

# **A COMBINED FEM/BEM DISCRETE NUMERICAL METHOD FOR SOLVING EXTERNAL SCATTERING PROBLEMS IN ACOUSTICS**

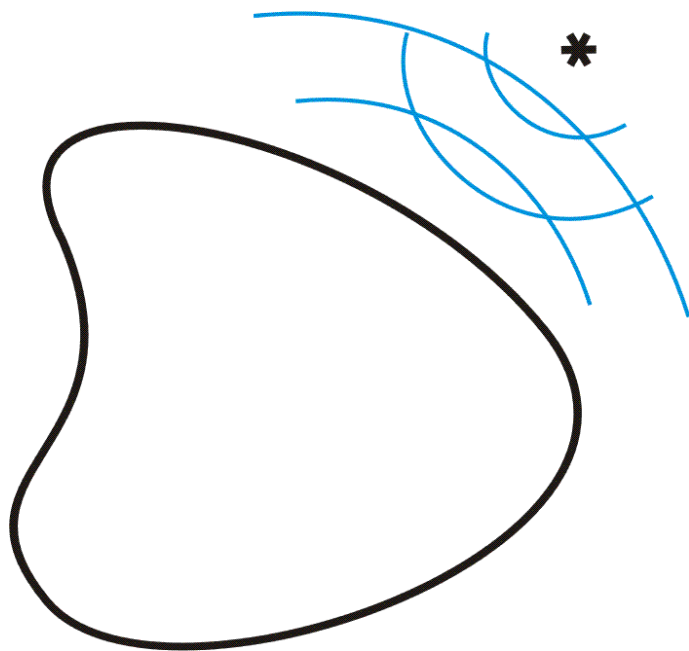
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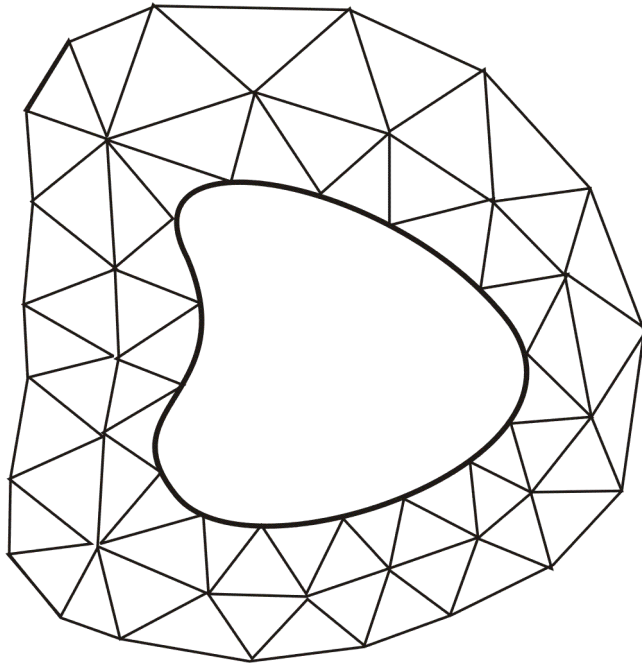
# What we would like to solve?

An external diffraction problem (stationary or non-stationary)  
with Neumann boundary conditions and some sources



# The main dilemma: FEM or BEM?

## FEM: Finite Elements Method



Equation solved: 
$$\Delta p - \frac{1}{c^2} \frac{\partial^2 p}{\partial t^2} = f(r, t)$$

Difficulties:

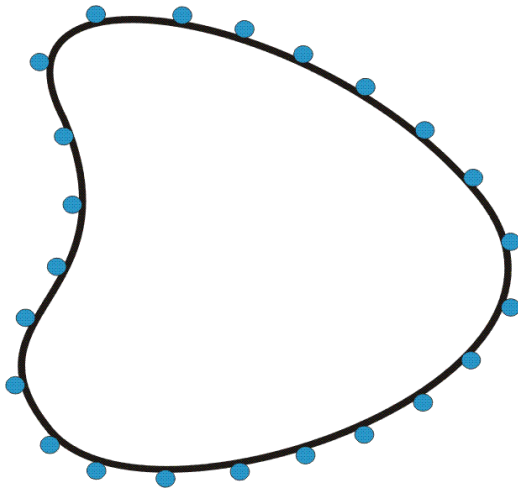
1. Too many nodes
2. Something should be done at infinity (evacuation of waves), say PML

## BEM: Boundary Elements Method

Equation solved:

$$\frac{1}{2}u(r) = \int_s [G(r, r')f(r') - u(r')\partial_n G(r, r')]ds + F(r)$$

(Direct Kirchhoff formulation,  
after application Green's theorem to  
Helmholtz equation)



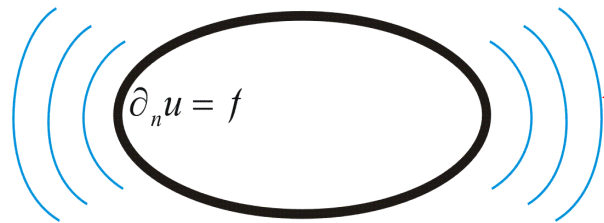
Benefits:

1. No problems at infinity
2. Relatively low amount of nodes
3. No dispersion

Problems:

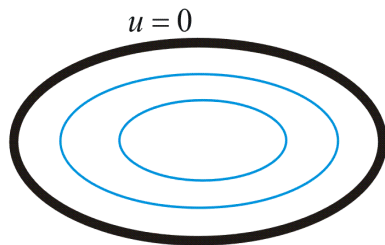
1. Singular and oversingular integrals,  
problems with accuracy
2. Spurious resonances  
(painful artifacts, see below)
3. Dense matrices

# Spurious frequencies (in BEM) and how to kill them



The problem we solve

on spurious frequencies the matrices become ill-conditioned



Spurious frequencies are the eigenvalues of the **dual** problem

The remedy: CFIE (Burton-Miller) approach

$$\begin{aligned}
 & \frac{1}{2}u(r) = \int_S [G(r, r')f(r') - u(r')\partial_{n'}G(r, r')]ds \\
 + & \frac{1}{2}f(r) = \int_S [\partial_n G(r, r')f(r') - u(r')\partial_{n'n}^2 G(r, r')]ds
 \end{aligned}$$

$\nu \partial_n$   
 any complex

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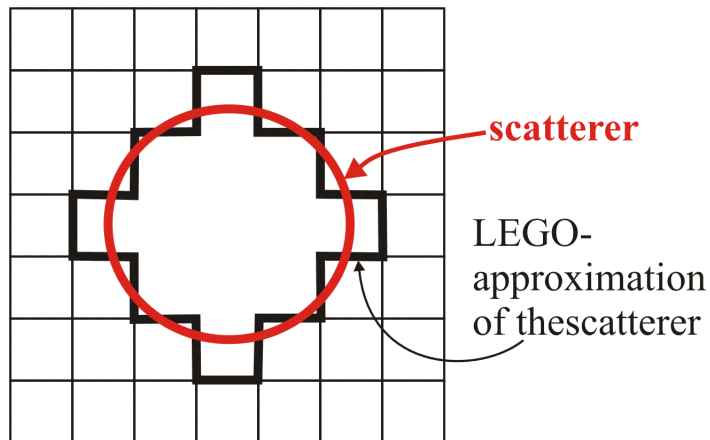

$$\frac{1}{2}u(r) + \int_S [u(r')\partial_{n'}G(r, r') - \nu u(r')\partial_{n'n}^2 G(r, r')]ds = \frac{\nu}{2}f(r) + \int_S [G(r, r')f(r') - \nu \partial_n G(r, r')f(r')]ds$$

painful singularity

# New method. **Step 1: LEGO-BEM** (BAE – Boundary Algebraic Equations)

The main problem of BEM comes from its continuous nature.

If we consider discrete formulation from the very beginning, the Green's function is not singular (and there should be no integration)



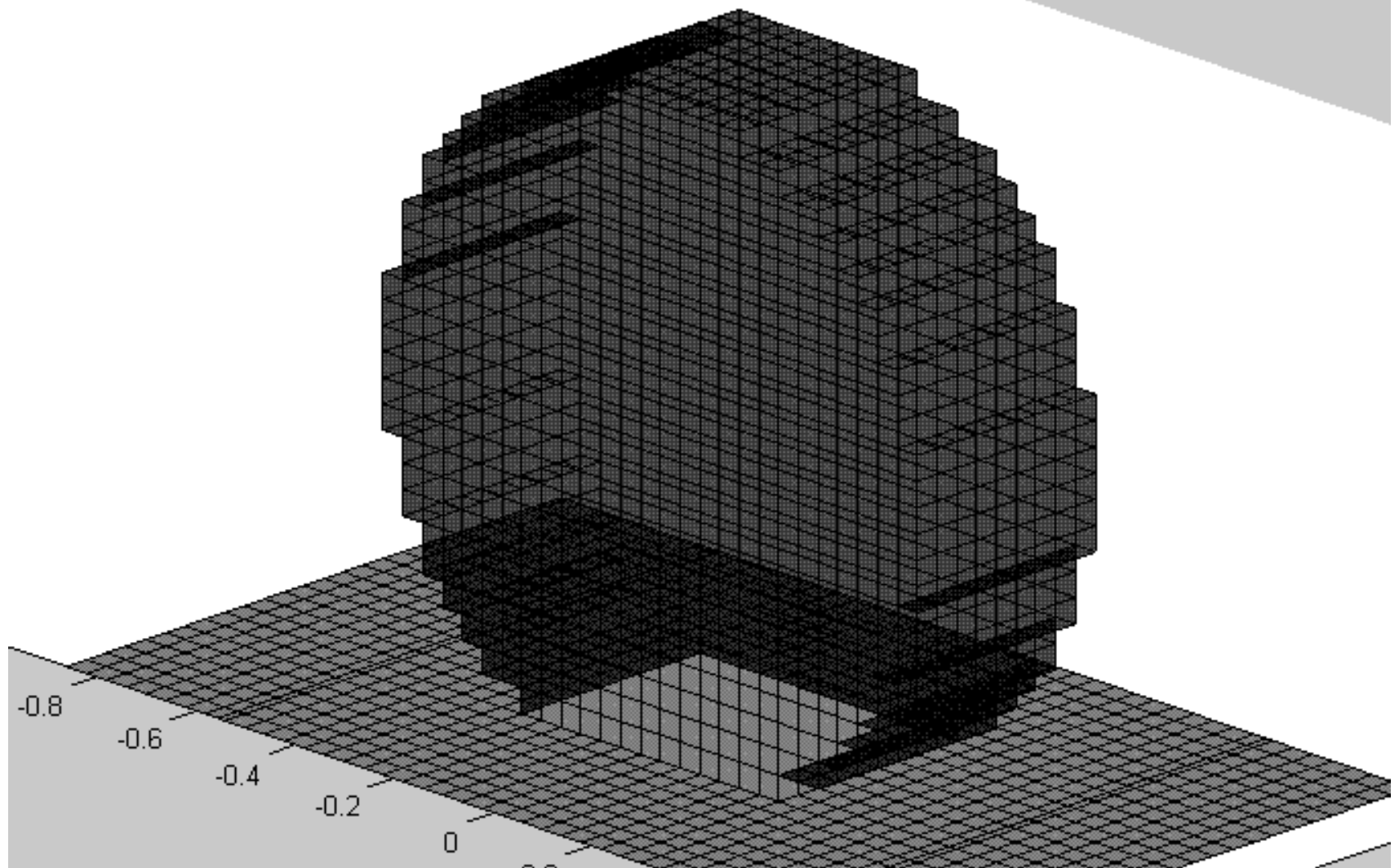
Our aim is to convert  
Burton-Miller (CFIE) technique to the  
discrete form

**The idea itself appears  
to be not new ☹**

**But we were the first  
to implement an analog of CFIE!**

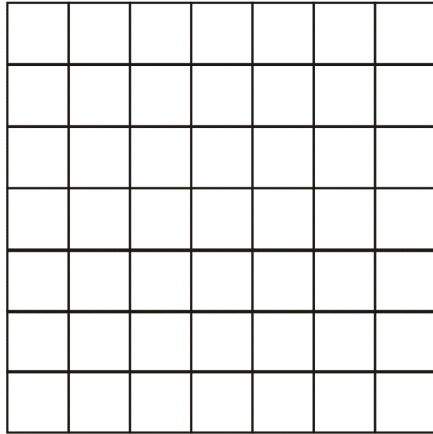
C.Saltzer, 1958;  
P-G. Martinsson et al, 2009, 2010  
I.Tsukerman, 2005, 2006, 2008, 2011

# An example of LEGO mesh



# Formalism of the method

uniform mesh

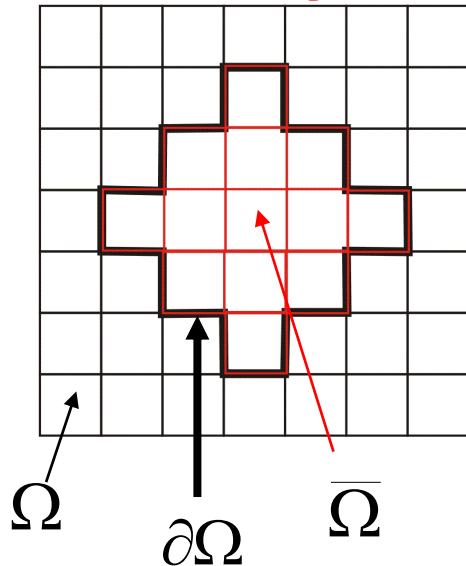


Equation to solve  $\Delta u + k^2 u = f$

Approximate form  $\beta_{i,j} u_j = f_i$

Green's function  $\beta_{i,j} G_{j,m} = \delta_{m,i}$

black mesh - outer problem  
red mesh - inner problem



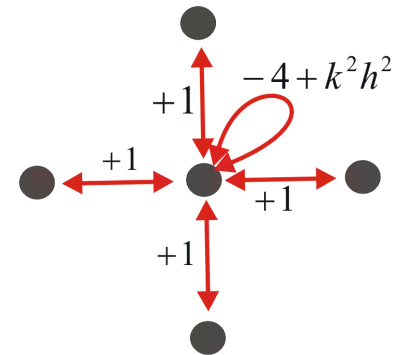
Split  $\beta_{i,j}$  between inner and outer mesh:

$$\beta_{i,j} = \beta_{i,j}^I + \beta_{i,j}^O$$

↑inner
 ↑outer

Both should be symmetrical:  $\beta_{i,j}^I = \beta_{j,i}^I$        $\beta_{i,j}^O = \beta_{j,i}^O$

example:  
5 point scheme  
non-zero values  
of  $\beta_{i,j}$





FEM equation we solve:

$$\sum_{j \in \Omega} \beta_{i,j}^o u_j = f_i, \quad f_i \neq 0 \quad \text{only for } i \in \partial\Omega$$

in operator form:  $\beta^{oo} u^o = f^o$  outer domain  $\Omega$

BAE (LEGO-BEM) formulation of the problem

$$\sum_{i \in \Omega} \sum_{j \in \partial\Omega} G_{m,i} \beta_{i,j}^o u_j = \sum_{j \in \partial\Omega} G_{m,j} f_j \quad m \in \partial\Omega$$

in operator form:  $G^{BO} \beta^{OB} u^B = G^{BB} f^B$  boundary nodes  $\partial\Omega$

It may be not clear from the first glance, but this is a direct analogue of the Kirchhoff BEM formulation!

# Interpretation of the boundary equation as the NtD (Neumann to Dirichlet) operator

$$G^{BO} \beta^{OB} u^B = G^{BB} f^B$$

$$f^B = \beta^{BO} u^O$$

Neumann boundary operator,  
discrete analogue of  $\partial_n u$

$$u^B = (G^{BO} \beta^{OB})^{-1} G^{BB} (\beta^{BO} u^O)$$

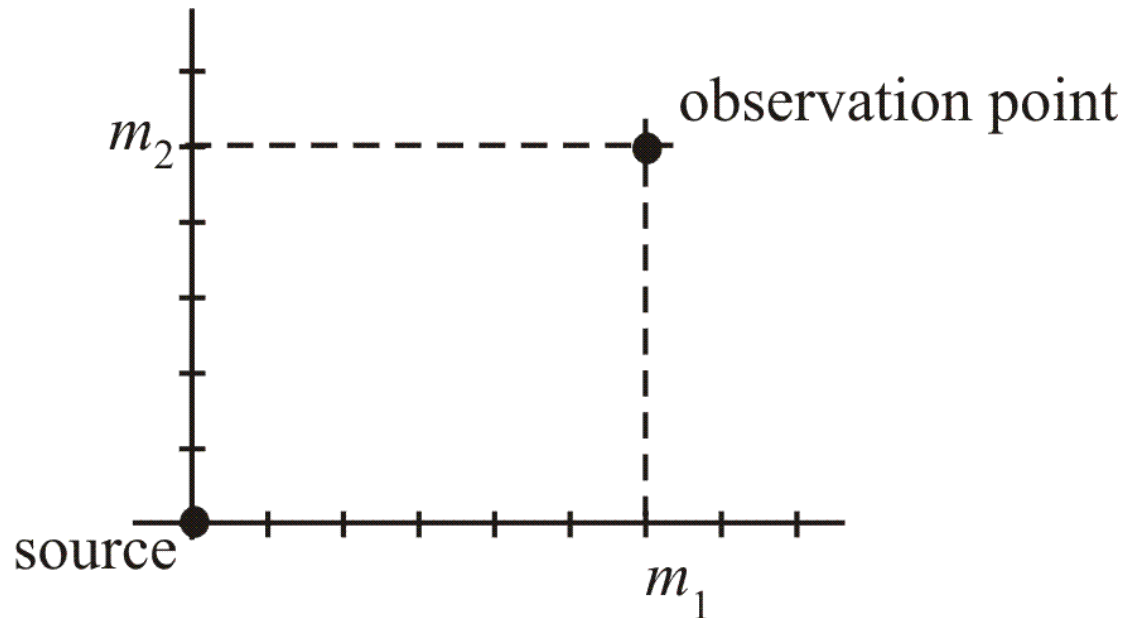
Dirichlet boundary data

NtD operator

Neumann boundary data

This relation guarantees that there are no waves coming from infinity!

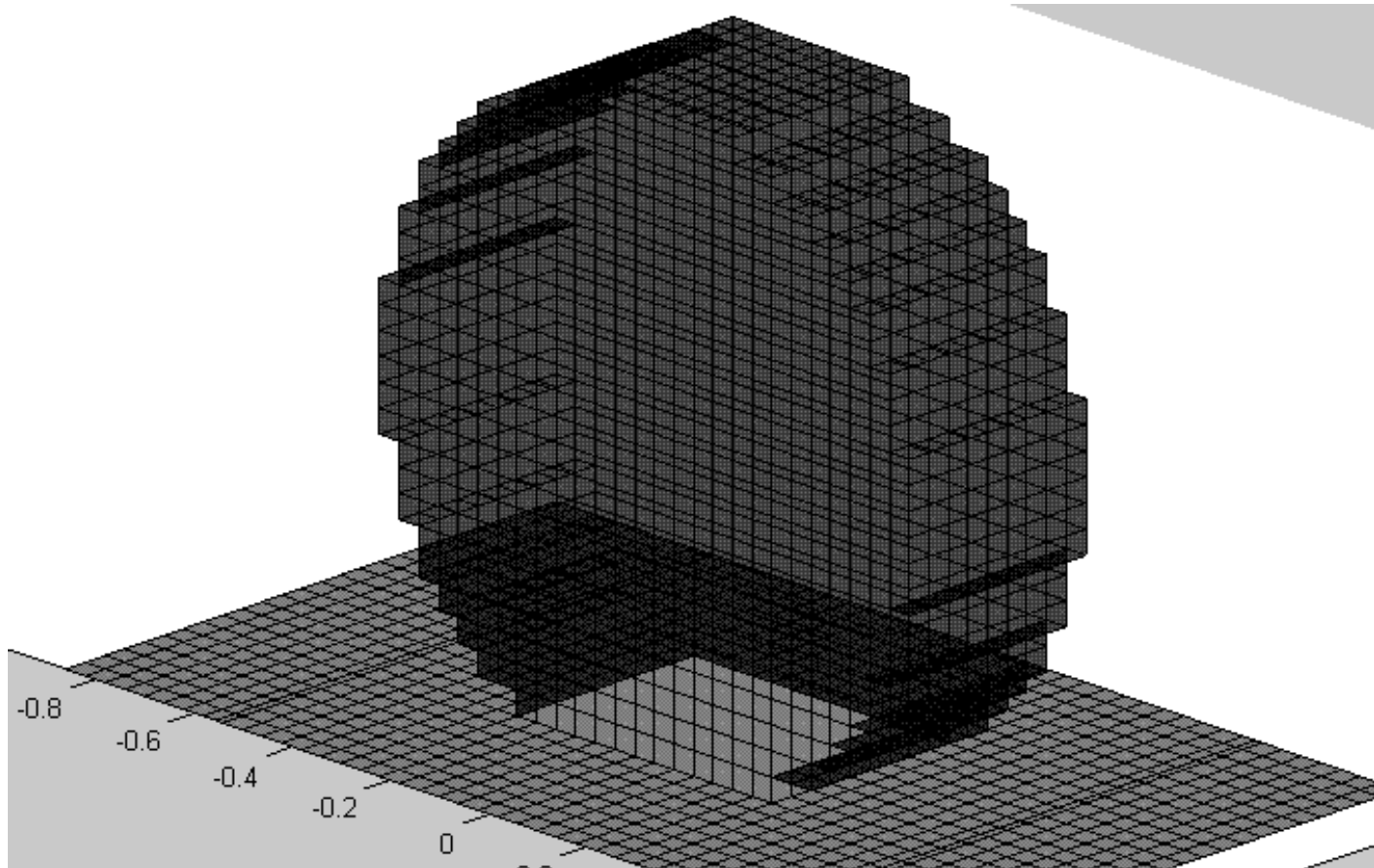
# The price we pay: computation (tabulation) of the Green's function



$$G(m_1, m_2) = \frac{h^2}{4\pi i} \int_{-\pi}^{\pi} \frac{\exp\{im_1\xi + i |m_2| \Xi(\xi)\}}{\Xi(\xi)} d\xi$$

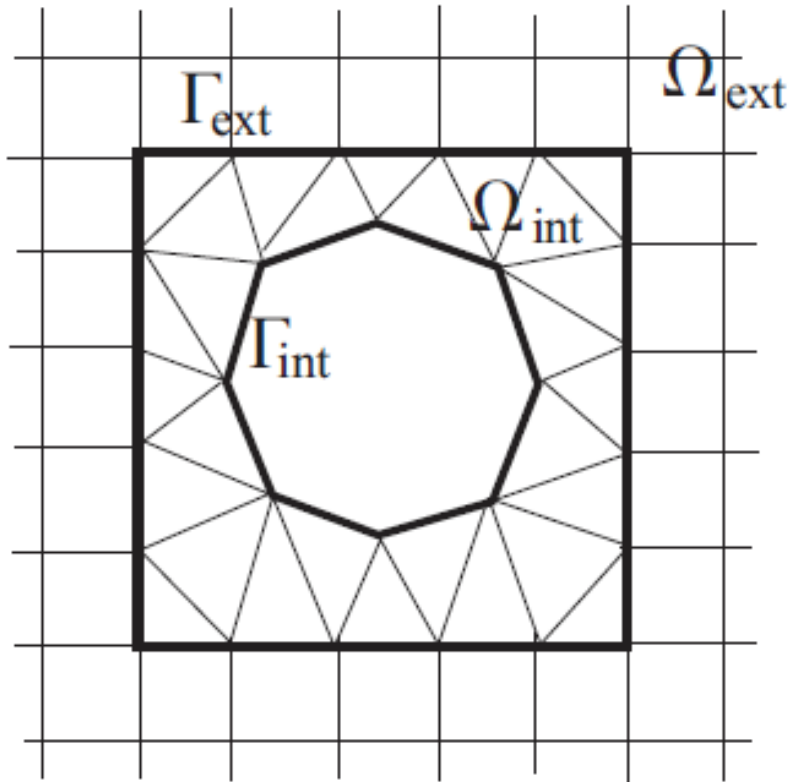
$$\Xi(\xi) = \arccos(2 - \cos \xi - h^2 k^2 / 2)$$

In principle, we can use BAE, but the accuracy is poor due to rough (“LEGO”) boundary modeling!



# New method, Step 2:

## Combining FEM with BAE

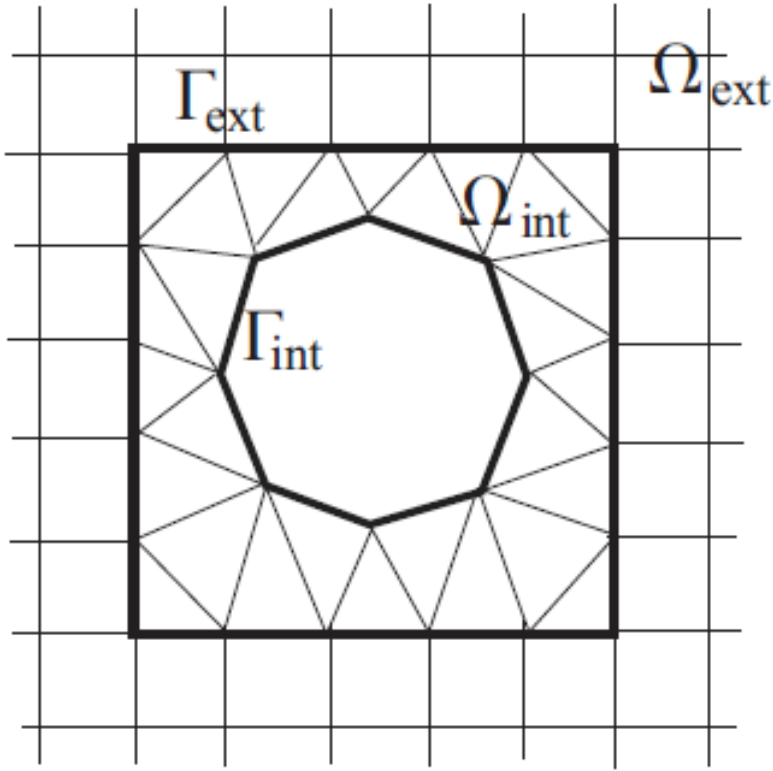


We apply BAE to  $\Omega_{ext}$

to find the link between  
Neumann and Dirichlet  
data on  $\Gamma_{ext}$

Then we apply FEM to  $\Omega_{int}$   
using BAE as a boundary  
condition on  $\Gamma_{ext}$

# Formalism of the method



$$I : \Omega_{\text{int}}$$

$$O : \Omega_{\text{ext}}$$

$$i : \Omega_{\text{int}} \setminus \Gamma_{\text{ext}}$$

$$B : \Gamma_{\text{ext}}$$

FEM on  $\Omega_{\text{int}}$  :

$$\beta^{iI} u^I = f^i \quad (*)$$

Boundary condition on  $\Gamma_{\text{ext}}$  :

Note that  $(\beta^{BO} + \beta^{BI})u = 0$  on  $\Gamma_{\text{ext}}$  (all sources are on  $\Gamma_{\text{int}}$ )

thus  $\beta^{BO} u^O = -\beta^{BI} u^I$  (Neumann boundary data)

$$u^B = -(G^{BO} \beta^{OB})^{-1} G^{BB} (\beta^{BI} u^I) \quad (**)$$

System (\*), (\*\*) is what we solve!

Example of mesh  $\Omega_{\text{int}}$

