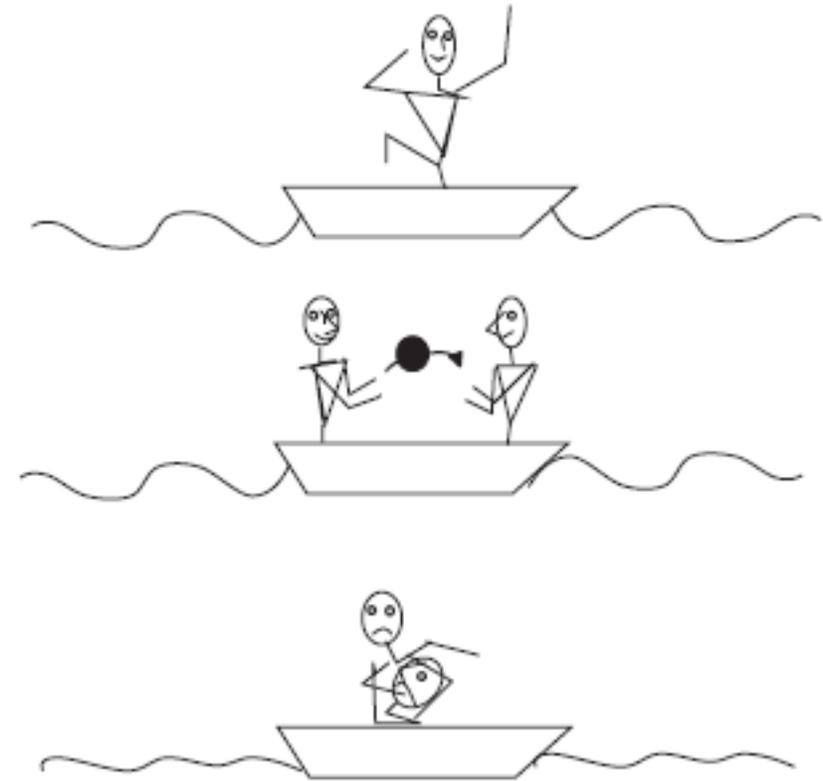
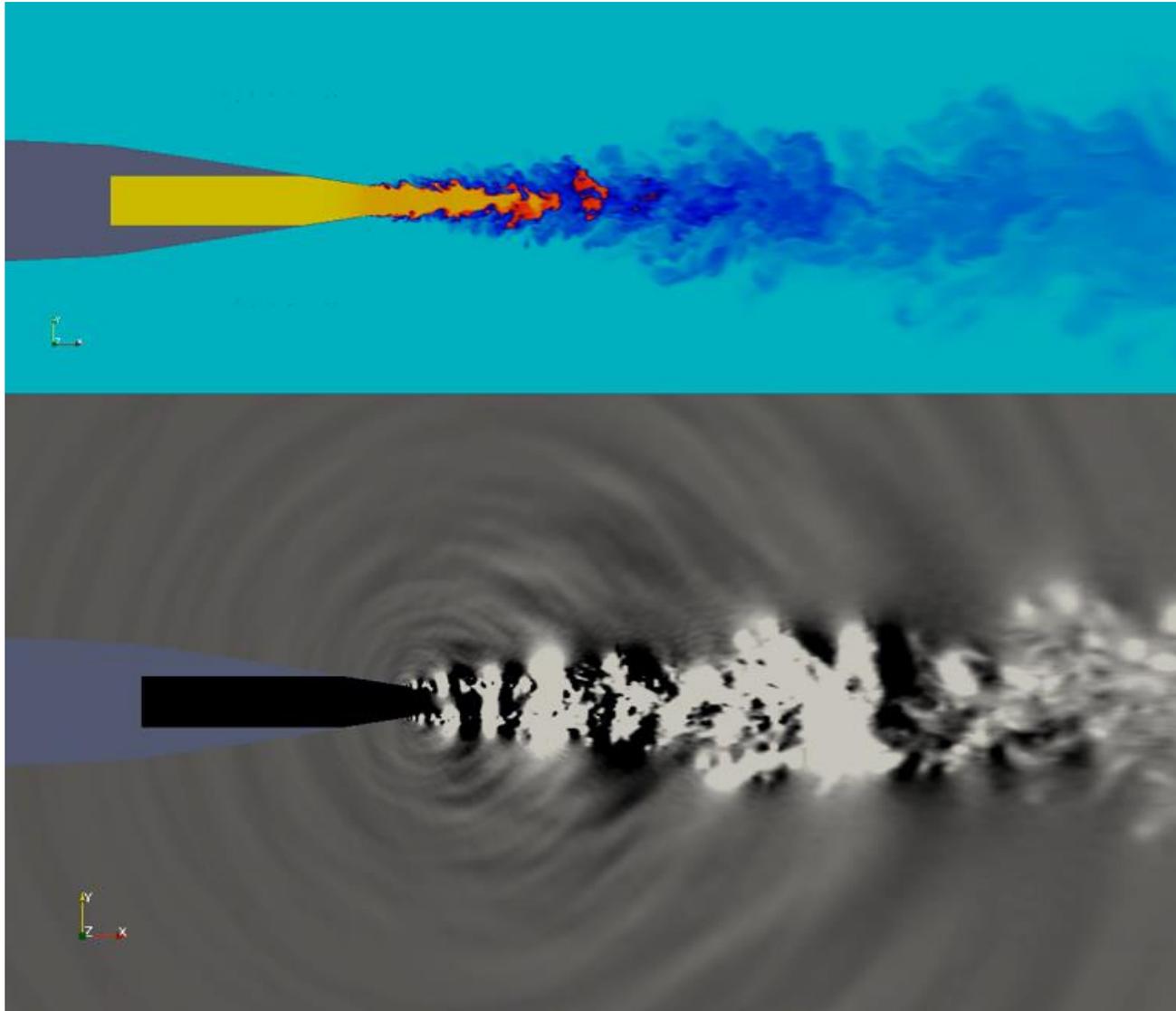


Figure 1: Monopole, dipole and quadrupole generating waves on the surface of the water around a boat.

# Von Neumann linear analysis for linear wave propagation

$$\frac{\partial u}{\partial t} + A \frac{\partial u}{\partial x} = 0$$

Represent the solution by a Fourier harmonics

$$u(x, t) = \frac{1}{2\pi} \sum_k u_k(t) e^{-ik \cdot x}$$

$$\frac{\partial u_k}{\partial t} + iA \cdot k \cdot u_k = 0 \Rightarrow u_k = C \cdot e^{i \cdot \omega \cdot t}, \underline{\omega = A \cdot k} \text{ Dispersion relation}$$

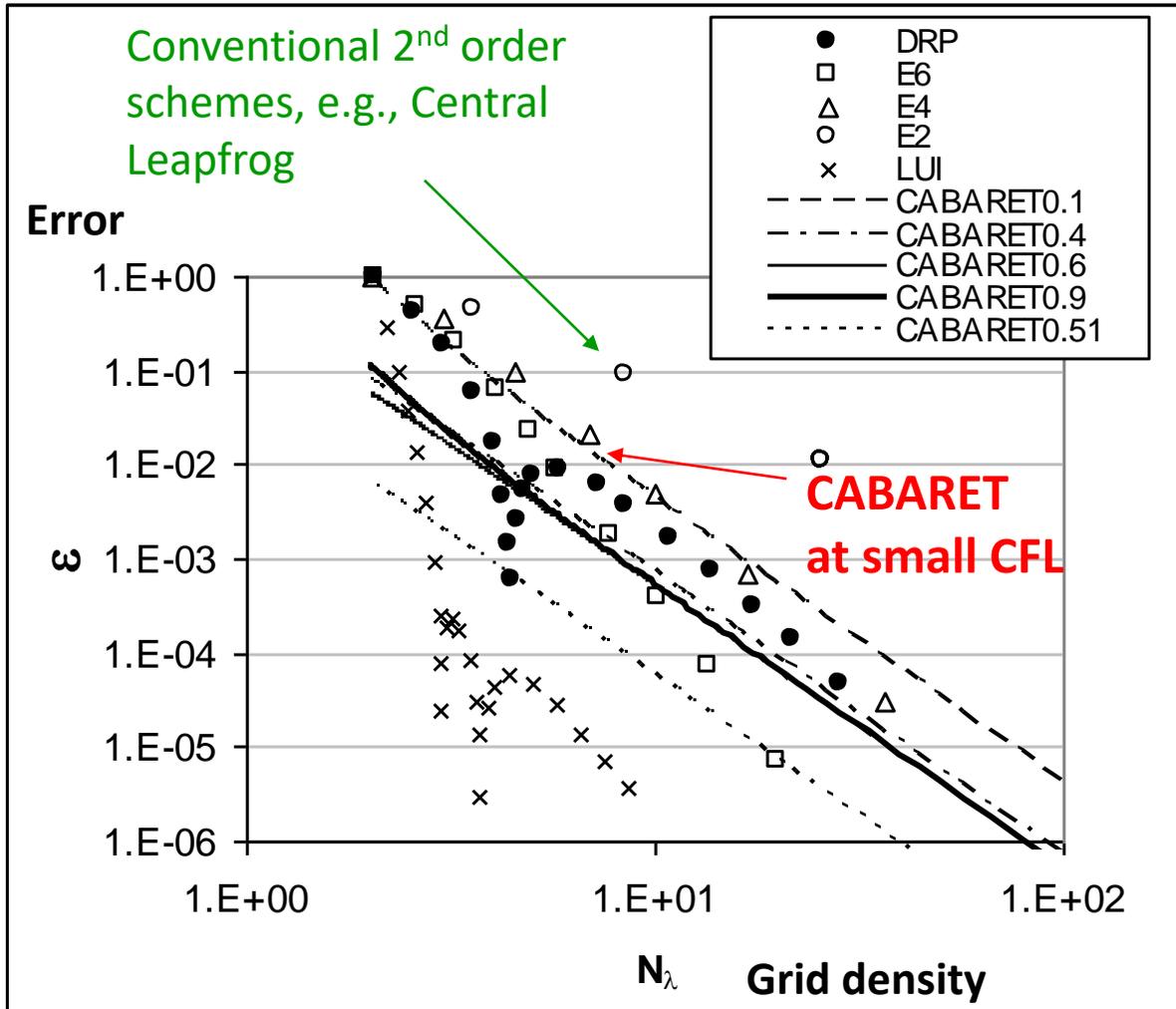
For real frequencies (A=real) the amplitude of simple waves is preserved

# Comparison of linear dispersion errors of several finite-difference schemes in terms of grid points per wavelength (P.P.W.)

Dispersion error due to numerical discretisation

$$\varepsilon = \left| \frac{\omega}{Ak} - 1 \right|$$

Conventional 2<sup>nd</sup> order schemes, e.g., Central Leapfrog



Ex –central differences of the x-th order  
 LUI – 6<sup>th</sup> order pentadiagonal compact scheme by Lui and Lele  
 DRP – 4-th order DRP scheme by Tam and Webb  
 CABARET<sub>x</sub>, CABARET scheme at CFL=x

PS: Can do a similar analysis for dissipation error

$$N_\lambda = \pi / (k \cdot h)$$

= Wavelength/ grid cell size

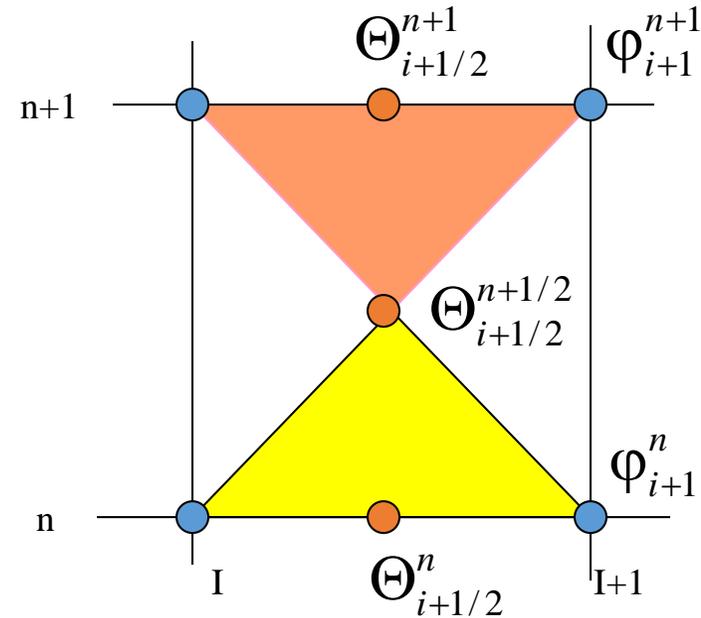
# Basic CABARET for linear advection using “active” conservation and flux variables

Goloviznin& Samarskii, 1998; Karabasov and Goloviznin, 2009

Calculate conservation variables at the mid-time step

$$\frac{\Theta_{i+1/2}^{n+1/2} - \Theta_{i+1/2}^n}{\tau/2} + c \cdot \frac{\varphi_{i+1}^n - \varphi_i^n}{h} = Q ;$$

$$\frac{\partial \varphi}{\partial t} + c \cdot \frac{\partial \varphi}{\partial x} = Q$$



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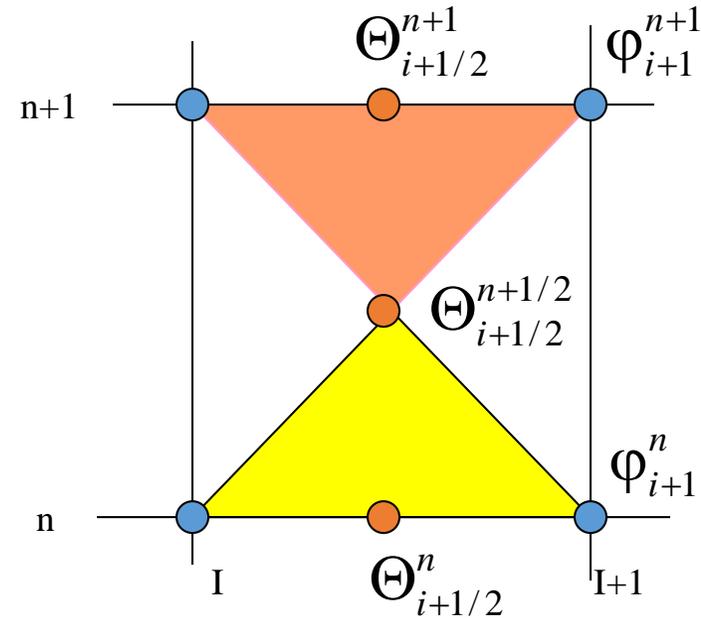
$$\frac{\Theta_{i+1/2}^{n+1/2} - \Theta_{i+1/2}^n}{\tau/2} + c \cdot \frac{\varphi_{i+1}^n - \varphi_i^n}{h} = Q ;$$

Upwind extrapolation of the flux variables

at the new time level

$$\varphi_{i+1}^{n+1} = 2 \cdot \Theta_{i+1/2}^{n+1/2} - \varphi_i^n \quad C > 0$$

$$\frac{\partial \varphi}{\partial t} + c \cdot \frac{\partial \varphi}{\partial x} = Q$$



# Basic CABARET for linear advection using “active” conservation and flux variables

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Upwind extrapolation of the flux variables

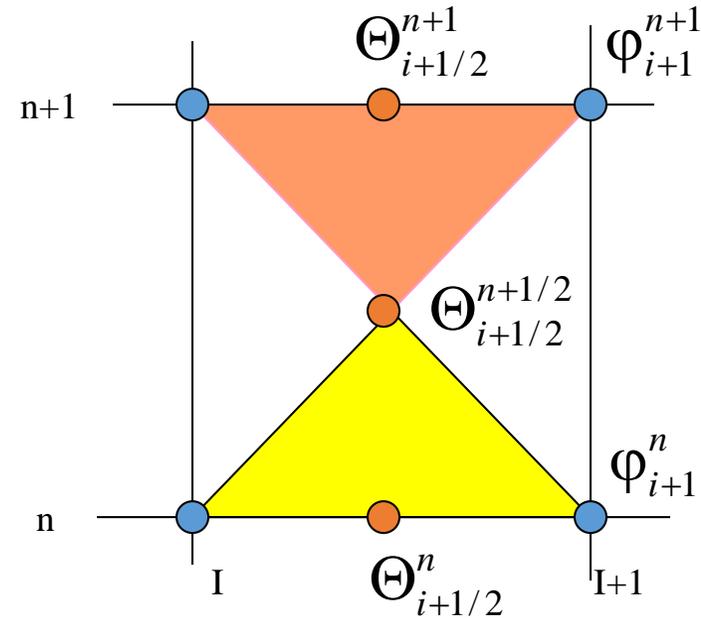
at the new time level

$$\varphi_{i+1}^{n+1} = 2 \cdot \Theta_{i+1/2}^{n+1/2} - \varphi_i^n \quad C > 0$$

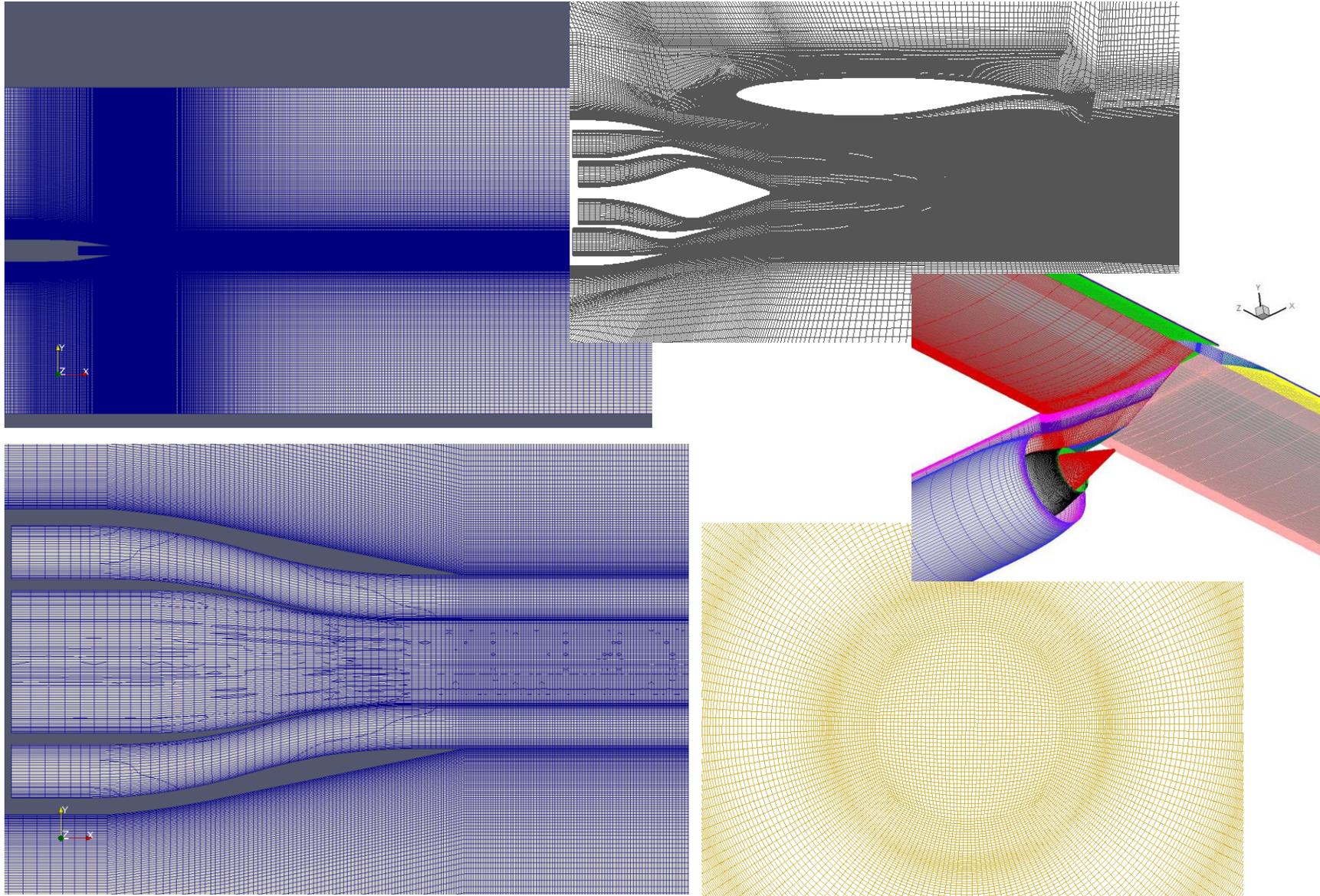
Update conservation variables at the new time step

$$\frac{\Theta_{i+1/2}^{n+1} - \Theta_{i+1/2}^{n+1/2}}{\tau/2} + c \cdot \frac{\varphi_{i+1}^{n+1} - \varphi_i^{n+1}}{h} = Q ;$$

$$\frac{\partial \varphi}{\partial t} + c \cdot \frac{\partial \varphi}{\partial x} = Q$$

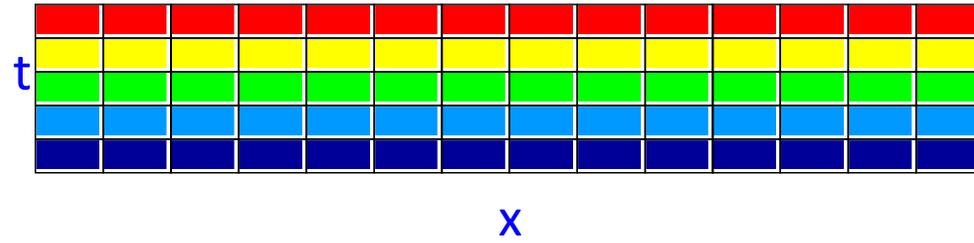


# Examples of grids

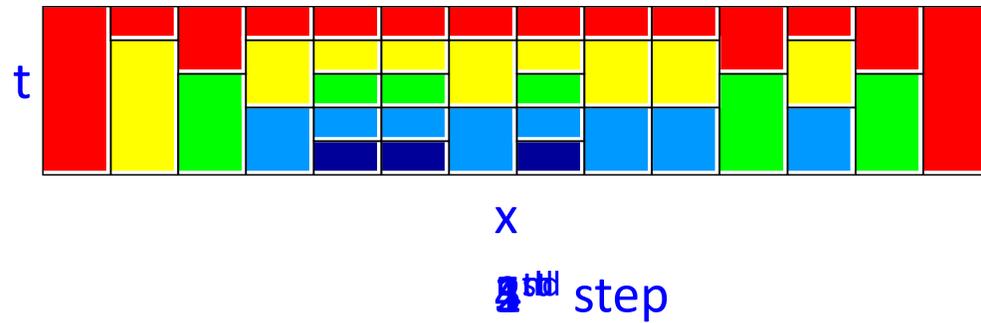


# Time-stepping

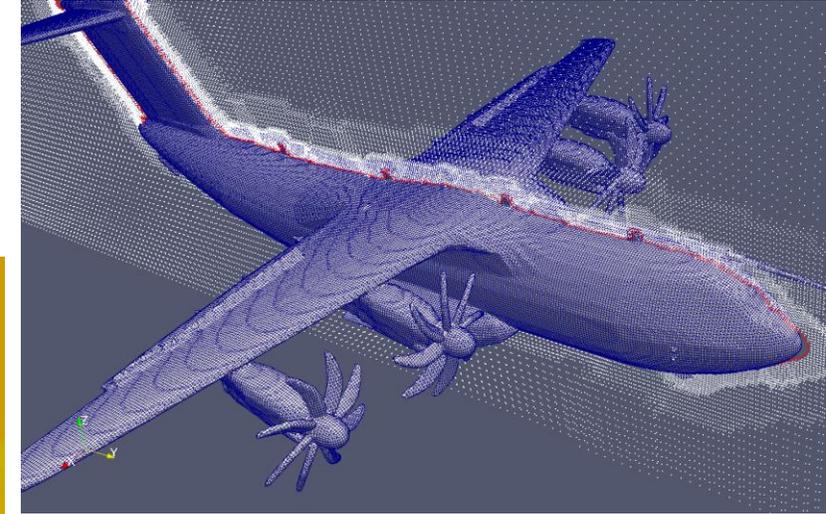
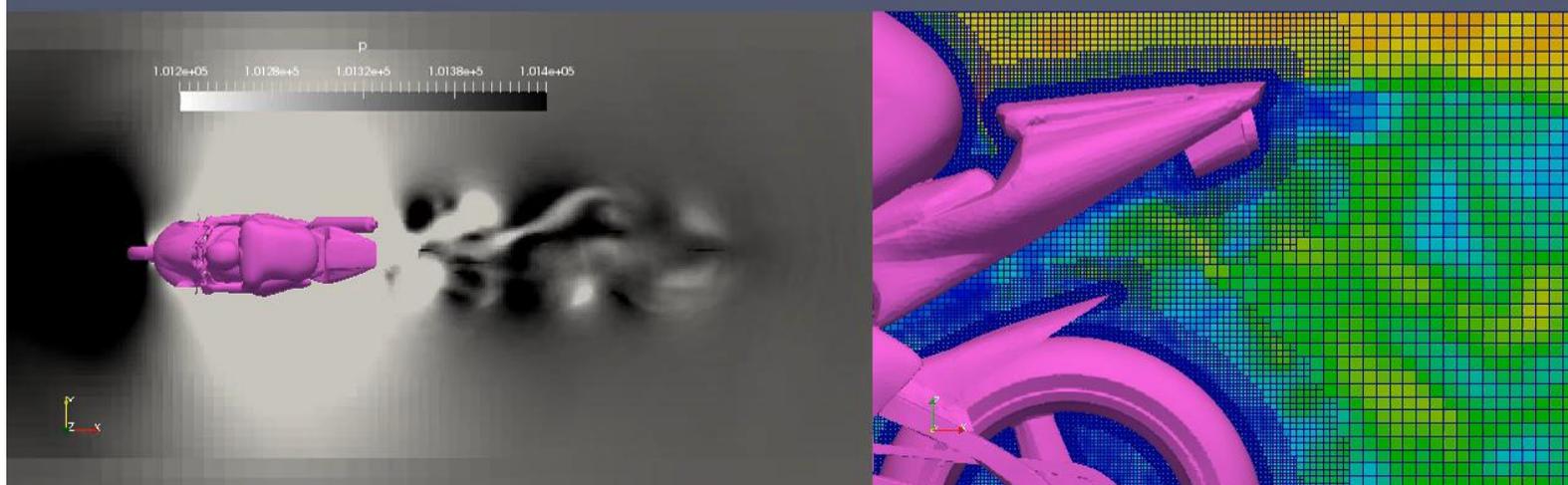
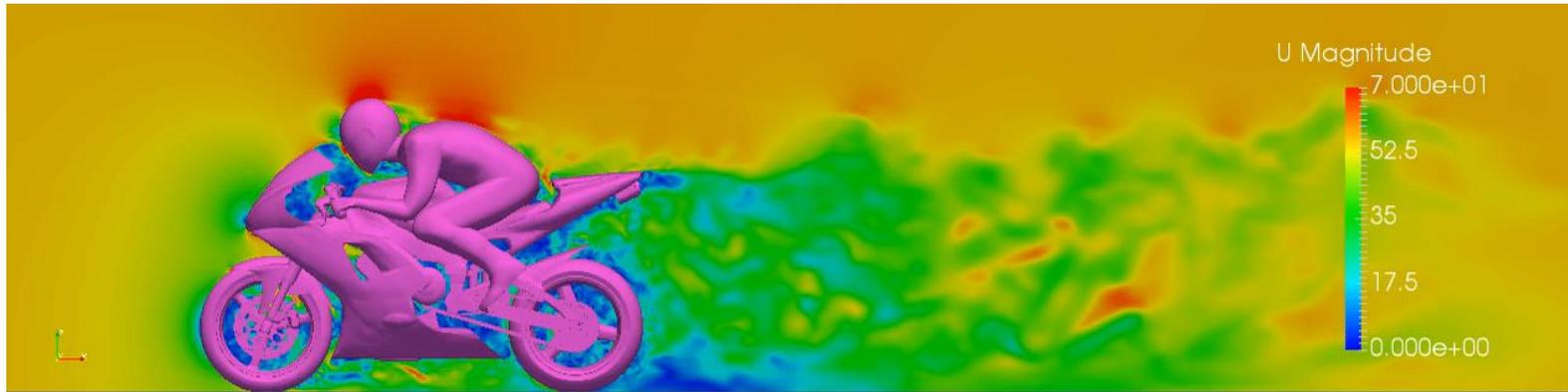
## Homogeneous



## Asynchronous



20 mln cell LES calculation on a snappy-hex mesh: 2GPU cards ~ 2 days simulate ~ 50 TUs  
“per time step, 2 GPUs are about a factor 55 quicker than a code like OpenFOAM on a 16-core cluster node”



# GPU Solver for Fast-Turn-Around Flow and Noise Calculations: our state-of-the art

*“HiFi prediction for validation of designs is mandatory”*

- Unstructured Hexagonal mesh + snappy-hexa + limited number of 1<sup>st</sup> order tetras is possible (OpenFOAM format)
- Some (old CPU) algorithms optimised for GPU usage
- Some algorithms completely rewritten from scratch
- Code is memory-optimised to fit on GPU and MPI-ed
  - 2.2 mln mesh per GB memory (6 GB~13mln)
- Asynchronous algorithm speeds-up computations
  - GPUs could do several magnitudes faster (>1000) than CPU
  - Same convergence rate of the original CABARET method
  - Decrease absolute error (CABARET’s optimal CFL condition)
- Typical grids we can handle (2 K80 GPUs) = 80–100 mln cells, typical run times (jet noise) = 2–4 days
- FWH acoustic modelling post-processor
- Goldstein acoustic analogy computation
- Results easily visualised with ParaView

Goal: optimise CABARET for linear dispersion properties at small CFL  
 (ok to tolerate a few new calibration parameters so long as they are well defined!)  
 while preserving conservation & flux correction

CABARET with an artificial dispersion term

$$\sim d^3u/dx^3$$

$$\frac{1}{2} \left( \frac{\hat{\varphi}_i - \varphi_i}{\tau} + \frac{\varphi_{i-1} - \check{\varphi}_{i-1}}{\tau} \right) + c \frac{\varphi_i - \varphi_{i-1}}{h} = \frac{c\mu}{h} [(\varphi_{i+1} - 2\varphi_i + \varphi_{i-1}) - (\varphi_i - 2\varphi_{i-1} + \varphi_{i-2})];$$

Goloviznin and Samarskii, 1998

Dispersion coefficient:  $\mu = \varepsilon (1-r)(1-2r);$

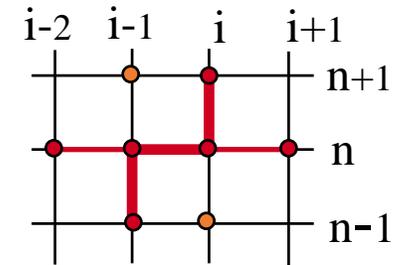
Characteristic equation:

$$q_k^2 - q_k (1 - e^{ikh}) [(1 - 2r) - 8\mu r (1 - \cos(kh))] - e^{ikh} = 0;$$

Non-dissipation (and stability) condition:

$$|q_1(kh, r)| \equiv 1 ; |q_2(kh, r)| \equiv 1 ; r \in [0, 1] ; kh \in [0, 2\pi];$$

satisfied for  $(-2 < \varepsilon < 0.29)$



$$\varepsilon = -0.08$$

# Dispersion improved CABARET in the flux form

Goloviznin and Samarskii, 1998

$$\frac{1}{2} \left( \frac{\hat{\varphi}_i - \varphi_i}{\tau} + \frac{\varphi_{i-1} - \check{\varphi}_{i-1}}{\tau} \right) + \frac{f_i - f_{i-1}}{h_i} = 0$$
$$f_i = \underbrace{c\varphi_i}_{\mathbf{F}} + c \frac{\mu_i}{\hbar_{i+1}} \left( \frac{\varphi_{i+1} - \varphi_i}{h_{i+1}} - \frac{\varphi_i - \varphi_{i-1}}{h_i} \right);$$

where

$$\hbar_i = 0.5(h_i + h_{i-1});$$

$$\mu_i = \varepsilon \left( h_i h_i - 3c\tau \hbar_i + 2c^2 \tau^2 \right);$$

$$\varepsilon = 0.08$$

# Dispersion improved CABARET in the flux form

Goloviznin and Samarskii, 1998

$$\frac{1}{2} \left( \frac{\hat{\varphi}_i - \varphi_i}{\tau} + \frac{\varphi_{i-1} - \check{\varphi}_{i-1}}{\tau} \right) + \frac{f_i - f_{i-1}}{h_i} = 0$$

$$f_i = c \varphi_i + c \frac{\mu_i}{\check{h}_{i+1}} \left( \frac{\varphi_{i+1} - \varphi_i}{h_{i+1}} - \frac{\varphi_i - \varphi_{i-1}}{h_i} \right);$$

where

$$\check{h}_i = 0.5(h_i + h_{i-1});$$

$$\mu_i = \varepsilon \left( h_i h_i - 3c\tau \check{h}_i + 2c^2 \tau^2 \right);$$

$$\varepsilon = 0.08$$

Extension to the vector case: (F=flux vector, U=vector variable, A=quasi-linear matrix, W=local Riemann invariants, L= left eigen matrix )

$$F^* = A \otimes U^*; W^* = L \otimes U^*; \Lambda = L \otimes A \otimes L^{-1}$$

$$L \otimes F^* = (L \otimes A \otimes L^{-1}) \otimes L \otimes U^* = \Lambda \otimes W^*$$

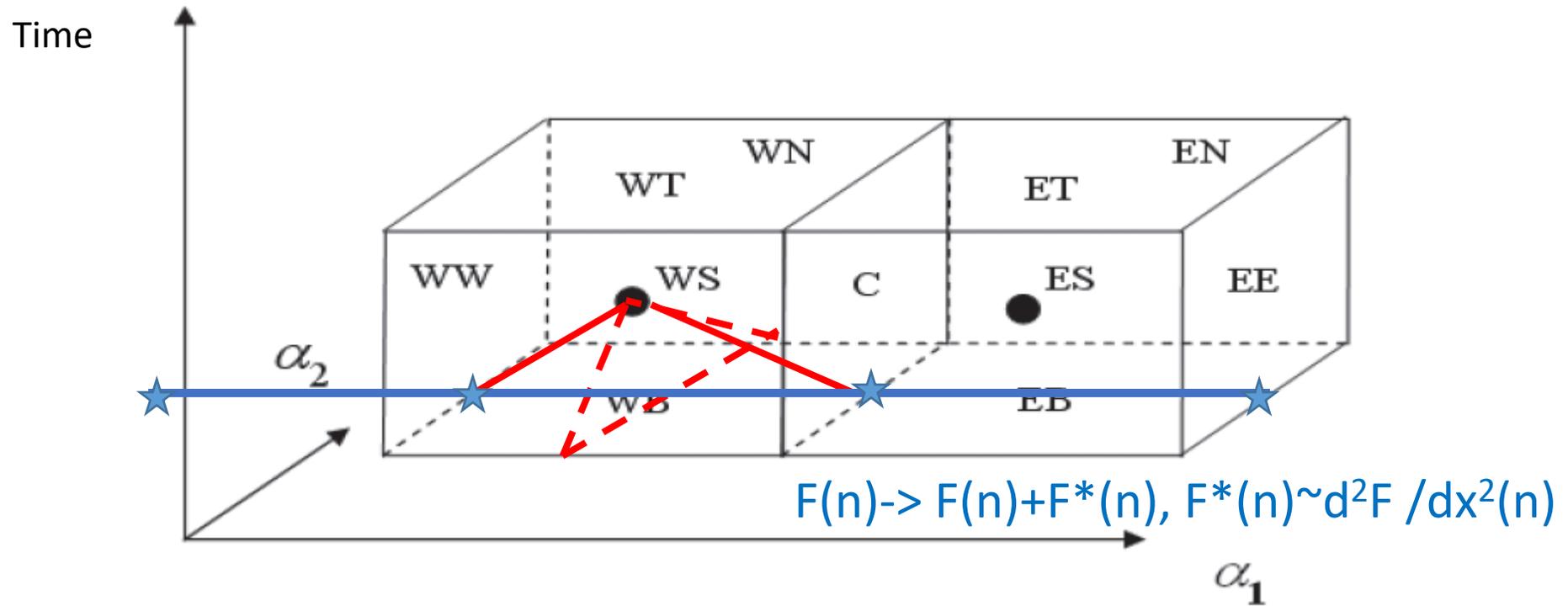
$$W^* \sim d^2 W / dx^2; L \otimes F^* \sim \Lambda \otimes d^2 W / dx^2$$

$$F^* = L^{-1} \otimes (L \otimes F^*) \sim L^{-1} \otimes (L \otimes A \otimes L^{-1}) \otimes d^2 W / dx^2$$

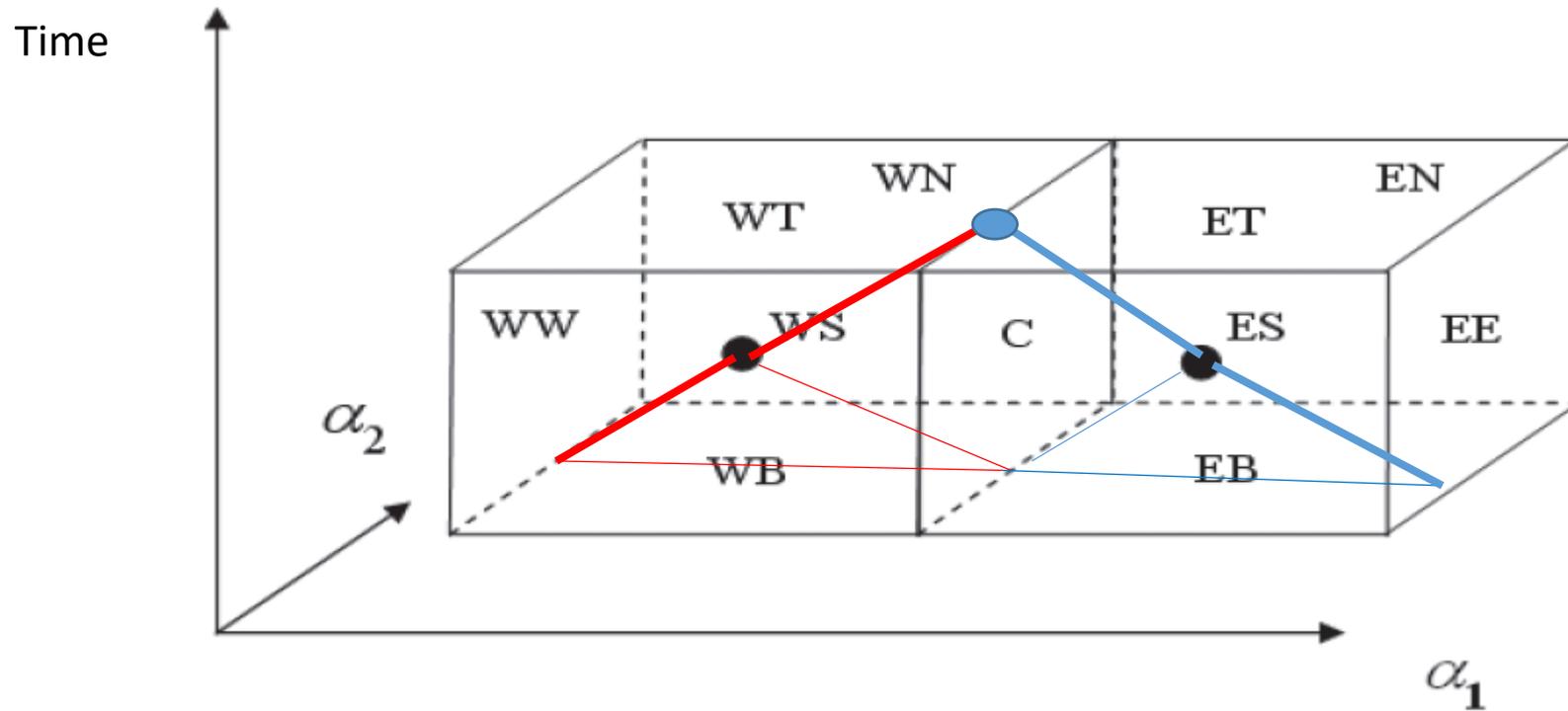
$$F^* \sim A \otimes d^2 U / dx^2 = d^2 F / dx^2$$

## 2D CABARET: predictor step

cell centre + cell face variables  $\rightarrow$  mid time level variables

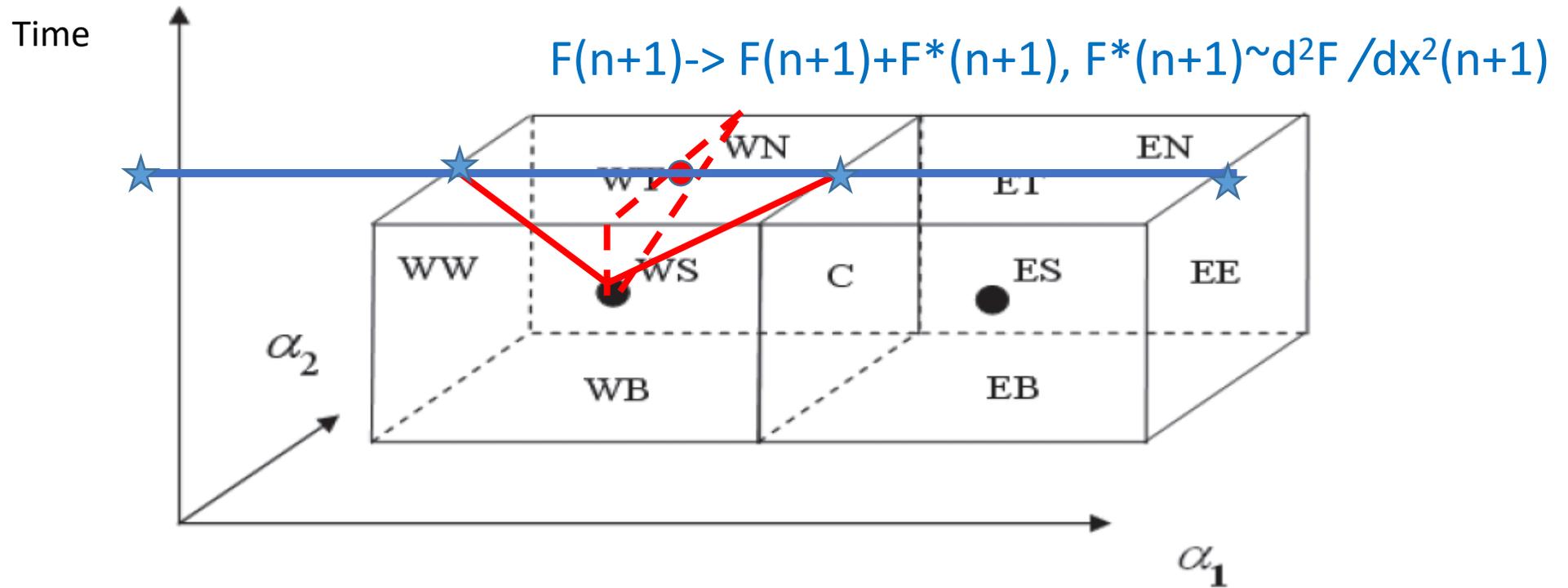


CABARET characteristic decomposition step,  
extrapolation and solution of the Riemann  
problem at the cell face in each grid direction  
=the same!



## 2D CABARET: corrector step

cell centre + cell face variables  $\rightarrow$  new time level variables



# Flux correction: Full Flux Corrected CABARET

- The P-CAB scheme is not monotone and allows non-physical oscillations.
- This is mitigated by applying a correction based on the maximum principle at the extrapolation stage

$$\text{if}(u_{i+\frac{1}{2}}^{n+1} > M)u_{i+\frac{1}{2}}^{n+1} = M \quad M = \max(u_{i+\frac{1}{2}}^n, u_i^n, u_{i-\frac{1}{2}}^n)$$

$$\text{if}(u_{i+\frac{1}{2}}^{n+1} < m)u_{i+\frac{1}{2}}^{n+1} = m \quad m = \min(u_{i+\frac{1}{2}}^n, u_i^n, u_{i-\frac{1}{2}}^n)$$

# Flux correction: Relax Flux Corrected CABARET

- The maximum principle limits are relaxed using a tunable relaxation parameter (epsilon).

$$M_{new} = (1 + \epsilon)M$$

$$m_{new} = (1 - \epsilon)m$$

$$if(M < 0)M_{new} = (1 - \epsilon)M$$

$$if(m < 0)m_{new} = (1 + \epsilon)m$$

# Flux correction: Modified Relax Flux Corrected CABARET

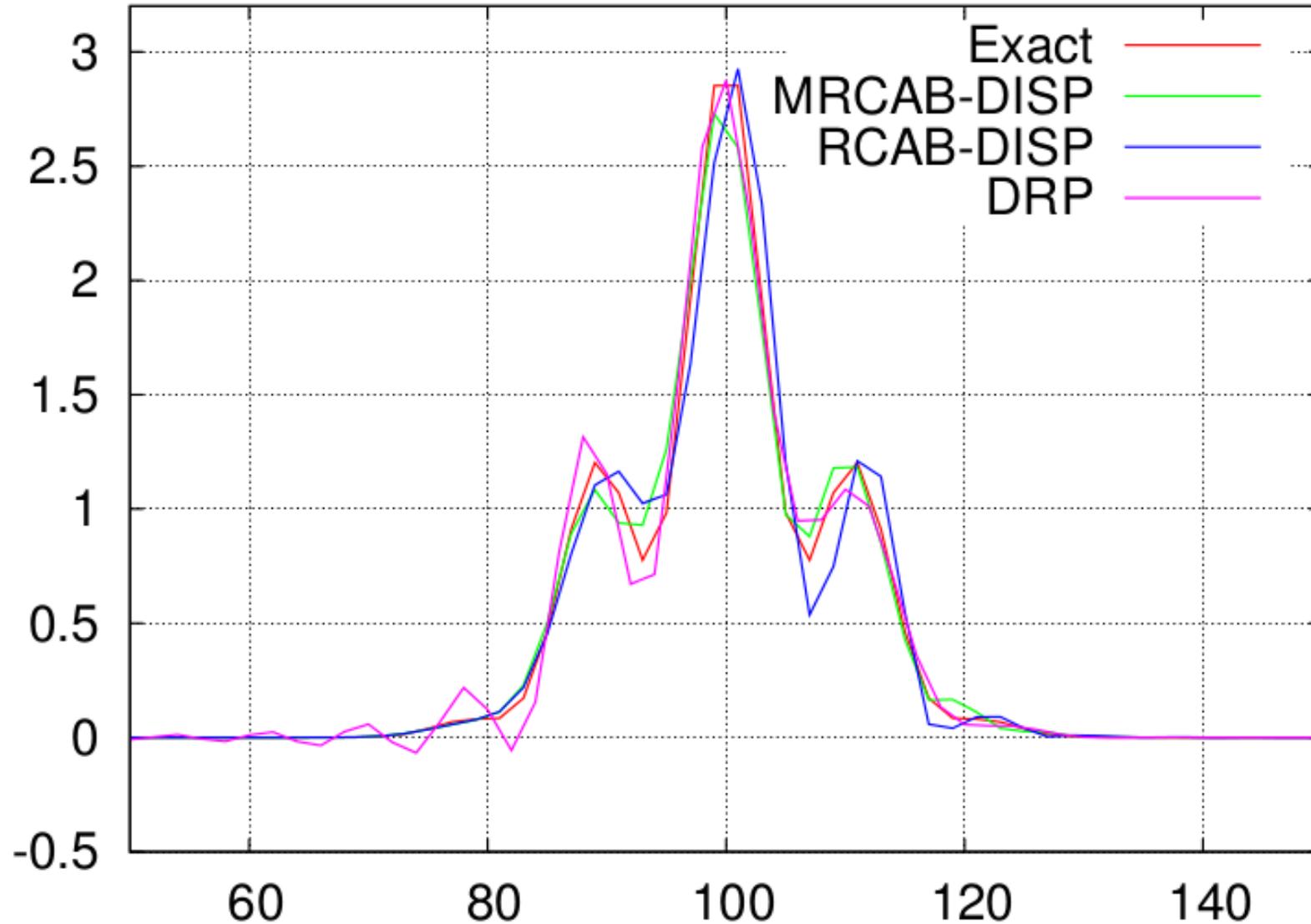
- The maximum principle limits are relaxed using a tunable relaxation parameter (epsilon) and also the local maxima and minima.

$$\delta_{rx} = \frac{(M - m)}{2}$$

$$M_{new} = M + \epsilon \delta_{rx}$$

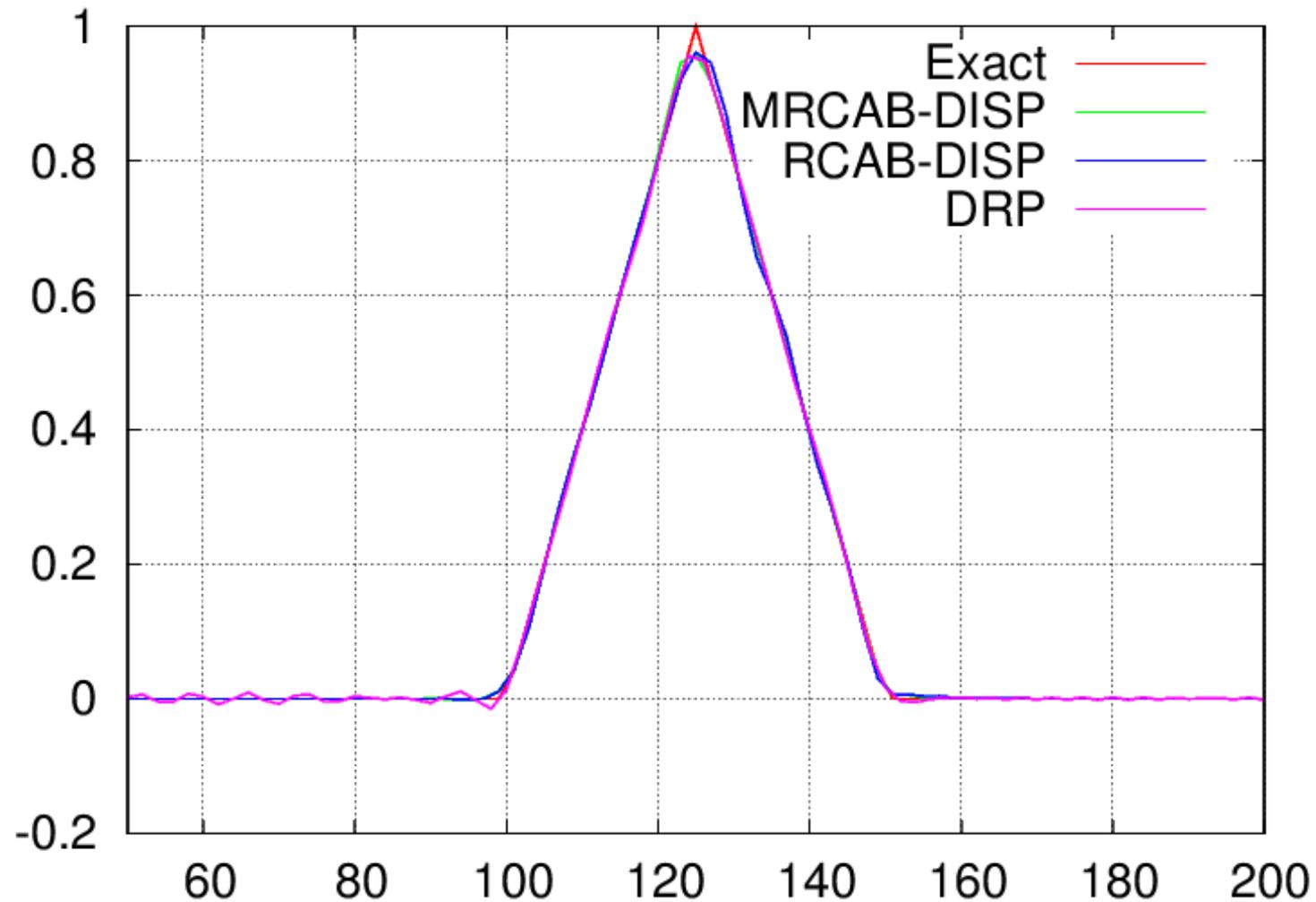
$$m_{new} = m - \epsilon \delta_{rx}$$

# Linear advection test#1: R-CAB,MR-CAB and DRP



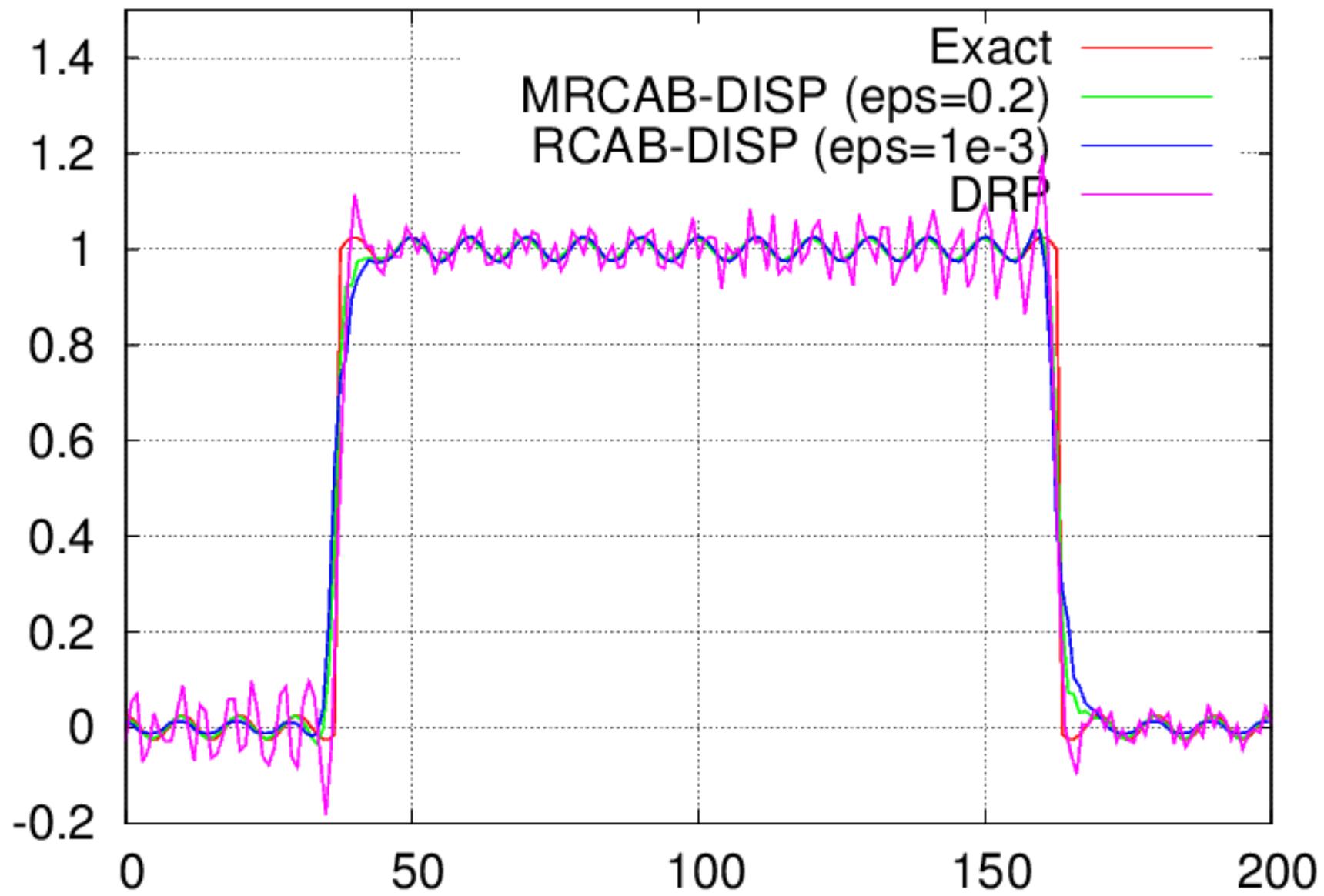
DRP is the winner for most grids

## Linear advection test#2: R-CAB,MR-CAB and DRP

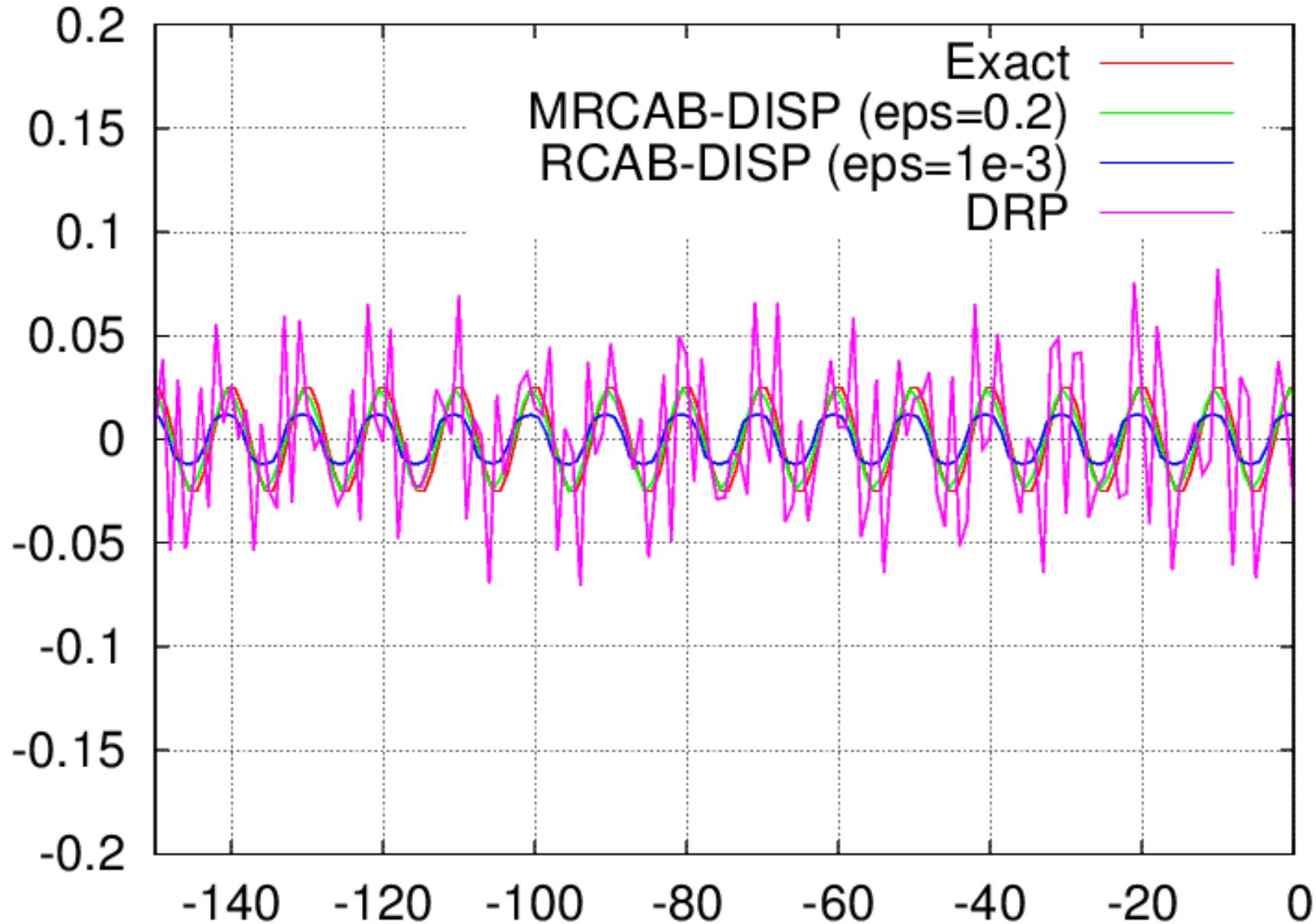


MRCAB-DISP  
~ DRP for all  
grids

# Linear advection test#3: R-CAB, MR-CAB and DRP schemes



## Test case#3 – floor



MRCAB-DISP  
is the winner  
for all grids

# Analytical test problem for CAA methods

Physics: acoustic wave propagation in a channel with periodic side walls and uniform flow at  $M=0.5$ ;

Model: polytropic & isothermal inviscid gas flow equations in 1D and 2D

## Linear wave equation

$$\frac{\partial^2}{\partial t^2} \rho' - c_0^2 \frac{\partial^2}{\partial x \partial x} \rho' - c_0^2 \frac{\partial^2}{\partial y \partial y} \rho' = 0$$

## Linear wave equation with constant flow

$$\left(\frac{\partial}{\partial t} - U \frac{\partial}{\partial x}\right) \left(\frac{\partial}{\partial t} - U \frac{\partial}{\partial x}\right) \rho' - c_0^2 \frac{\partial^2}{\partial x \partial x} \rho' - c_0^2 \frac{\partial^2}{\partial y \partial y} \rho' = 0$$

Travelling wave solution  $\rho' \sim \rho'_k e^{i(\omega t - k_x \cdot x - k_y \cdot y)}$

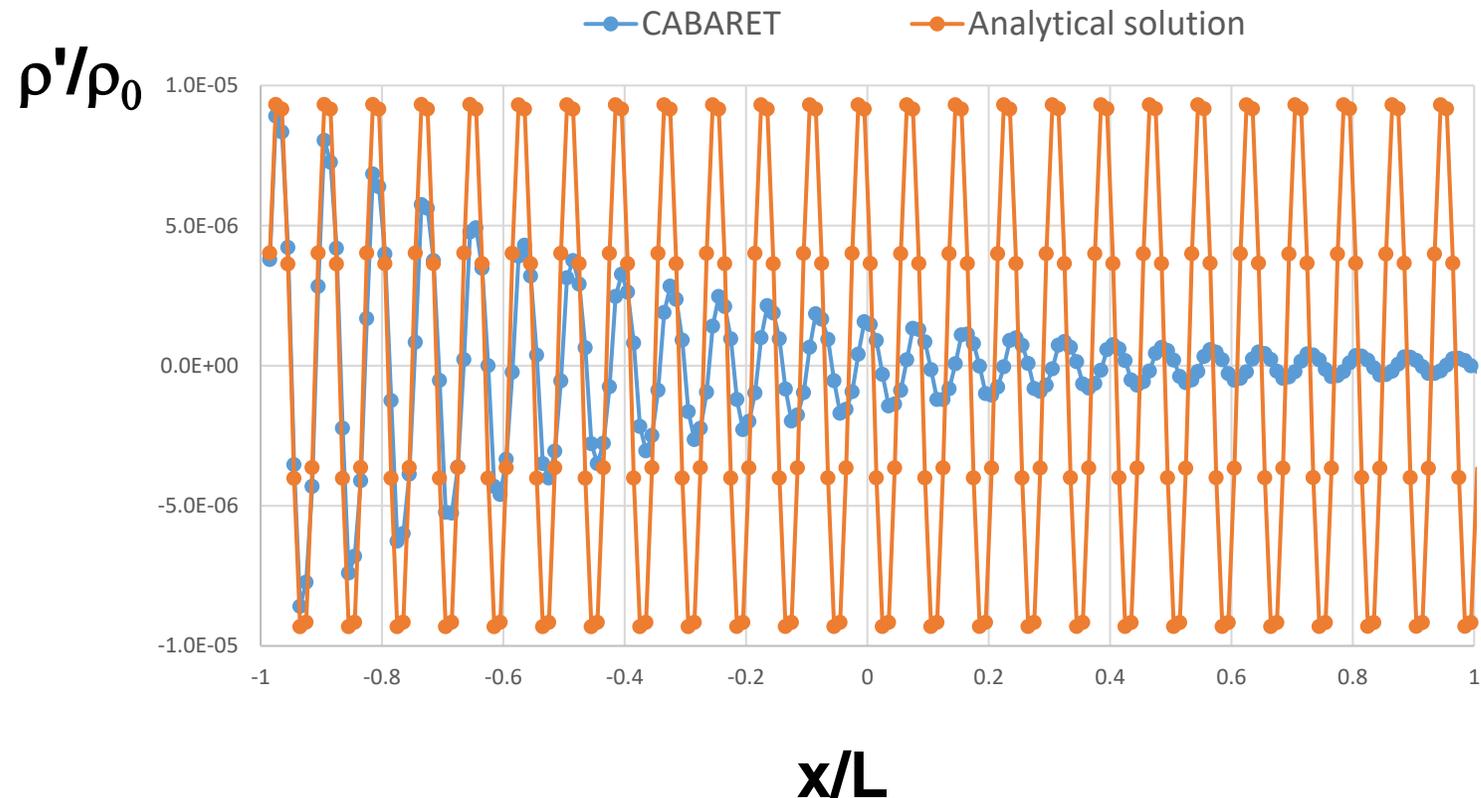
## Dispersion relations

$$\omega^2 - c_0^2 k_x^2 - c_0^2 k_y^2 = 0$$

$$\omega^2 + 2k_x U + k_x^2 U^2 - c_0^2 k_x^2 - c_0^2 k_y^2 = 0$$

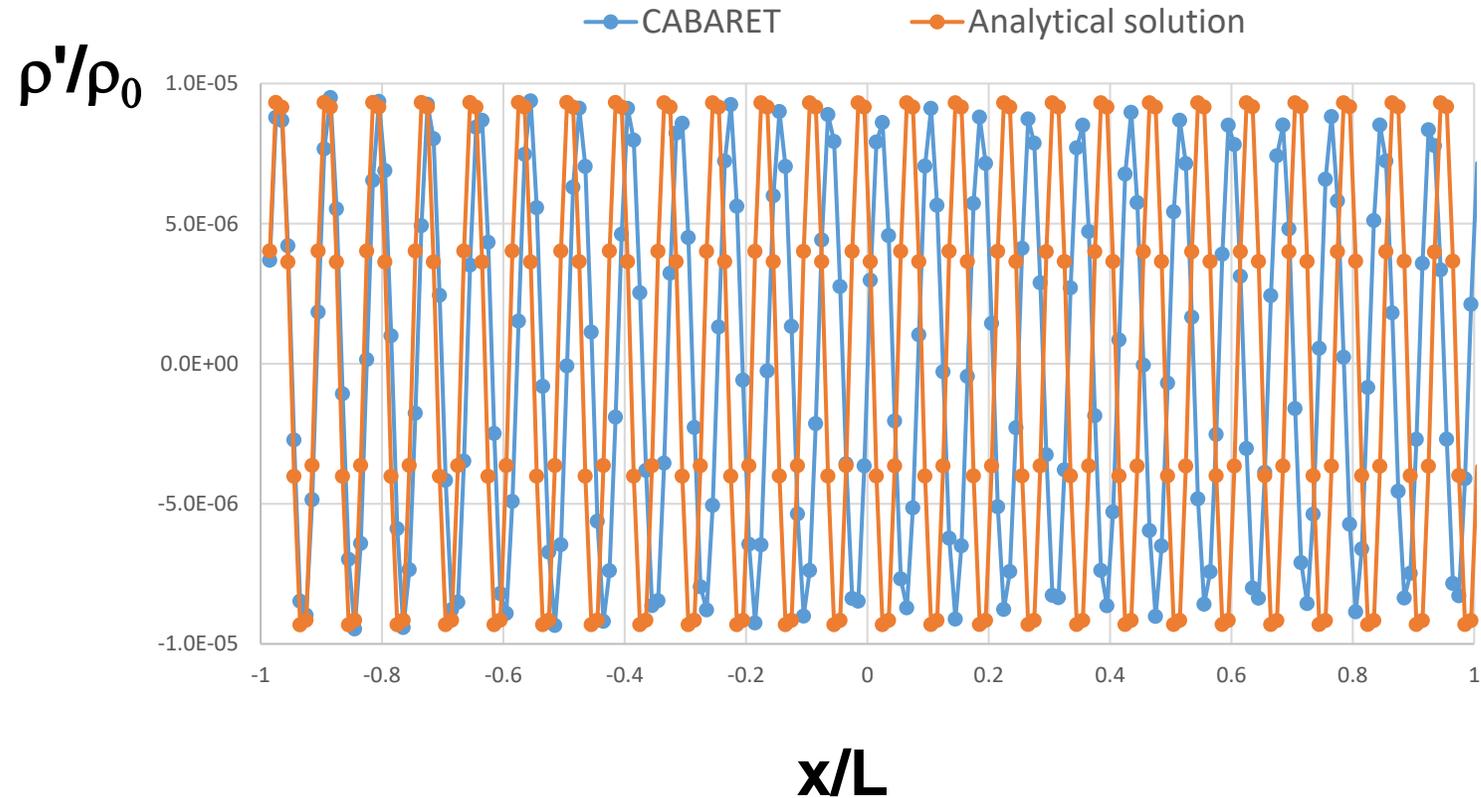
# Planar wave (1D), 8ppw, CFL=0.1

## Full flux correction, original CABARET



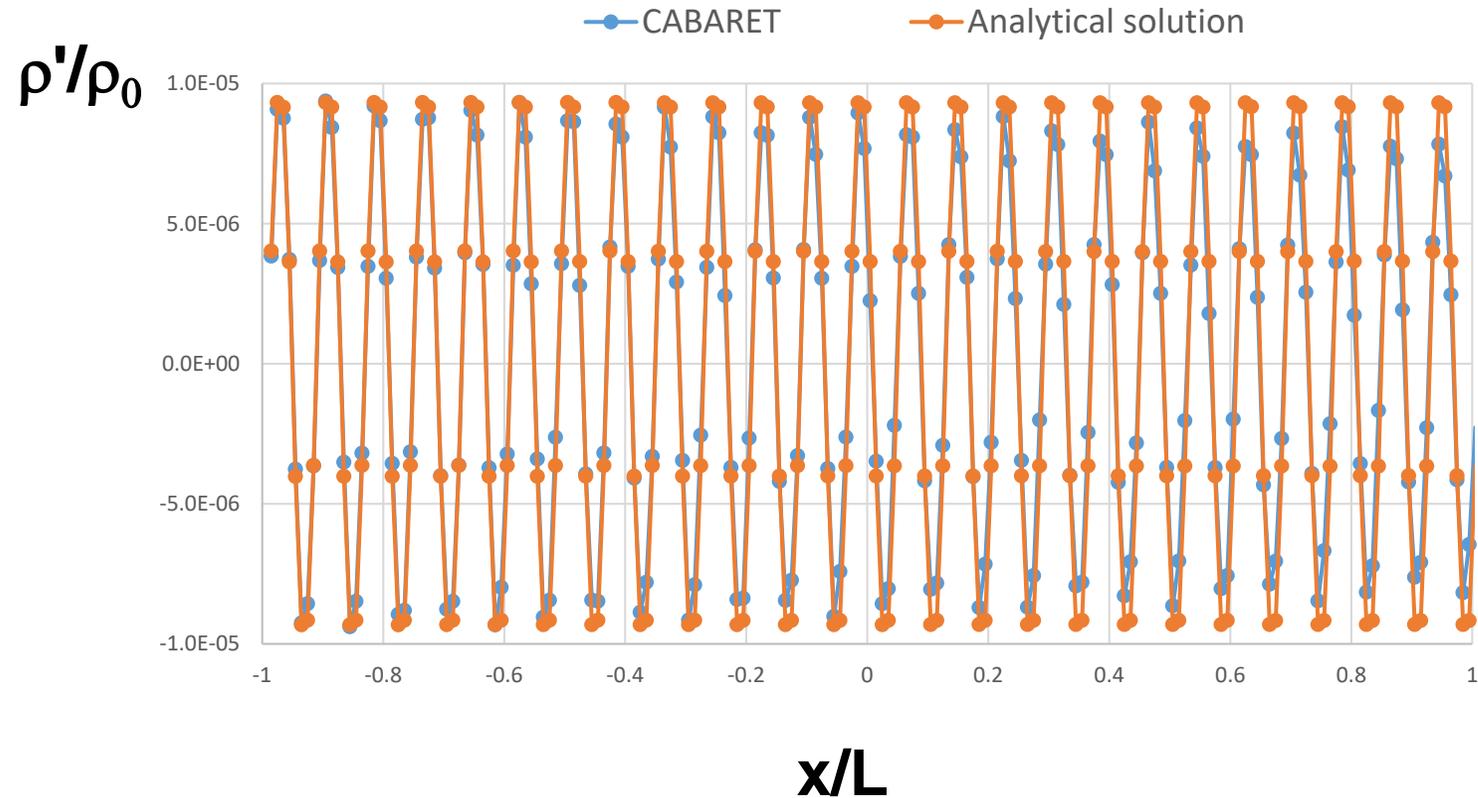
Planar wave (1D), 8ppw, CFL=0.1

New relaxed flux correction ( $\varepsilon=0.2$ ), original CABARET

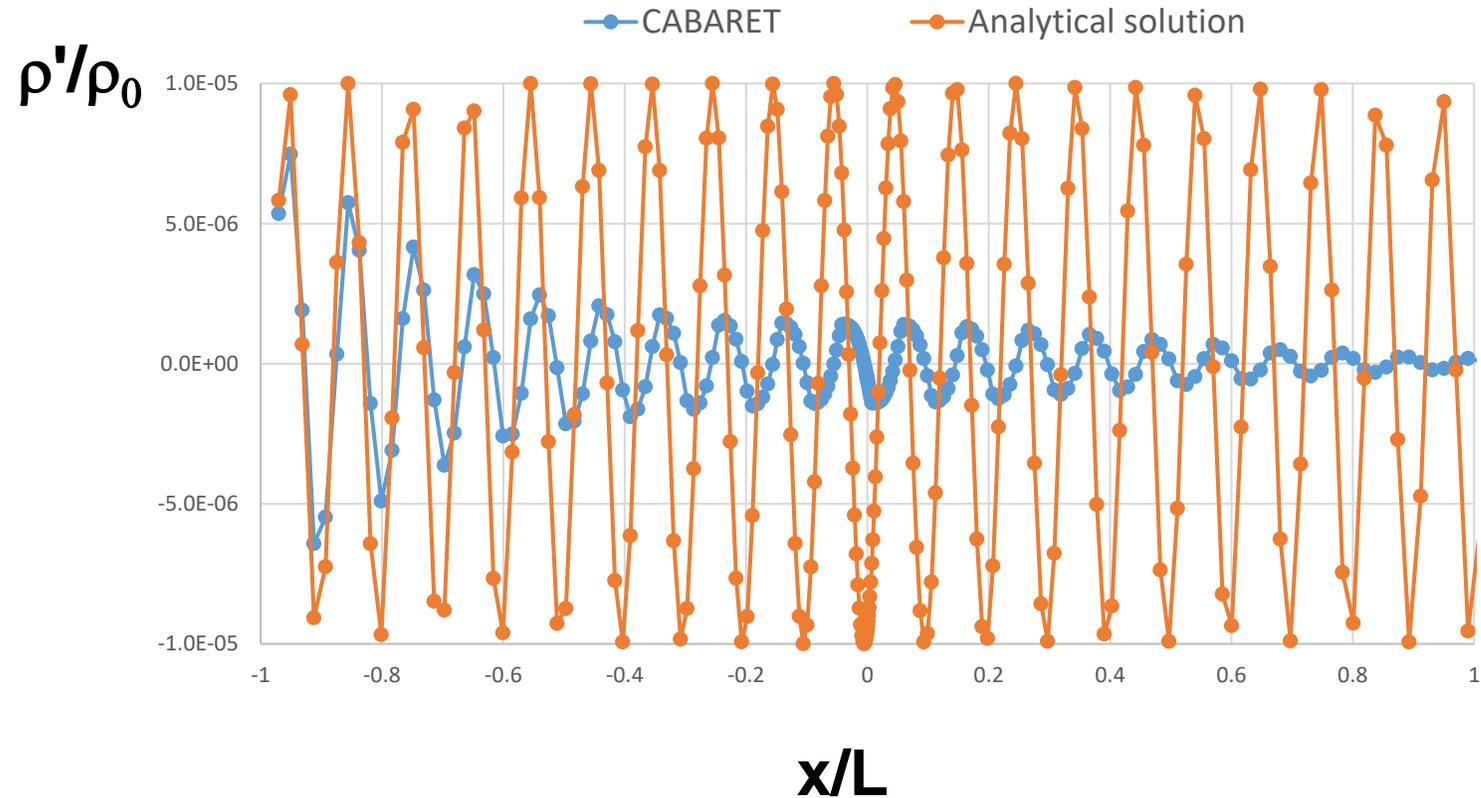


Planar wave (1D), 8ppw, CFL=0.1

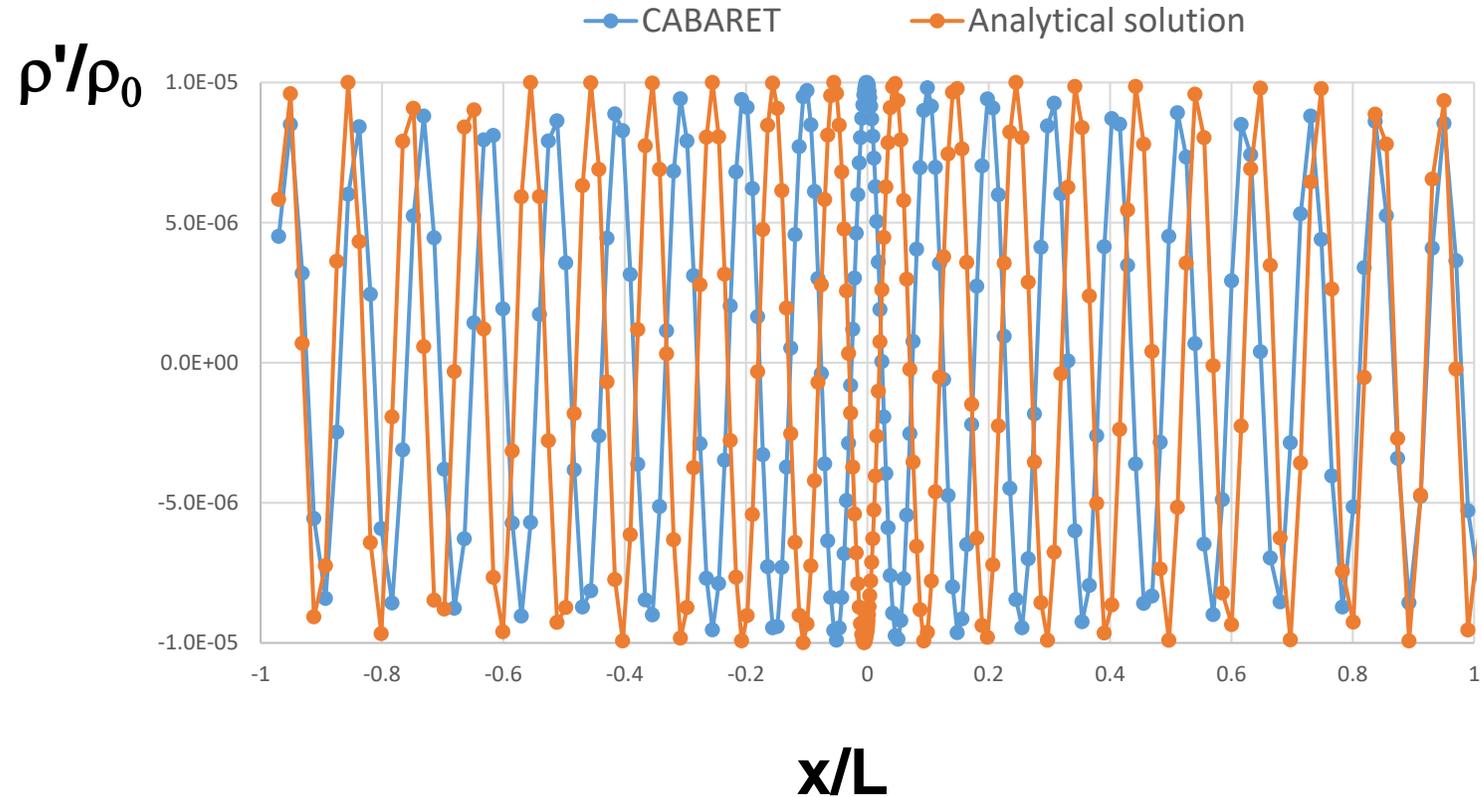
New relaxed flux correction ( $\varepsilon=0.2$ ), dispersion improved CABARET



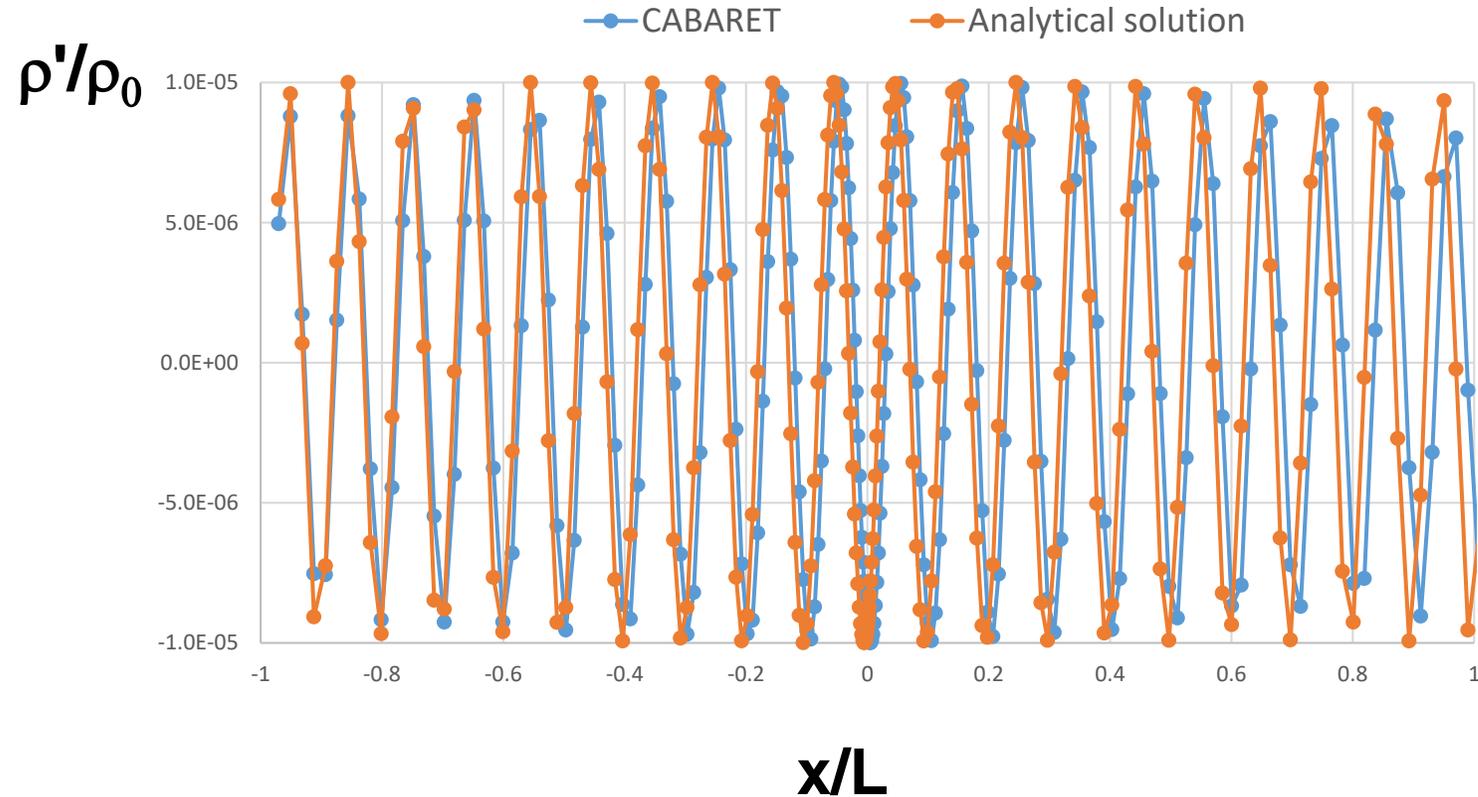
# Planar wave (1D), non-uniform grid, $CFL_{\max}=0.8$ Full flux correction, original CABARET



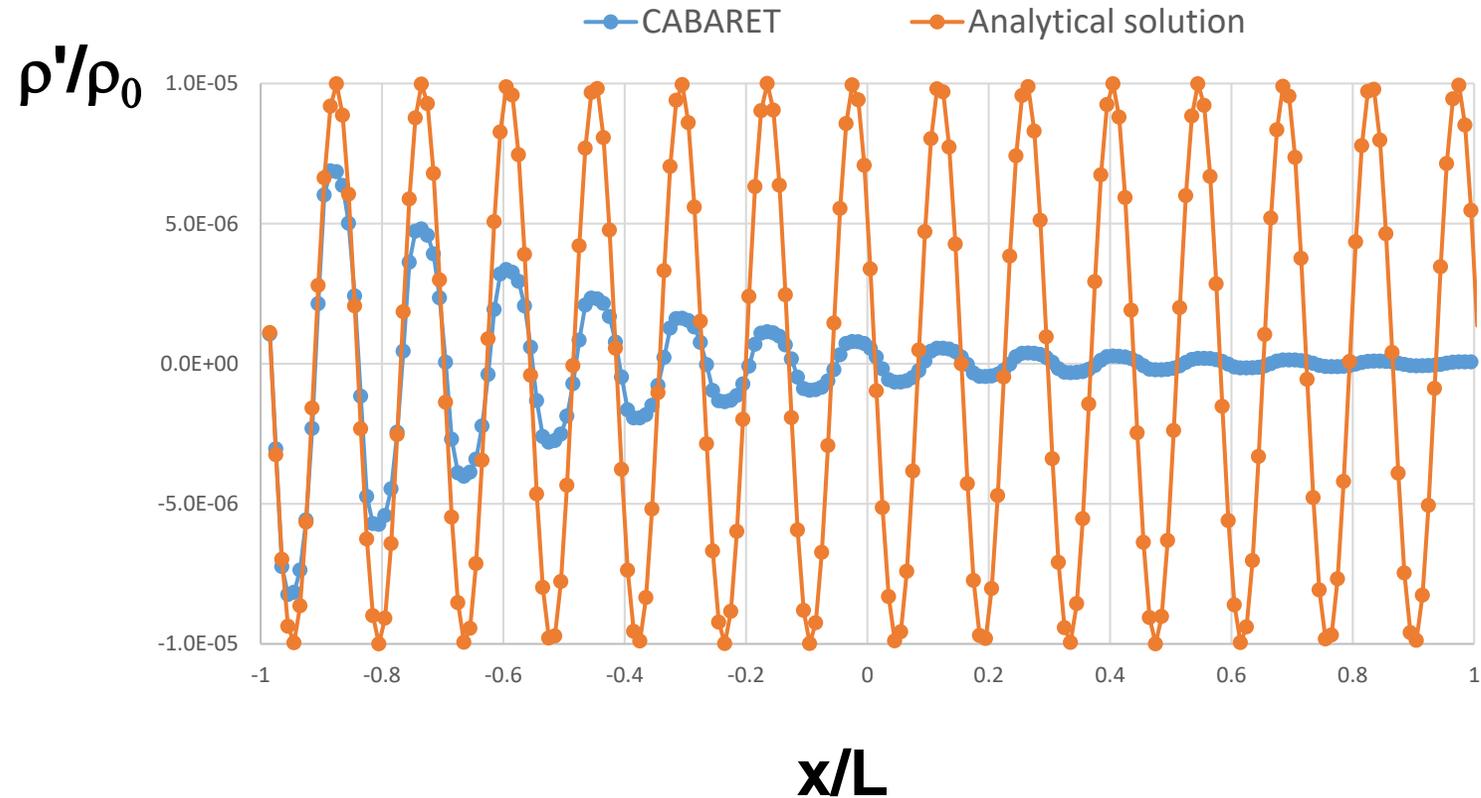
Planar wave (1D), non-uniform grid,  $CFL_{\max}=0.8$   
New relaxed flux correction ( $\varepsilon=0.2$ ), original CABARET



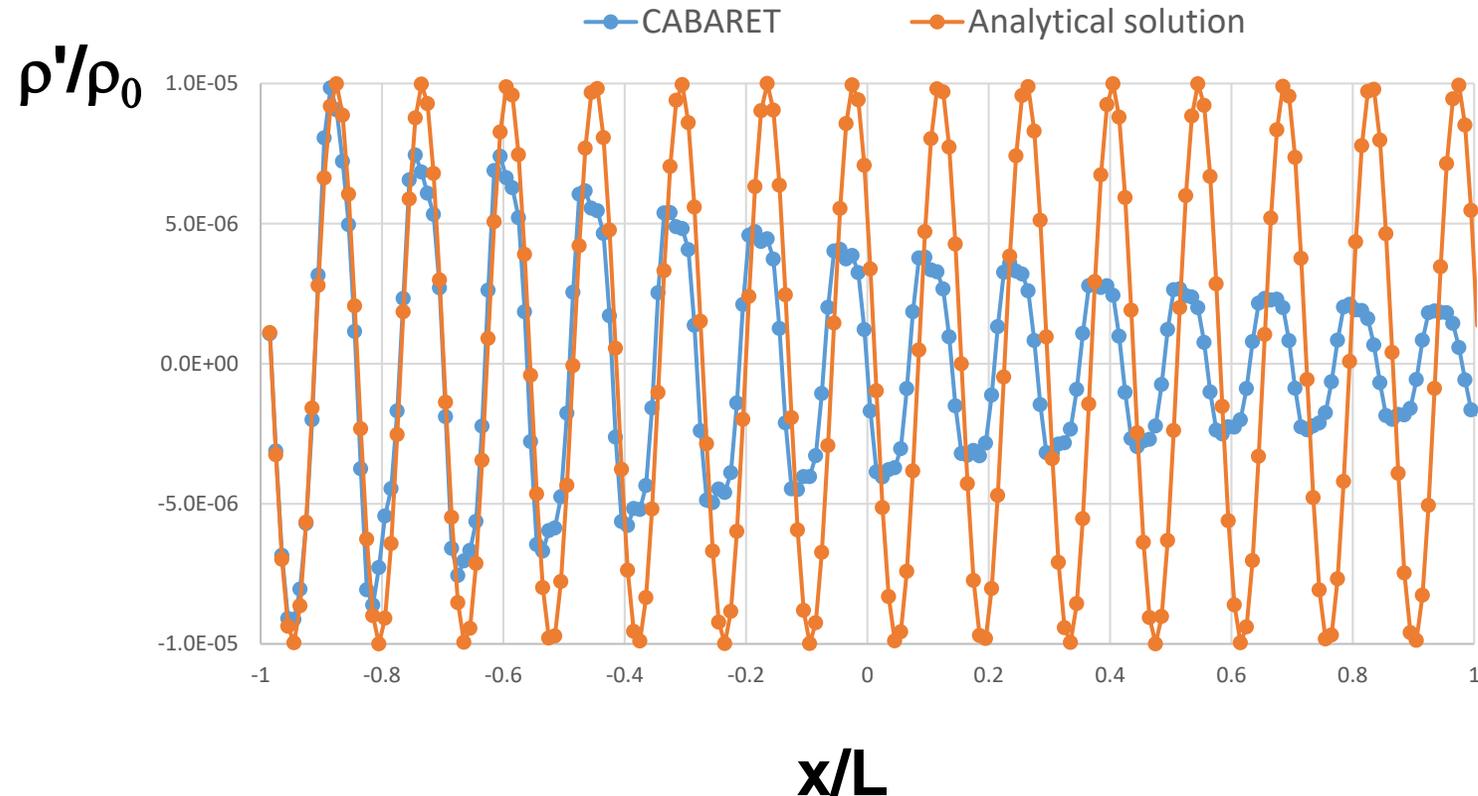
Planar wave (1D), non-uniform grid,  $CFL_{\max}=0.8$   
New relaxed flux correction ( $\varepsilon=0.2$ ), dispersion  
improved CABARET



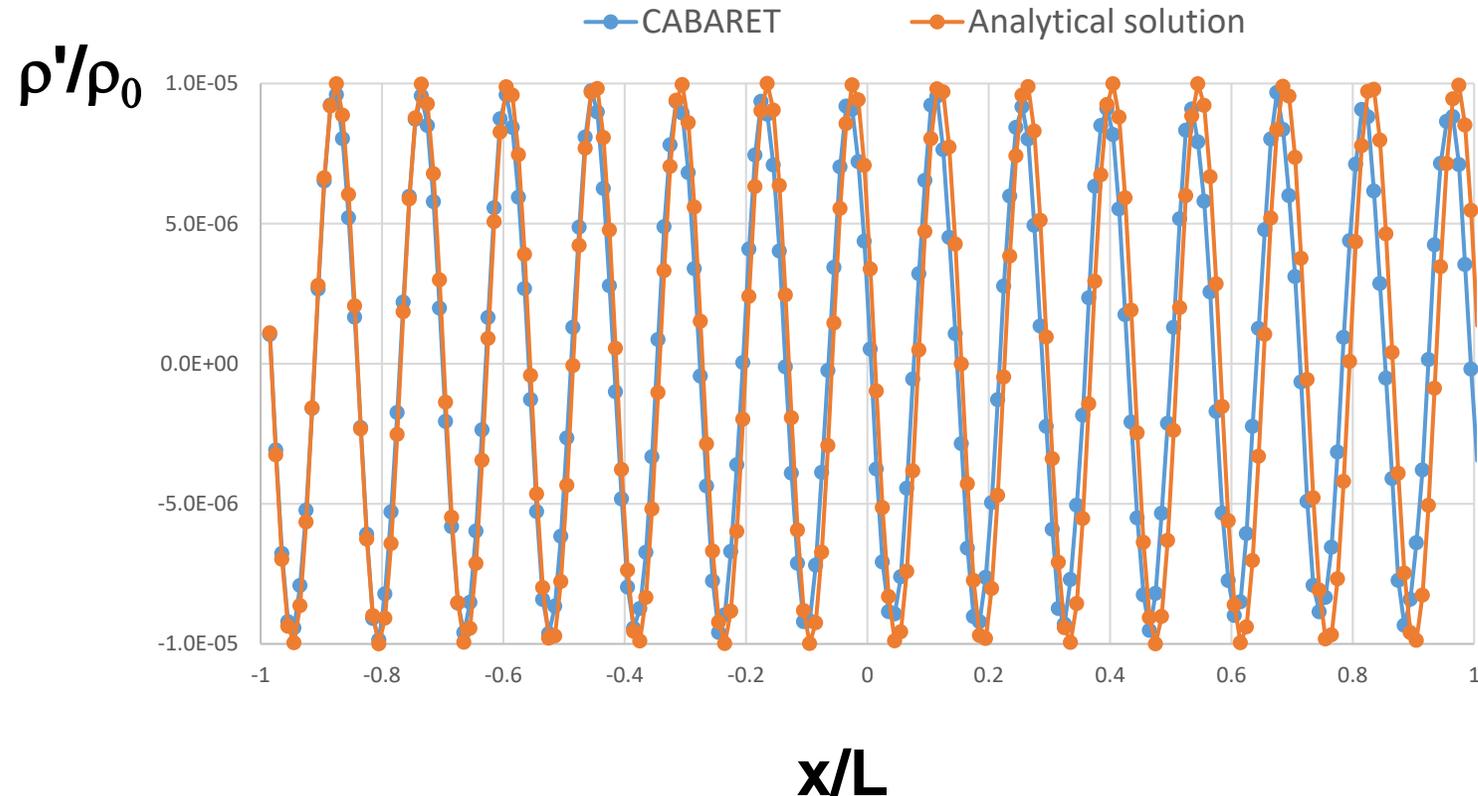
2D wave, CFL=0.1,  $ppw_{x,y}=\{15,9\}$ ,  $\text{atan}(k_y/k_x)=58^\circ$   
Full flux correction, original CABARET



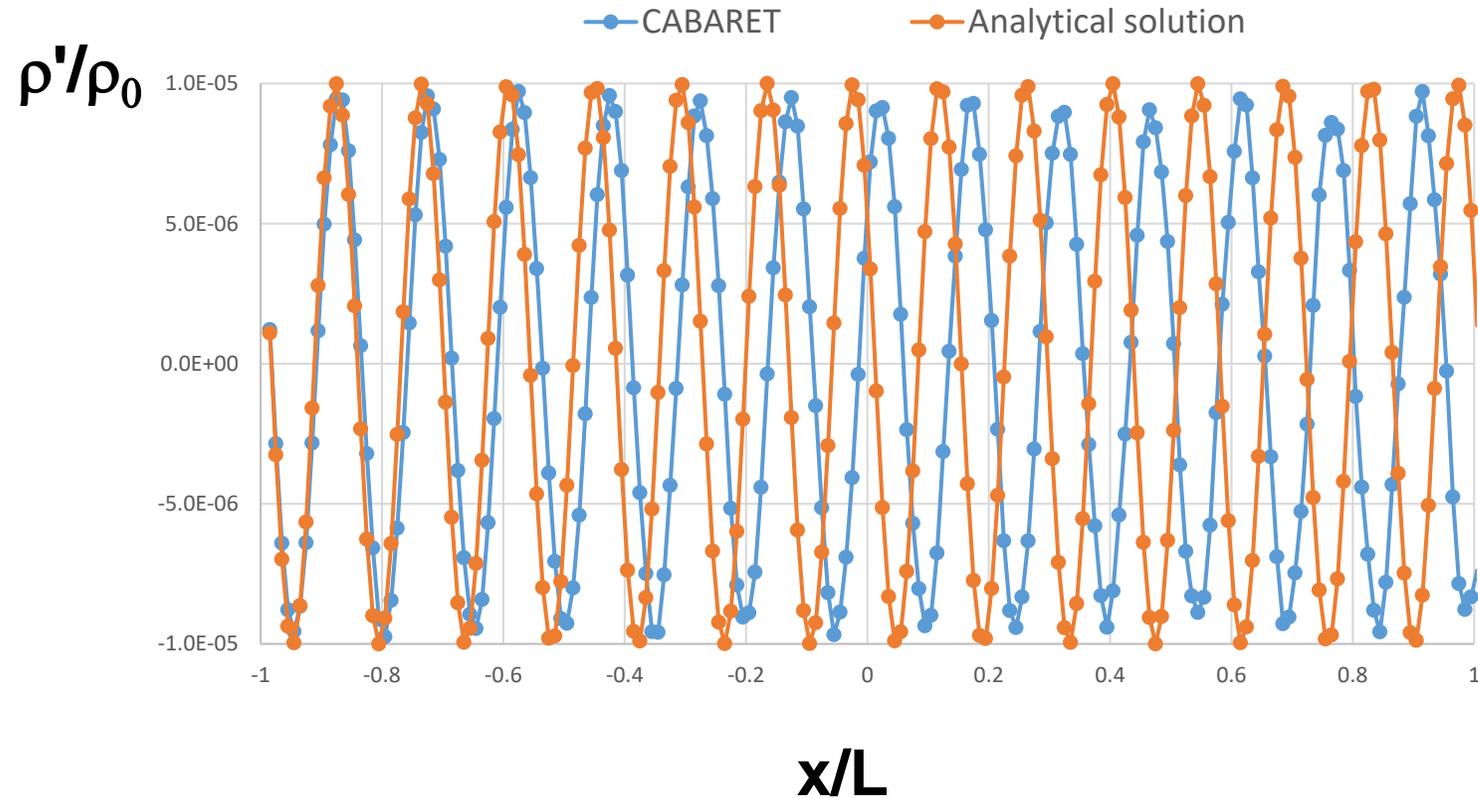
2D wave, CFL=0.1, ppw<sub>x,y</sub>={15,9}, atan(k<sub>y</sub>/k<sub>x</sub>)=58°  
New relaxed flux correction (ε=0.2), dispersion  
improved CABARET



2D wave, CFL=0.1, ppw<sub>x,y</sub>={15,9}, atan(k<sub>y</sub>/k<sub>x</sub>)=58°  
New relaxed flux correction (ε=0.4), dispersion  
improved CABARET

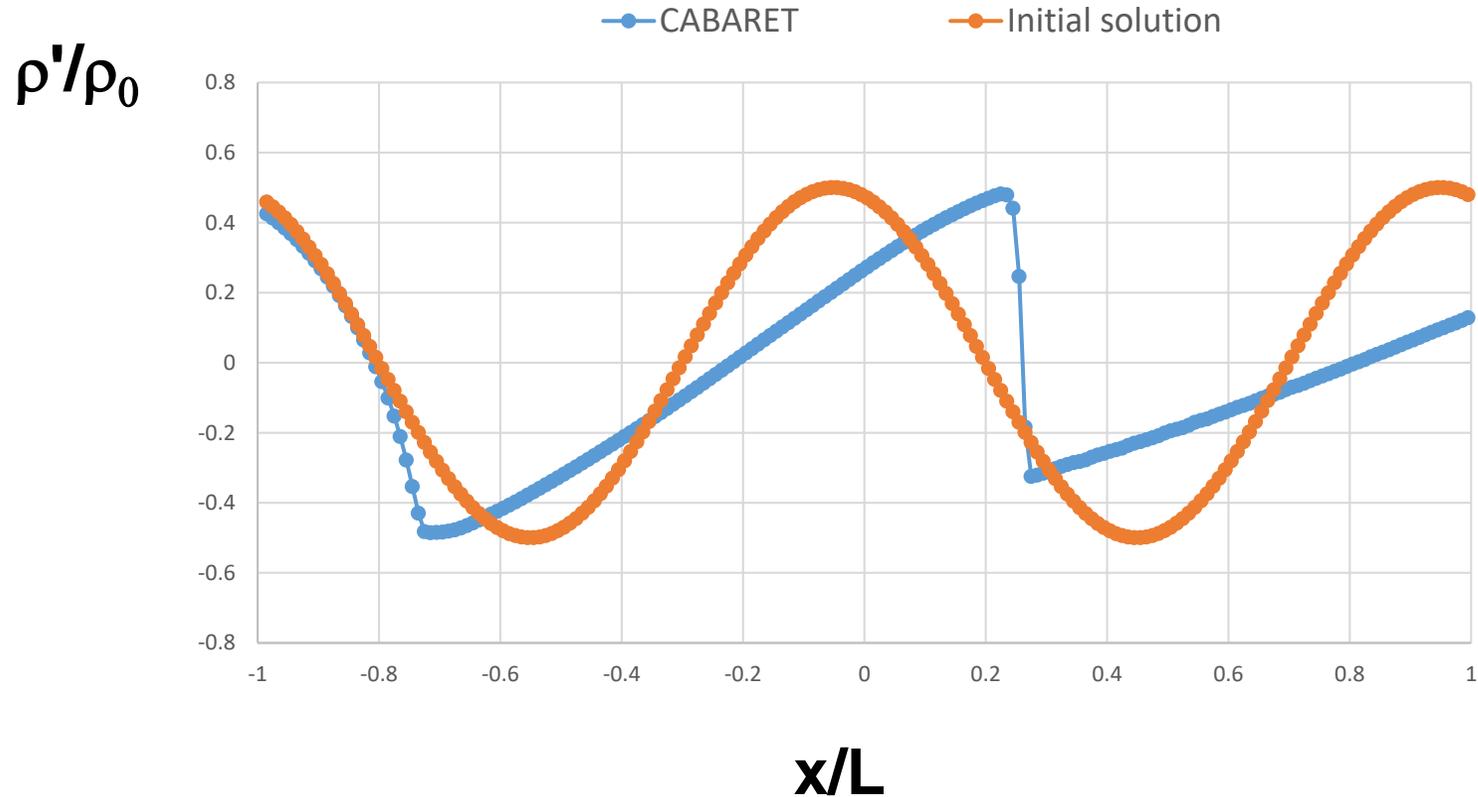


2D wave, CFL=0.1, ppw<sub>x,y</sub>={15,9}, atan(k<sub>y</sub>/k<sub>x</sub>)=58°  
New relaxed flux correction (ε=0.4), standard  
CABARET

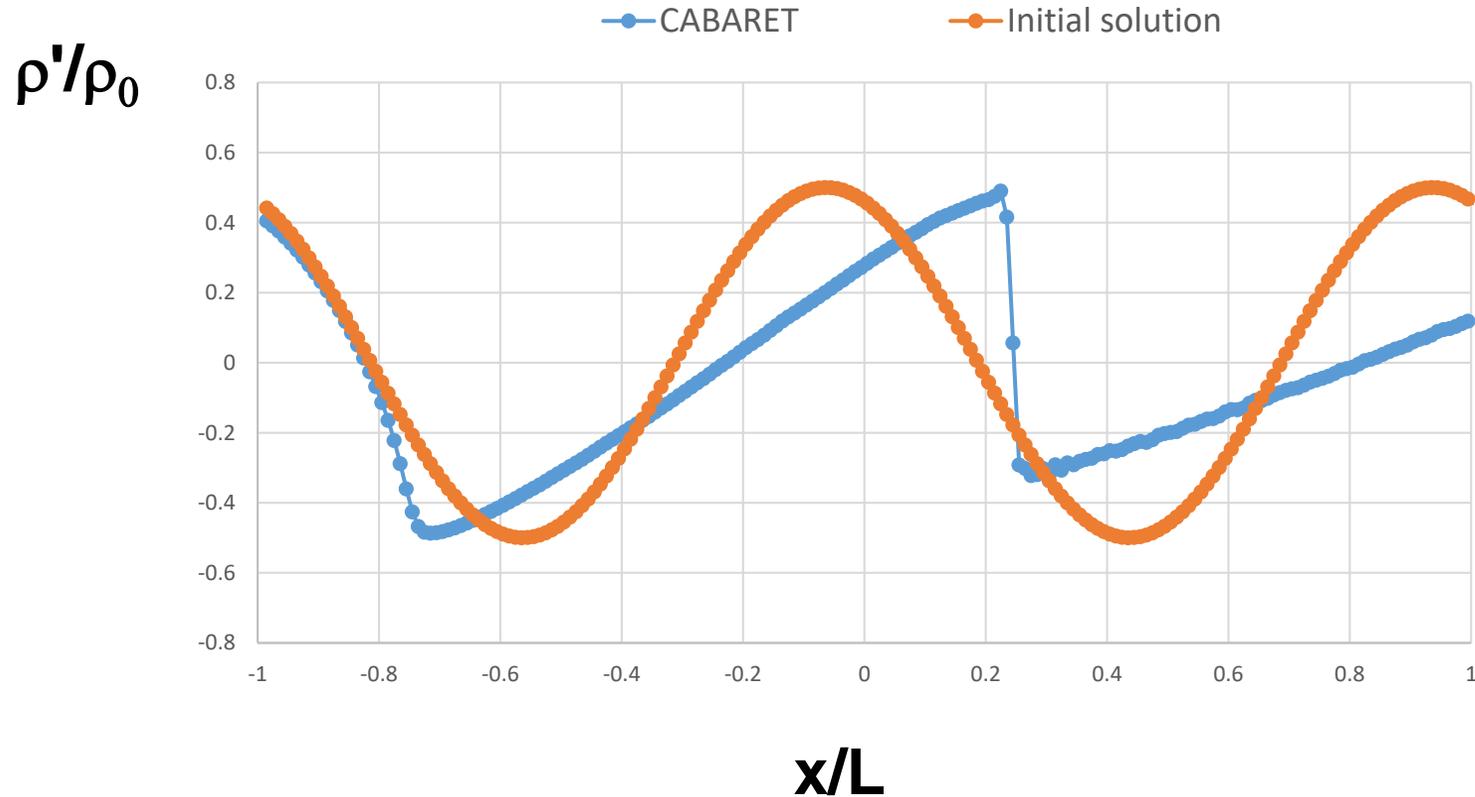


# Nonlinear planar wave (1D), 100ppw, CFL=0.1

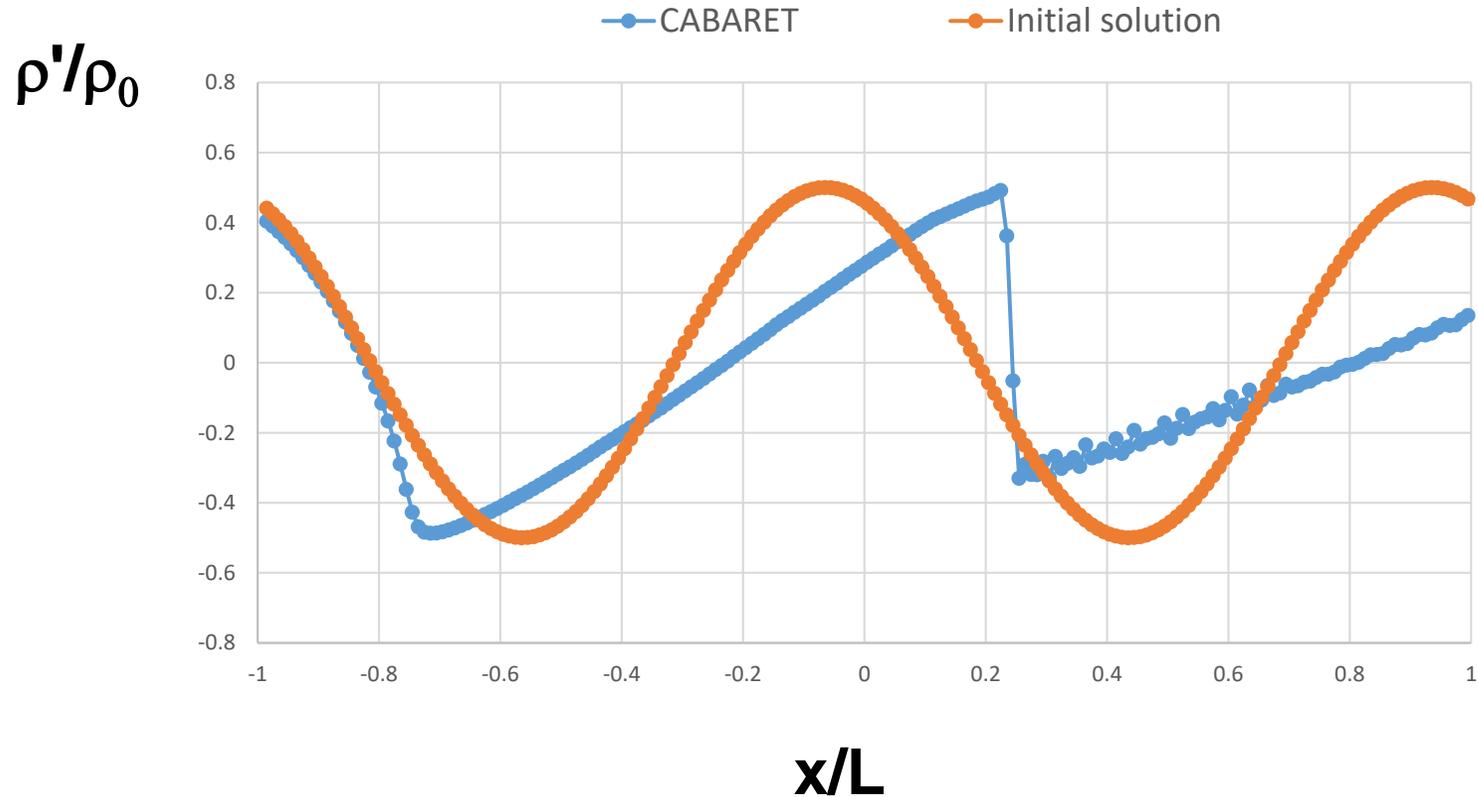
## Full flux correction, original CABARET



Nonlinear planar wave (1D), 100ppw, CFL=0.1  
New relaxed flux correction ( $\varepsilon=0.2$ ), dispersion  
improved CABARET

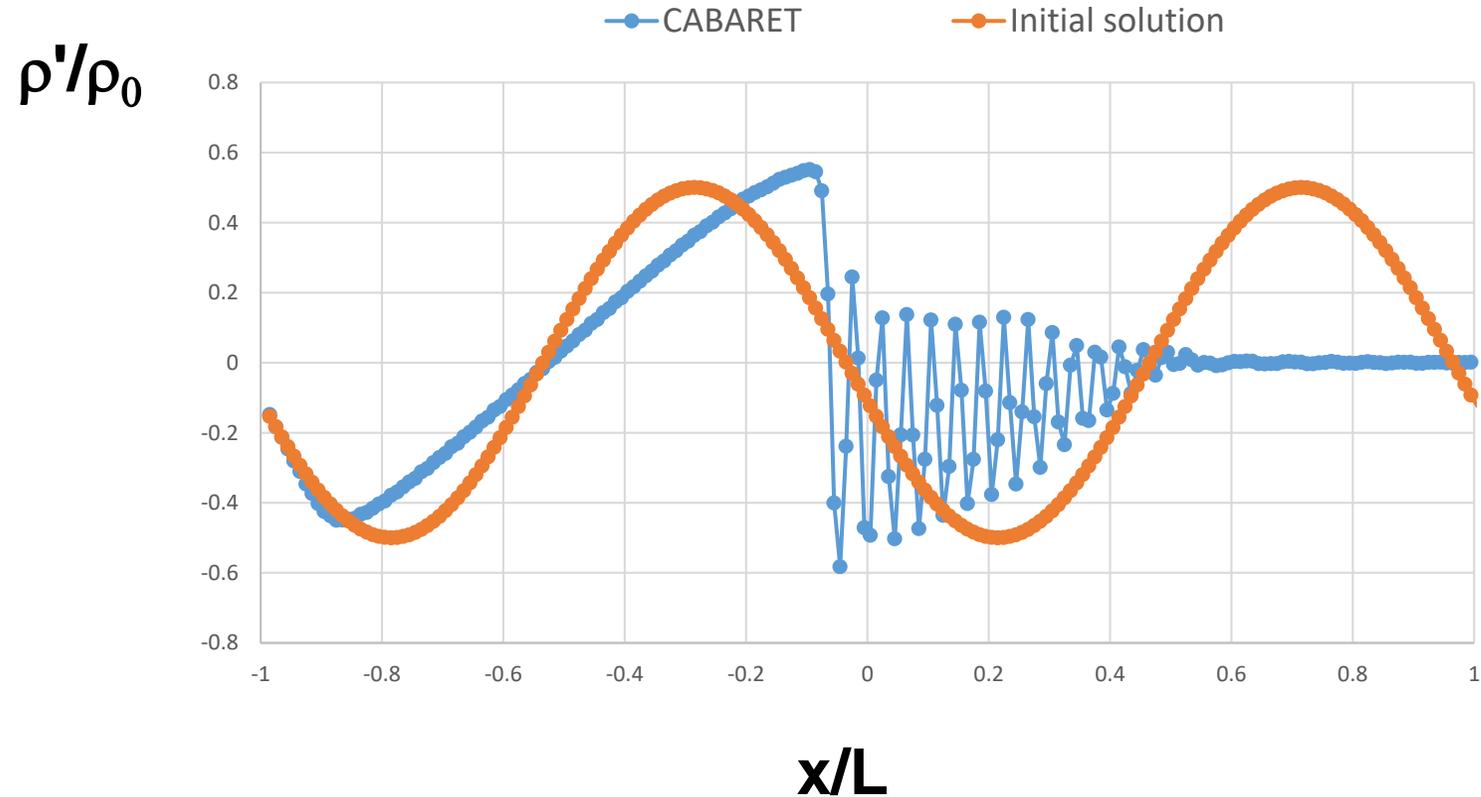


Nonlinear planar wave (1D), 100ppw, CFL=0.1  
New relaxed flux correction ( $\varepsilon=0.4$ ), dispersion  
improved CABARET



# Nonlinear planar wave (1D), 100ppw, CFL=0.1

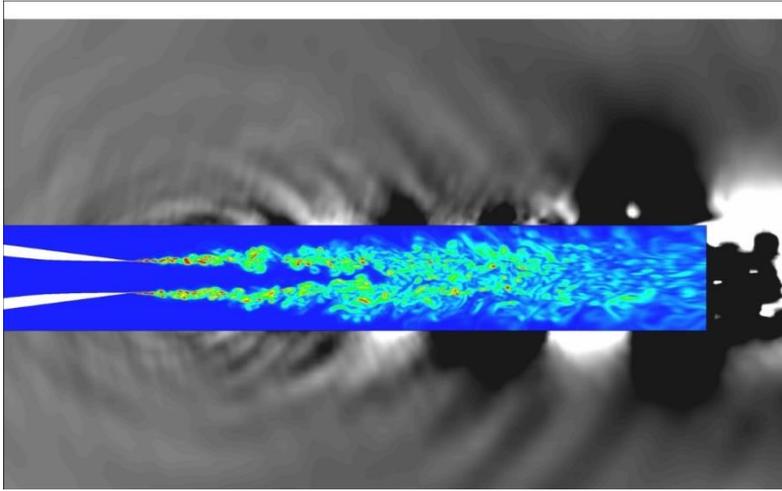
## No flux correction, 'pure' CABARET



# Conclusion

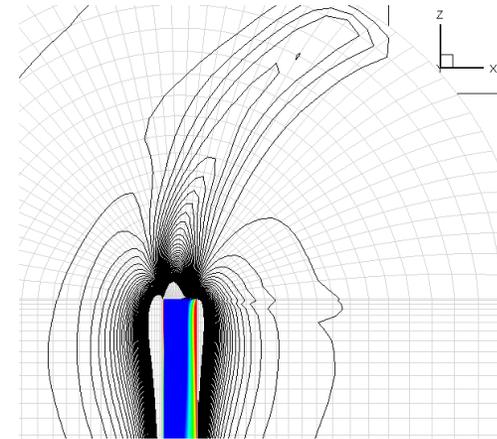
- Dispersion improved CABARET of Goloviznin and Samarskii (1998) is extended to the compact two-time-level predictor-corrector form
- The compact formulation allows further generalisations to include non-linear flux correction (including a new relaxed correction method) and extensions to multiple dimensions with gas dynamics.
- Numerical solutions for 1d linear advection, 1d and 2d linear acoustic propagation, and 1d non-linear z-wave propagation are considered
- A considerable increase in accuracy over the standard CABARET (not to mention DRP;) is reported. In combination with the new relaxed correction algorithm within a small range of calibration parameter  $\varepsilon=0.2-0.4$ , superior results are reported for both linear and nonlinear problems.
- Further work will be towards implementing the improved CABARET in the current 3D GPU CABARET NS solver

# Aeroacoustics of engineering flows: challenges



Faranosov et al. 2013

*Vorticity and acoustic pressure field of a turbulent jet*



Morgans et a. 2005

*Acoustic pressure contours of a high-speed helicopter blade*

Aerodynamic scales  $\ll$  Acoustic scales

$$\delta \text{ b.layer}/D \sim 0.01-0.001$$

$$L/D \sim 100-1000$$

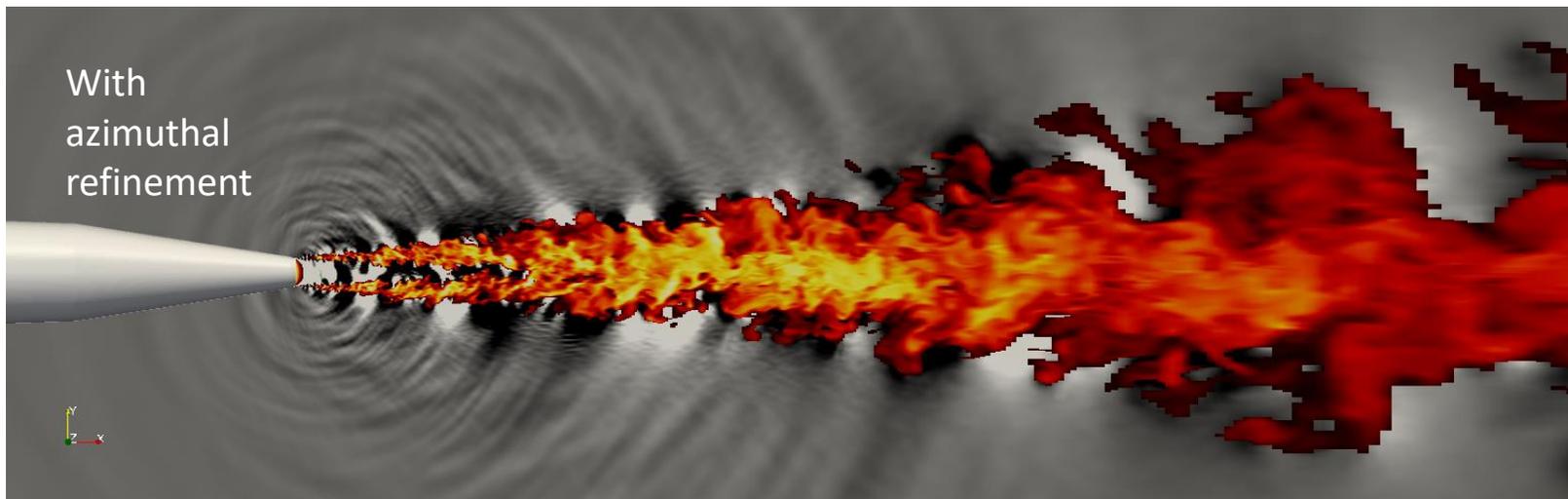
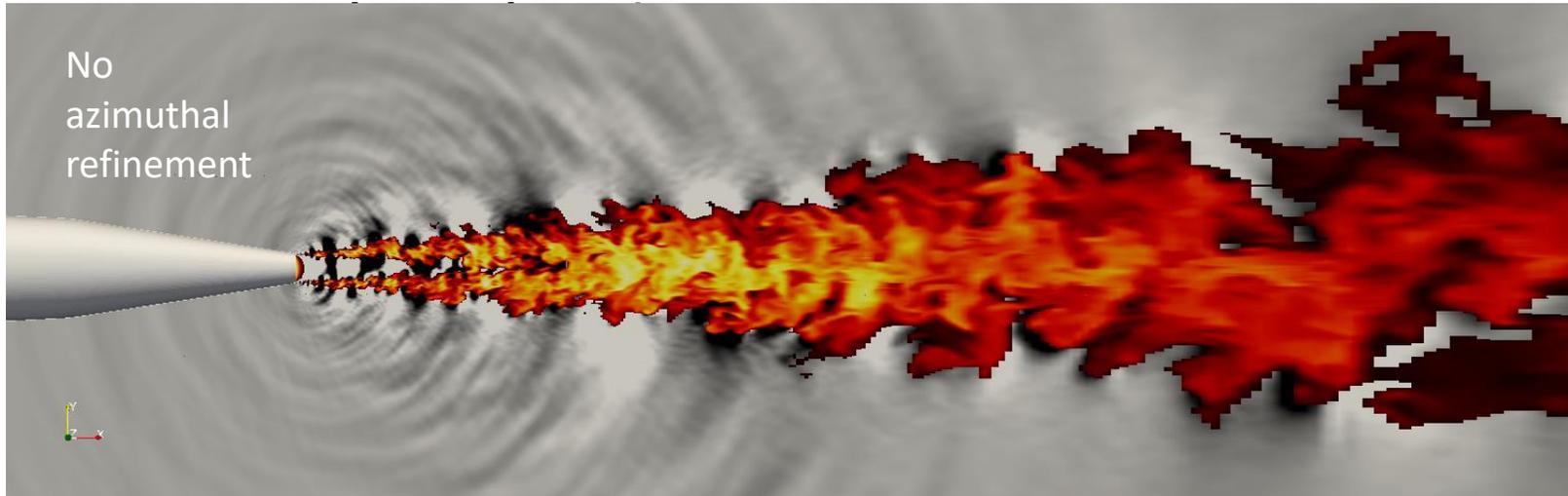
Acoustic fluctuations are so tiny that they are typically measured at log scale ( $\text{dB} = 10 \log_{10}(\text{r.m.s.}(p')^2/p_{\text{ref}}^2)$ ) -- the same as the human hearing works!

# GPU CABARET: CPU VS GPU – 1 Billion challenge

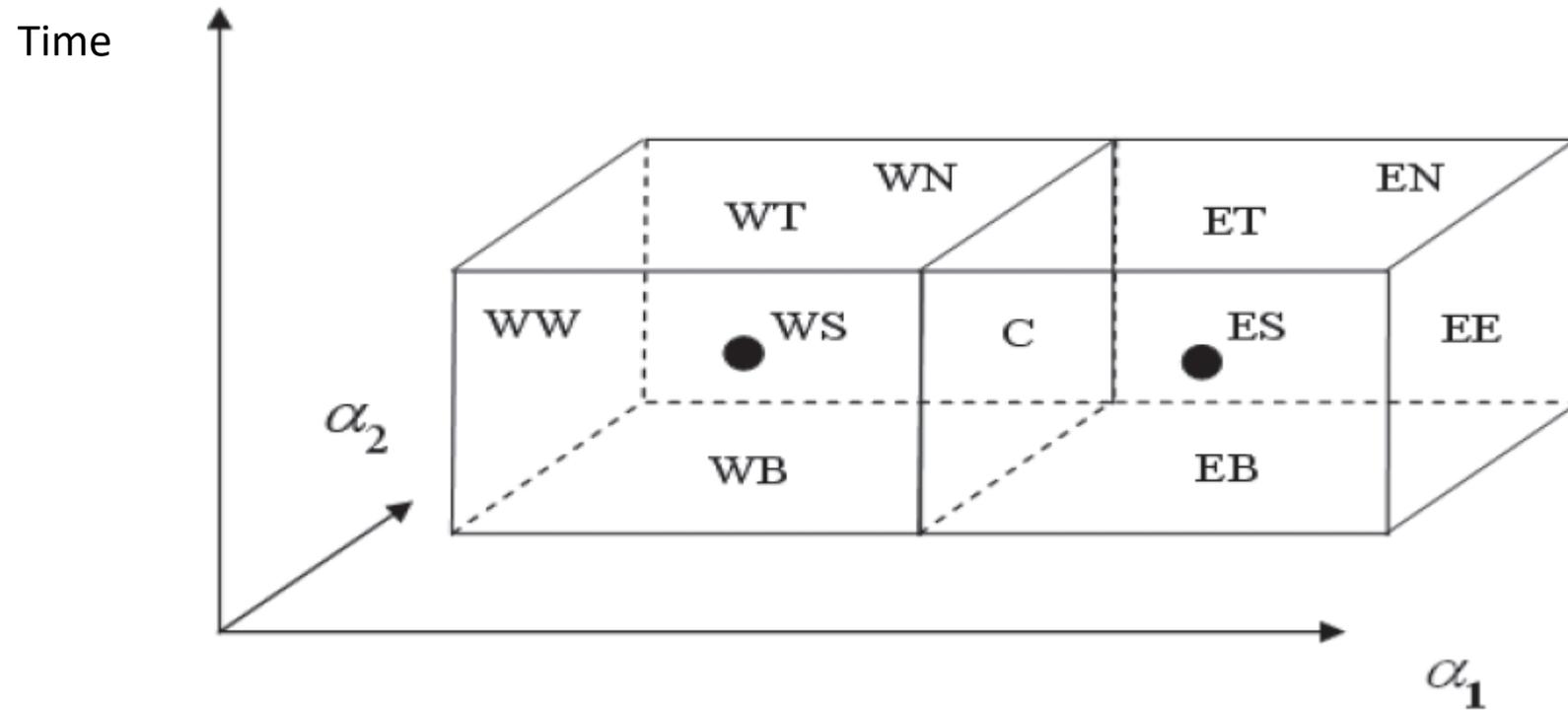
- CPU only:
  - Simple up-scaling: ~50000 cores, ~0.6 MW, many “core hours”
  - Heavily resource bounded (power requirements, communication infrastructure, only possible on a “national” level, multi-million Dollar/Pound/Euro investment, or as part of a large research grant)
- GPU only:
  - Likely memory bounded, GPUs have limited on-board memory
  - Current GPU-CABARET code can handle ~2 million cells per GB
  - Possible GPUs: AMD W9100 32GB / NVIDIA K80 24GB
    - Both sell for “only” \$5,000 and we need  $1e9/2e6=500$  GB
  - GPUs needed: ~16 W9100 or ~20 K80s == \$80,000-\$100,000
  - This is a similar amount of money spent in a medium-sized company on software licenses only!

# Isothermal Static SILOET Jet

- Two  $20 \times 10^6$  cells grids: effect of the azimuthal resolution

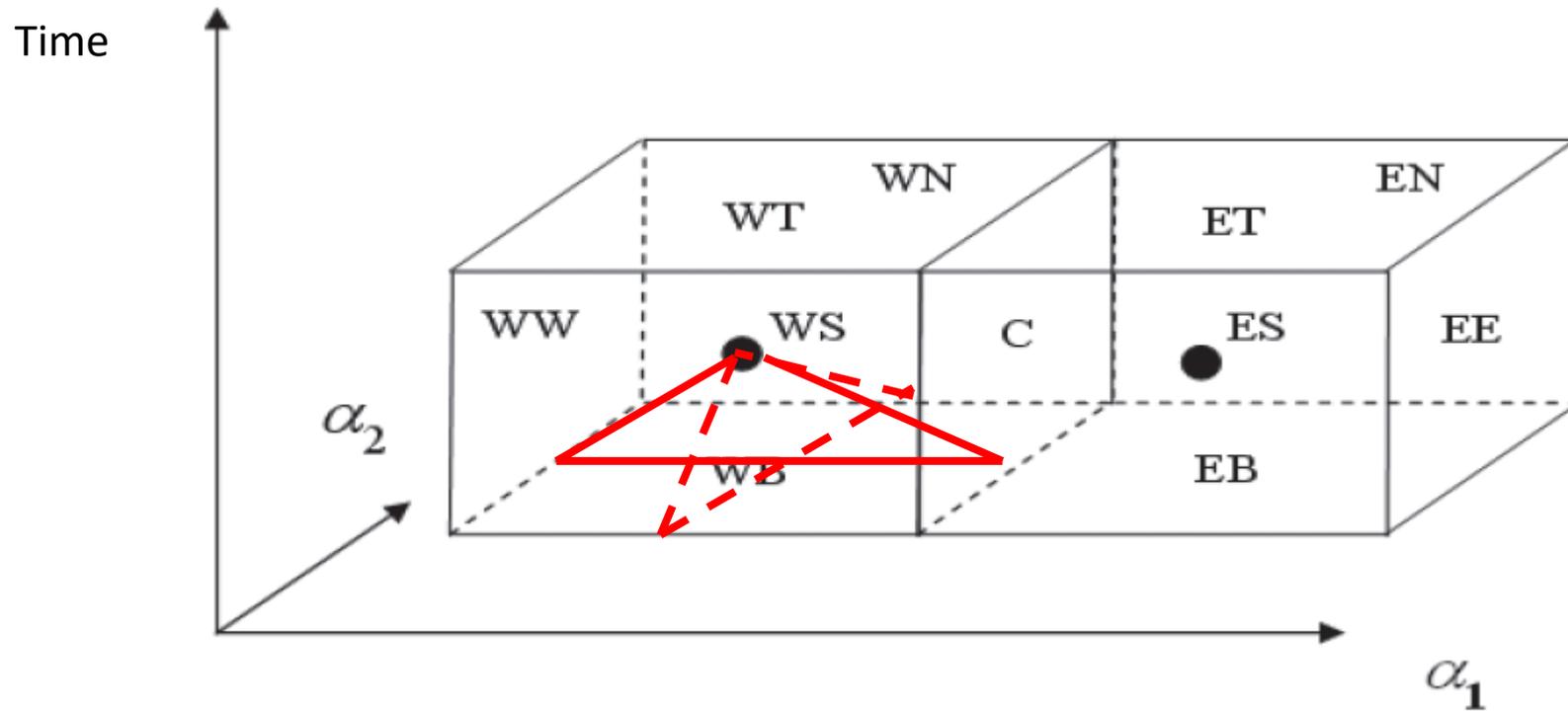


# 2D CABARET space stencil in grid coordinates

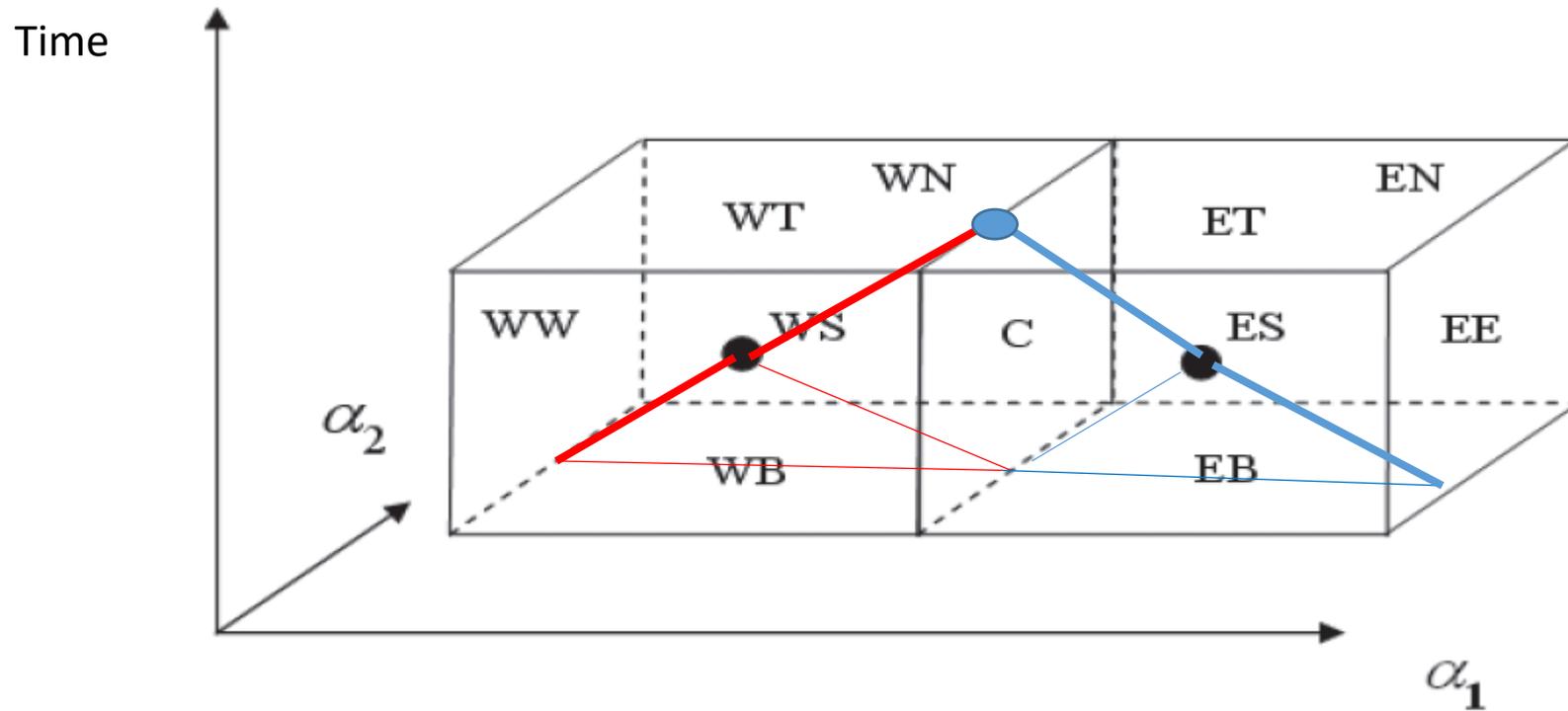


## 2D CABARET: predictor step

cell centre + cell face variables  $\rightarrow$  mid time level variables

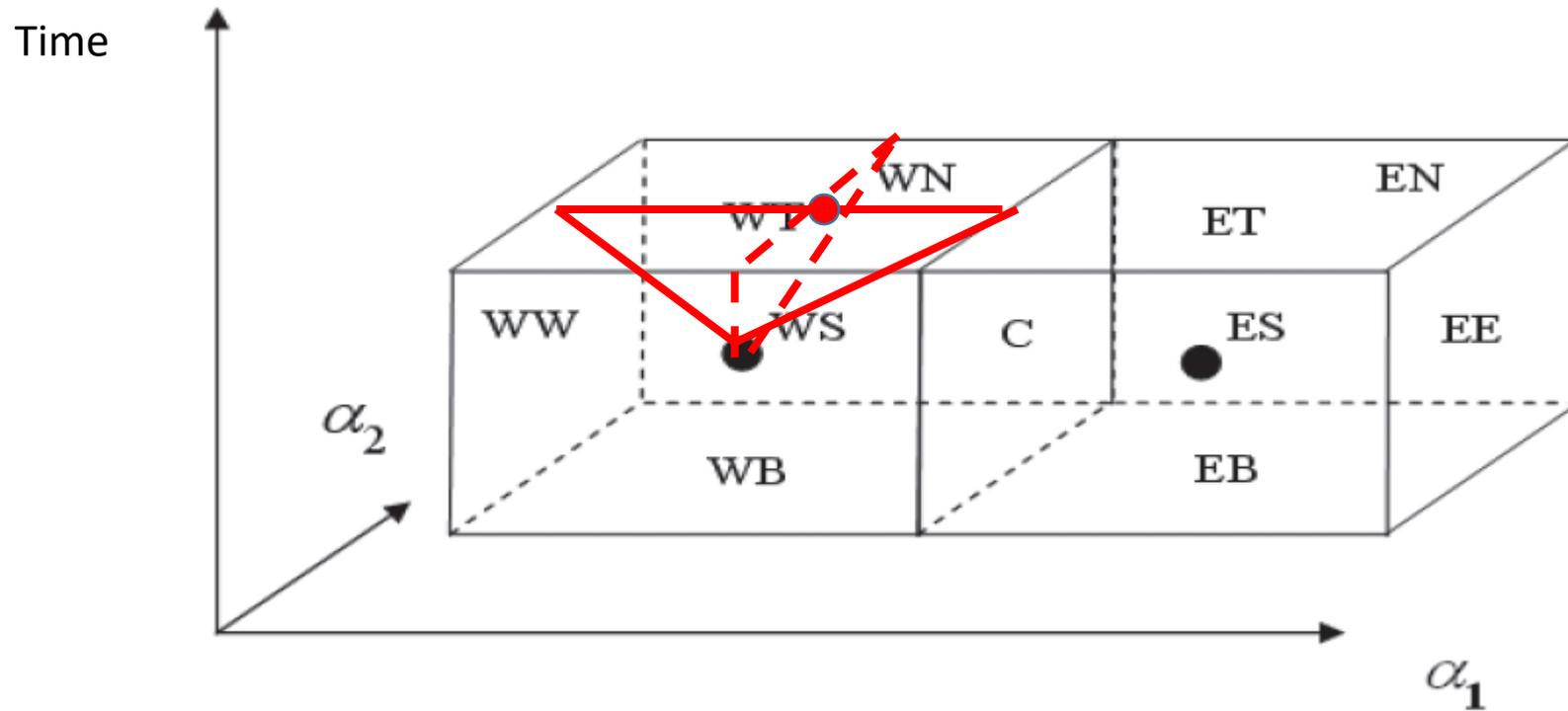


CABARET characteristic decomposition step,  
extrapolation and solution of the Riemann  
problem at the cell face in each grid direction



## 2D CABARET: corrector step

cell centre + cell face variables  $\rightarrow$  new time level variables



# Flux Correction directly based on the maximum principle

$$\tilde{\varphi}_{i+1}^{n+1} = 2 \cdot \psi_{i+1/2}^{n+1/2} - \varphi_i^n$$

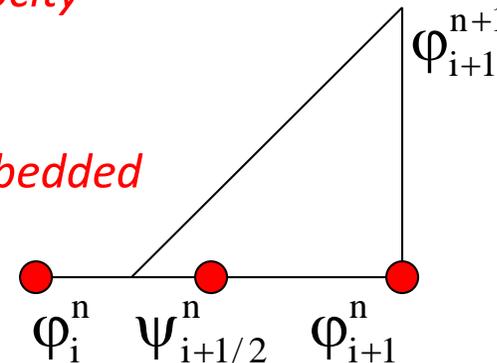
Maximum principle

$$\varphi_{i+1}^{n+1} \leq \max(\varphi), \quad x \in [x_i, x_{i+1}], t = t_n$$

$$\varphi_{i+1}^{n+1} \geq \min(\varphi), \quad x \in [x_i, x_{i+1}], t = t_n$$

*MILES approach: the capability of capturing sharp velocity gradients for smallest resolved scales emulates high-wavenumber end of the inertial subrange region characterised by thin filaments of intense vorticity embedded in a weak background vorticity*

Fureby et al, Phys. Fluids 1997



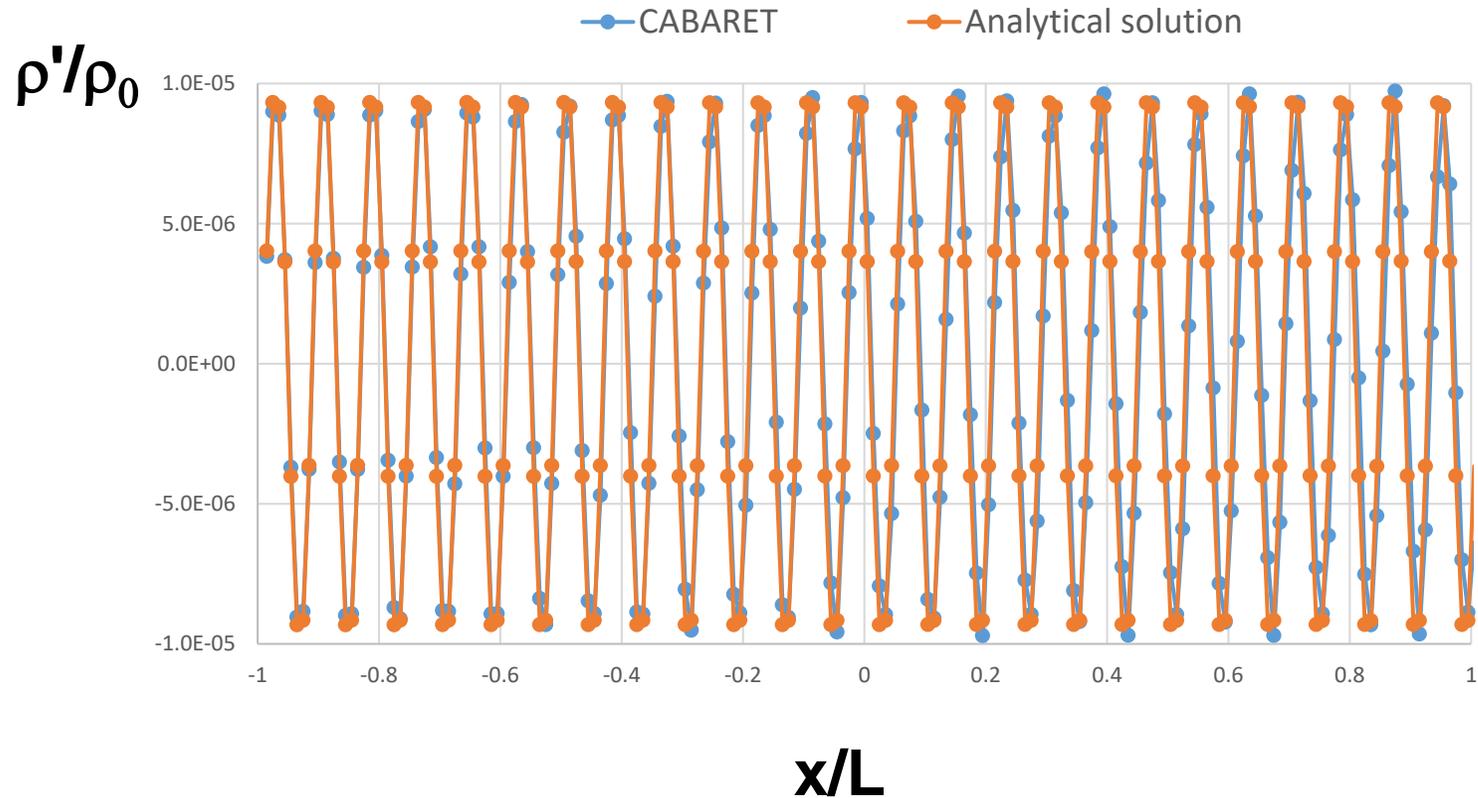
We constrain the solution so that

$$\varphi_{i+1}^{n+1} \leq \max_{x \in [x_i, x_{i+1}], t = t_n} (\varphi) = \left[ \max(\varphi) \right]_{i+1/2}^n = \max(\varphi_{i+1}^n, \psi_{i+1/2}^n, \varphi_i^n)$$

$$\varphi_{i+1}^{n+1} \geq \min_{x \in [x_i, x_{i+1}], t = t_n} (\varphi) = \left[ \min(\varphi) \right]_{i+1/2}^n = \min(\varphi_{i+1}^n, \psi_{i+1/2}^n, \varphi_i^n)$$

Planar wave (1D), 8ppw, CFL=0.1

New relaxed flux correction ( $\varepsilon=0.4$ ), dispersion improved CABARET



Planar wave (1D), non-uniform grid,  $CFL_{\max}=0.8$   
New relaxed flux correction ( $\varepsilon=0.4$ ), dispersion  
improved CABARET

